


Communication

Sol-Gel Synthesized Amorphous $(\text{In}_x\text{Ga}_{1-x})_2\text{O}_3$ for UV Photodetection with High Responsivity

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Abstract: $\beta\text{-Ga}_2\text{O}_3$ photodetectors have the advantages of low dark current and strong radiation resistance in UV detection. However, the limited photocurrent has restricted their applications. Herein, MSM UV photodetectors based on $(\text{In}_x\text{Ga}_{1-x})_2\text{O}_3$ ($x = 0, 0.1, 0.2, 0.3$) by a sol-gel method were fabricated and studied. The doping of indium ions in Ga_2O_3 leads to lattice distortion and promotes the formation of oxygen vacancies. The oxygen vacancies in $(\text{In}_x\text{Ga}_{1-x})_2\text{O}_3$ can be modulated by various proportions of indium, and the increased oxygen vacancies contribute to the enhancement of electron concentration. The results show that the amorphous $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ photodetector exhibited improved performances, including a high light-to-dark current ratio (2.8×10^3) and high responsivity (739.2 A/W). This work provides a promising semiconductor material $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ for high-performance MSM UV photodetectors.

Keywords: ultraviolet photodetector; high responsivity; $(\text{In}_x\text{Ga}_{1-x})_2\text{O}_3$; oxygen vacancies



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1. Introduction

Ultraviolet (UV) photodetectors have attracted widespread attention due to their potential applications in solar spectrum detection, UV warning, flame detection, ozone monitoring and so on [1–4]. Metal oxide semiconductors have attracted considerable attention in the field of ultraviolet photodetectors due to their outstanding electronic and photoelectric properties. $\beta\text{-Ga}_2\text{O}_3$ is a semiconductor with a direct wide bandgap of 4.9 eV, excellent chemical and thermal stability, and a high breakdown field strength of 8 MV/cm, second only to diamond, making it suitable for the fabrication of solar-blind ultraviolet photodetectors [1]. Xu et al. fabricated a solar-blind ultraviolet photodetector based on MSM structure using $\beta\text{-Ga}_2\text{O}_3$ thin film, which exhibits a low dark current of <10 pA at 10 V [5]. The solar-blind ultraviolet photodetector based on $\beta\text{-Ga}_2\text{O}_3$, fabricated by Shen et al., demonstrates a rapid response with a response time and decay time of 0.10 s [6]. However, existing UV photodetectors still have limitations in terms of photocurrent and responsivity. Therefore, enhancing the performance of UV photodetectors has become a focal point of current research [7–11].

Doping is an effective method for improving the performance of UV photodetectors [12–15]. Doping can increase the carrier concentration in materials, thereby enhancing their conductivity and optoelectronic performance [16–20]. In studies, the photocurrent of Ga_2O_3 devices can be improved by doping with In_2O_3 , which increases the device's responsivity [21,22]. By altering the In content in $(\text{In}_x\text{Ga}_{1-x})_2\text{O}_3$ (abbreviated as IGO) thin films, the optical bandgap can be adjusted from 2.9 eV to 4.9 eV, making it useful for the fabrication of ultraviolet and visible-blind photodetectors with different cutoff

wavelengths [23]. Simultaneously, the content of In plays a crucial role in regulating the phase transition process of thin films. There is a dearth of literature on IGO-based UV photodetectors. Yoshihiro Kokubun et al. used a sol-gel method to tune the indium (In) content in polycrystalline indium gallium oxide ($\text{In}_x\text{Ga}_{1-x}\text{O}_3$) and found that increasing the In content led to an increase in lattice constant and a decrease in bandgap [24]. Metal-semiconductor-metal (MSM) UV photodetectors were fabricated using radio frequency (RF) sputtered ($\text{In}_{0.9}\text{Ga}_{0.1}\text{O}_3$) thin films, and it was observed that the device switched from photoconductive to Schottky contact mode with increasing oxygen pressure, with a peak responsivity of 0.31 A/W [23]. Kuan-Yu Chen et al. prepared amorphous indium gallium oxide (InGaO) MSM photodetectors using magnetron sputtering and demonstrated a peak responsivity of 3.83 A/W at 5 V bias and 280 nm illumination [25]. Isa Hatipoglu et al. grew polycrystalline monoclinic indium gallium oxide using metal organic chemical vapor deposition (MOCVD) and tuned the peak responsivity of MSM photodetectors by increasing the indium content incorporated into the Ga_2O_3 lattice, from 0.79 A/W (pure Ga_2O_3) to 319.1 A/W, 66.1 A/W, and 27.7 A/W for ($\text{In}_{0.203}\text{Ga}_{0.797}\text{O}_3$), ($\text{In}_{0.177}\text{Ga}_{0.823}\text{O}_3$), and ($\text{In}_{0.106}\text{Ga}_{0.894}\text{O}_3$), respectively [26]. However, there is no research yet on effectively controlling the film's crystal phase by increasing the In content to make the film transition from $\beta\text{-Ga}_2\text{O}_3$ phase to an amorphous structure and then to In_2O_3 phase. By leveraging the dual effects of crystal phase transition and increased In content, the photodetector's performance can be improved.

In this study, we prepared IGO UV photodetectors based on a metal-semiconductor-metal (MSM) structure using a sol-gel method. By changing the indium content in the film, the crystal phase of the film can be effectively controlled. Taking advantage of the high oxygen vacancy density of amorphous materials, the lifetime of carriers was significantly extended, and the light-to-dark current suppression ratio as well as the responsivity of the device were improved. Among them, devices prepared with $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ thin film showed a high light-to-dark current ratio (2.8×10^3) and high responsivity (739.2 A/W). Additionally, this research explores the working mechanisms and performance-influencing factors of the detectors, providing a theoretical basis for further enhancing the performance of UV photodetectors.

2. Materials and Methods

In this work, IGO thin films were prepared on quartz substrate using a sol-gel method. The specific experimental steps are as follows:

A certain amount of gallium nitrate trihydrate $\text{Ga}(\text{NO}_3)_3 \cdot x\text{H}_2\text{O}$ (A.R. from Aladdin (Shanghai, China, CAS No. 69365-72-6) and indium nitrate tetrahydrate $\text{In}(\text{NO}_3)_3 \cdot x\text{H}_2\text{O}$ (A.R. from Aladdin, CAS No. 207398-97-8) were dissolved in 30 mL of deionized water. The metal ions in the solution were controlled at 0.1 M, and the molar ratios of In and Ga were 0:10, 1:9, 2:8 and 3:7, respectively. The corresponding amounts of $\text{Ga}(\text{NO}_3)_3 \cdot x\text{H}_2\text{O}$ are 3.835 g, 3.452 g, 3.068 g and 2.685 g, and the corresponding amounts of $\text{In}(\text{NO}_3)_3 \cdot x\text{H}_2\text{O}$ are 0 g, 0.451 g, 0.903 g and 1.654 g. Then, while maintaining a constant temperature of 75 °C, 30 μL of polyoxyethylene lauryl ether ($\text{HO}(\text{CH}_2\text{CH}_2\text{O})_n(\text{CH}_2)_{11}\text{CH}_3$, Brij[®]35, from Aladdin, CAS No. 9002-92-0) was added, and the mixture was stirred for 30 min to obtain the IGO precursor solution. The IGO precursor solution was spin-coated on a pre-cleaned quartz substrate at a rotation speed of 3000 rpm for 30 s, followed by heating at 110 °C for 10 min to complete the gelation process, which was repeated 10 times. After the spin coating was completed, the samples were annealed in a furnace at 700 °C for 30 min. By adjusting the In content in the precursor solution to 0%, 10%, 20%, and 30%, four types of thin films were obtained, namely $\beta\text{-Ga}_2\text{O}_3$, $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$, $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, and $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$. Au interdigital electrodes were fabricated on the thin films using photolithography and radio frequency magnetron sputtering techniques (Vacuum degree was 7×10^{-3} Pa, sputter pressure was 2.4×10^0 Pa, sputtering power was 60 W, sputtering time was 6 min, and turntable speed was 10 rpm). From top of view, the finger width and the spacing of the

interdigital electrodes were both 20 μm , and the effective detection area of the device was 0.38 mm^2 . The schematic diagram of the IGO MSM photodetector is shown in Figure 1.

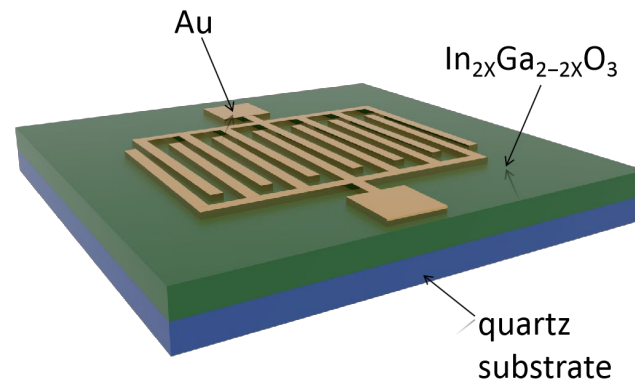


Figure 1. Schematic structure of a metal-semiconductor-metal (MSM) photodetector.

The X-ray diffraction patterns of the thin films were observed using $\text{Cu K}\alpha$ radiation on a Shimadzu XRD-6000 diffractometer (Shimadzu, Columbia, MD, USA). X-ray photoelectron spectroscopy (XPS) measurements were carried out using an ESCALAB 250 photoelectron spectrometer (Thermo Fischer, Waltham, MA, USA). Absorption coefficient spectra were measured using a Shimadzu UV-3600 Pharma Spec UV-Vis spectrophotometer. Current-voltage characteristics were measured using a Keithley 2450 source meter (Keithley Instruments, Cleveland, OH, USA). The photodetector's light response was measured using a monochromator equipped with a 30 W xenon lamp.

3. Results and Discussion

After testing the performance of the prepared UV detectors based on $\beta\text{-Ga}_2\text{O}_3$, $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$, $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ and $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$, it was found that the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ detector had the best performance. The SEM images of the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ film are presented here, and there is no obvious difference among the SEM images of all films. Figure 2a illustrates the top-view scanning electron microscopy (SEM) images of the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ film at different magnification scales. Nanoparticles of $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, with an average diameter of approximately 15 nm, are densely packed on a quartz substrate, forming a compact film. Figure 2b is the cross-view SEM image of the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ film. The thickness is about 84 nm. These results indicate the successful preparation of a flat $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ film.

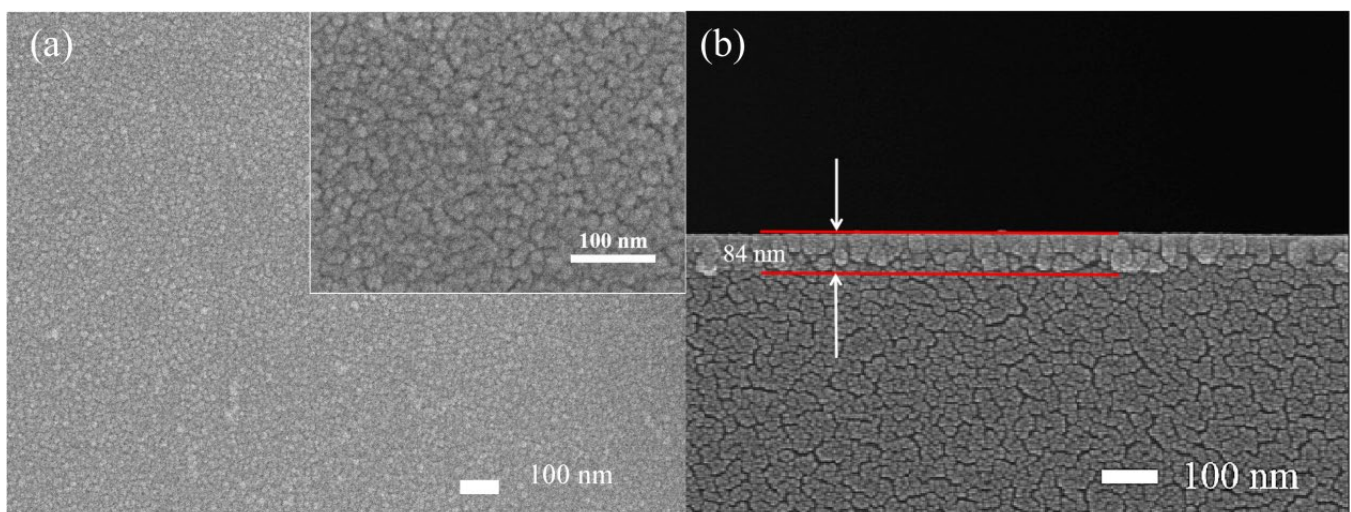


Figure 2. (a) Top-view SEM image of the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ film, the inset is the high magnification. (b) Cross-view SEM image of the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ film.

Figure 3 displays the X-ray diffraction patterns of the IGO thin films. The β -Ga₂O₃ film with an In doping concentration of 0% is crystalline, and the observed (−110), (−202), (111), (−311), (−312) and (−421) main peaks are consistent with the β -Ga₂O₃ phase, which match well with the JCPDS card number 41-1103 [27]. As the In doping concentration increases, the peak intensities decrease, the full width at half maximum (FWHM) broadens, and the surface crystal quality deteriorates. In the In_{0.6}Ga_{1.4}O₃ films, the observed (222), (321), (440), and (622) main peaks are consistent with the In₂O₃ phase with the JCPDS card number 06-0416 [28]. As the In doping concentration increases from 0% to 30%, a phase transition occurs from the β -Ga₂O₃ phase to the In₂O₃ phase [21,29].

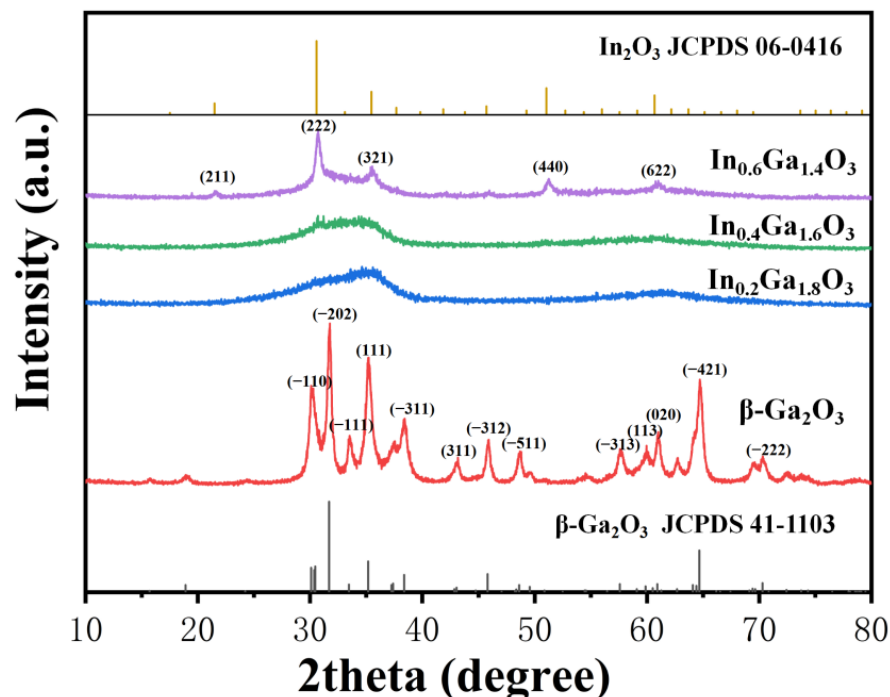


Figure 3. XRD pattern of the β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃ and In_{0.6}Ga_{1.4}O₃ films.

Figure 4a presents the optical absorption spectra of the β -Ga₂O₃ and IGO films. All films exhibit significant light absorption in the ultraviolet region. As the In content increases, the absorbance of the films increases, indicating enhanced ultraviolet light absorption by the films. Additionally, it can be roughly observed that as the In content increases, the absorption edge of the curve experiences a redshift. The following formula, Tauc equation, is used to calculate the material's bandgap width based on the absorption spectra, providing a detailed comparison of how the increase in In content affects the material's bandgap width [30]:

$$\alpha = A \frac{(h\nu - E_g)^{\frac{n}{2}}}{h\nu} \quad (1)$$

where α is absorption coefficient, A is a constant, h is Planck's constant, ν is the frequency of the light, E_g is the bandgap energy and the exponent n is denoted as the nature of transitions.

As can be seen from Figure 4b, the approximate bandgap width of the thin film materials continuously decreases with increasing In doping concentration. The energy value of the films decreases progressively, with the β -Ga₂O₃ film having 4.84 eV, the In_{0.2}Ga_{1.8}O₃ film possessing 4.46 eV, the In_{0.4}Ga_{1.6}O₃ film exhibiting 4.28 eV, and finally, the In_{0.6}Ga_{1.4}O₃ film featuring 4.18 eV, the trend is the same in other researchers' work [21,26,31].

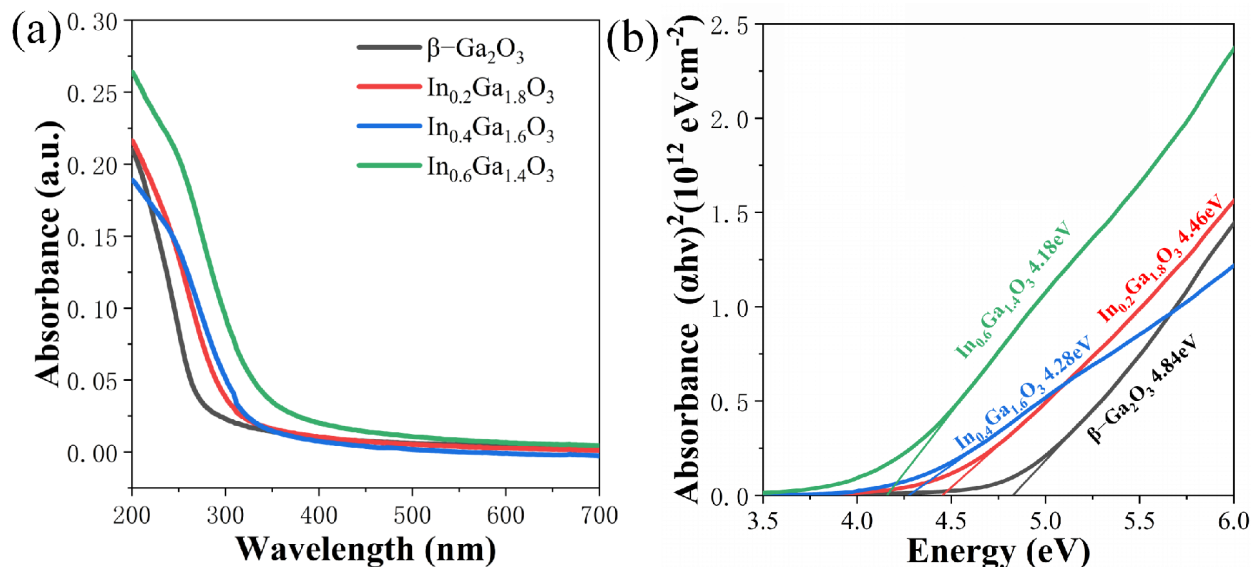


Figure 4. (a) Absorbance spectra of β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃ and In_{0.6}Ga_{1.4}O₃ films. (b) Bandgap widths of β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃, and In_{0.6}Ga_{1.4}O₃ films.

The chemical composition of IGO films was investigated using X-ray photoelectron spectroscopy (XPS). Figure 5 shows the full XPS spectrum, and Figure 6 presents the O1s spectra for the four types of thin films. The actual chemical compositions of the β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃ and In_{0.6}Ga_{1.4}O₃ films are Ga_{0.2938}O_{0.7062}, In_{0.0747}Ga_{0.2796}O_{0.6457}, In_{0.1256}Ga_{0.2624}O_{0.6124} and In_{0.2092}Ga_{0.1761}O_{0.6147} from XPS analysis. With In doping, the amount of In increases from 0.0747, 0.1256 to 0.2093, which is approximately equal to 1:2:3. The amount of Ga decreases from 0.2938, 0.2796, 0.2624 to 0.1761, which is approximately equal to 1:0.9:0.8:0.7. To distinguish the XPS measurements, Gaussian fitting was used to divide the typical O1s peak into three peaks. The deconvoluted peak near 530 eV corresponds to lattice oxygen, while the deconvoluted peak near 531 eV corresponds to defect oxygen [21]. The deconvoluted peak near 532 eV corresponds to adsorbed oxygen [32]. By comparing the four thin film devices, it was found that as the In doping concentration increases, the proportion of lattice oxygen decreases, and the proportion of defect oxygen increases. In the O1s spectrum of In_{0.6}Ga_{1.4}O₃, the proportion of lattice oxygen is smaller than that of defect oxygen. Combined with the XRD characterization of In_{0.6}Ga_{1.4}O₃, it indicates that as the In content increases to a certain extent, In aggregation occurs on the surface of In_{0.6}Ga_{1.4}O₃, altering the chemical environment of the material surface.

Figure 7 displays the I-V curves of β -Ga₂O₃ and IGO devices under dark and UV illumination (β -Ga₂O₃ @260 nm, 4.8 μ W/cm², and IGO @260 and 270 nm, 6.8 and 9.7 μ W/cm²) at biases ranging from -5 to 5 V. By comparing the photocurrents of the β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃, and In_{0.6}Ga_{1.4}O₃ devices, it reveals that as the In doping increases, the photocurrent gradually increases, and the increase is quite substantial. At a 5V bias, the photocurrents of the devices with In doping concentrations of 0%, 10%, 20%, and 30% are 0.2 nA, 53.2 nA, 23.4 μ A, and 758.5 μ A, respectively. It can be observed that as the In doping concentration increases, the rate of increase in the photocurrent of the device slows down. At a 5V bias, the dark currents of the β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃, and In_{0.6}Ga_{1.4}O₃ devices are 0.05 nA, 0.1 nA, 8.2 nA, and 3.0 μ A, respectively. It can be seen that as In doping concentration increases, the rate of increase in the dark current of the device accelerates.

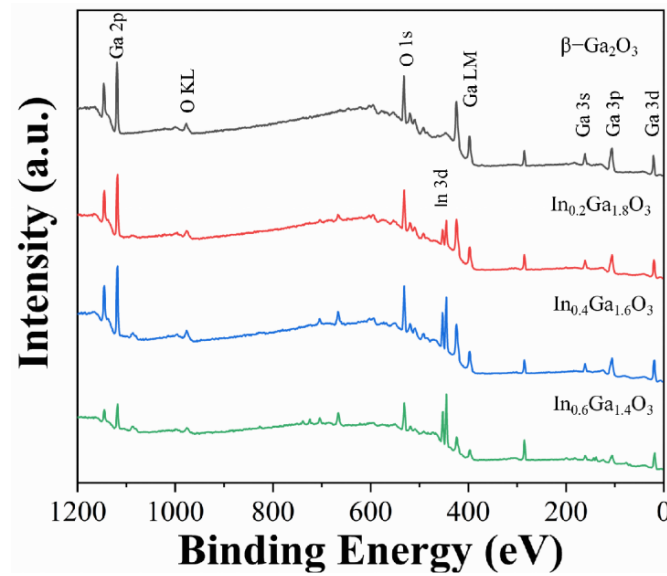


Figure 5. Survey spectra of $\beta\text{-Ga}_2\text{O}_3$, $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$, $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, and $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$ films.

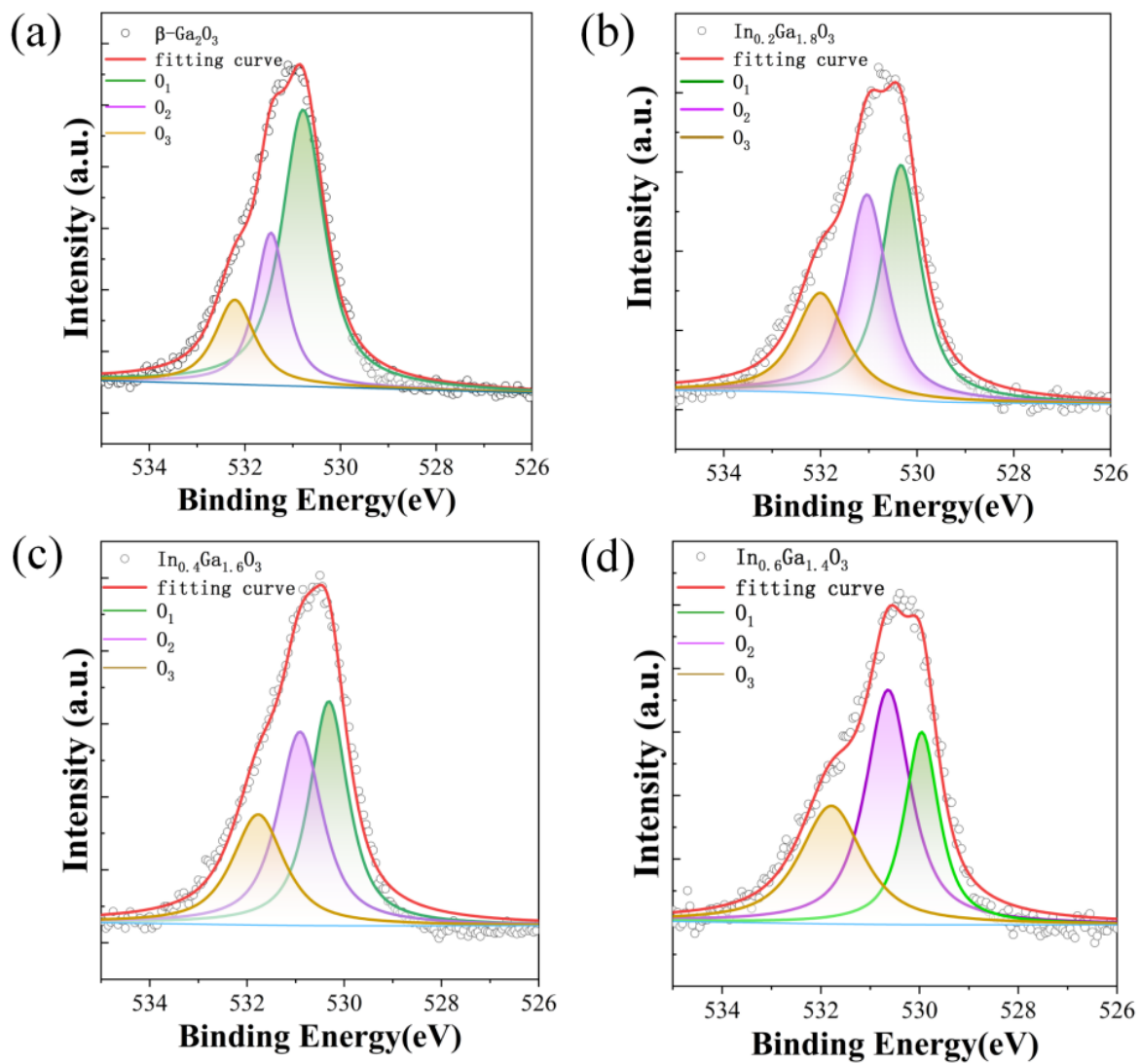


Figure 6. X-ray photoelectron spectroscopy (XPS) spectra of O1s of (a) $\beta\text{-Ga}_2\text{O}_3$, (b) $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$, (c) $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, and (d) $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$ films.

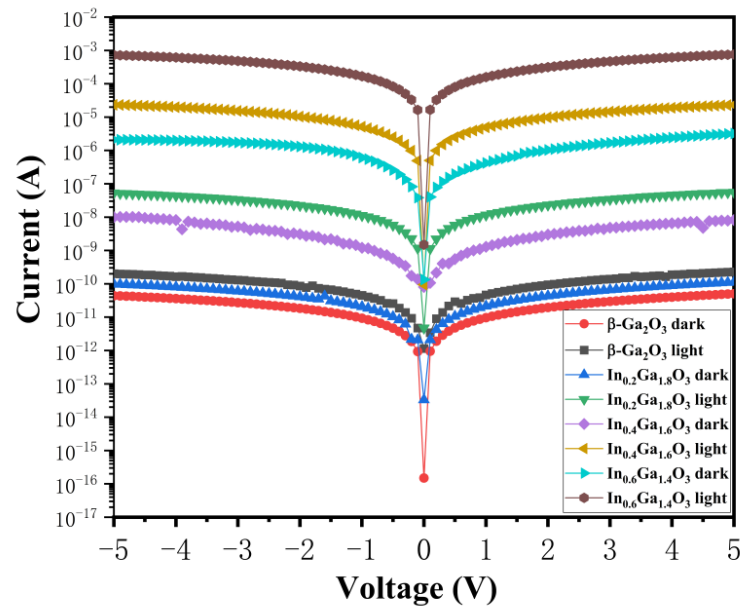


Figure 7. I-V characteristics of $\beta\text{-Ga}_2\text{O}_3$ and IGO devices under dark and UV illumination.

Based on the I-V curves, it was found that as the In doping concentration increases, the photocurrent and dark current of the devices exhibit different rates of increase. To illustrate this trend, Figure 8 shows the light current, dark current, and light-to-dark current ratio of the four devices at a 5V bias. Among them, the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ thin film device demonstrates the highest ratio, at 2.8×10^3 , which is an order of magnitude higher than the other two devices.

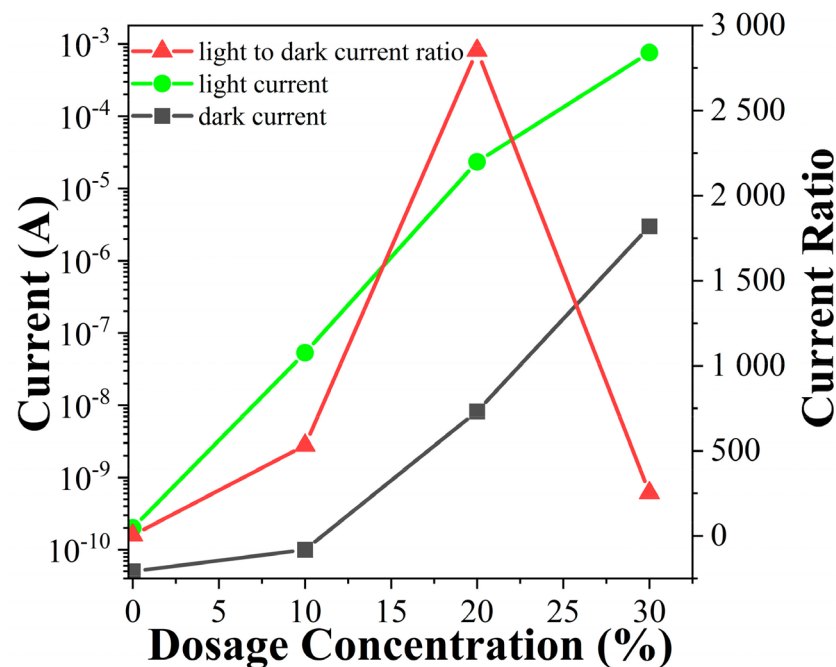


Figure 8. The I_l , I_d and I_l to I_d ratio values of devices with different In dosage concentration under 5 V bias.

Figure 9a shows the characteristic photoresponse spectra of the photodetector at 5 V bias. The responsivity can be expressed as [33]:

$$R_\lambda = (I_{ph}(\lambda) - I_d(\lambda)) / (P(\lambda)A) \quad (2)$$

where $I_{ph}(\lambda)$ and $I_d(\lambda)$ are photocurrent and current in dark, λ is wavelength of light, A and $P(\lambda)$ are effective area of the devices and luminous power, respectively. Under 260 nm UV light irradiation, the β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃ and In_{0.4}Ga_{1.6}O₃ devices achieve their maximum responsivity, at 0.00433 A/W, 1.86 A/W and 739.2 A/W, respectively. Under 270 nm UV light irradiation, the In_{0.6}Ga_{1.4}O₃ device reaches its maximum responsivity, at 20,579 A/W. Noise equivalent power (NEP) is defined as the minimum incident radiation power demanded to realize an SNR (signal to noise ratio) of 1 in a 1 Hz bandwidth. NEP can be expressed as [33]:

$$NEP = \sqrt{2qI_d}/R(\lambda) \quad (3)$$

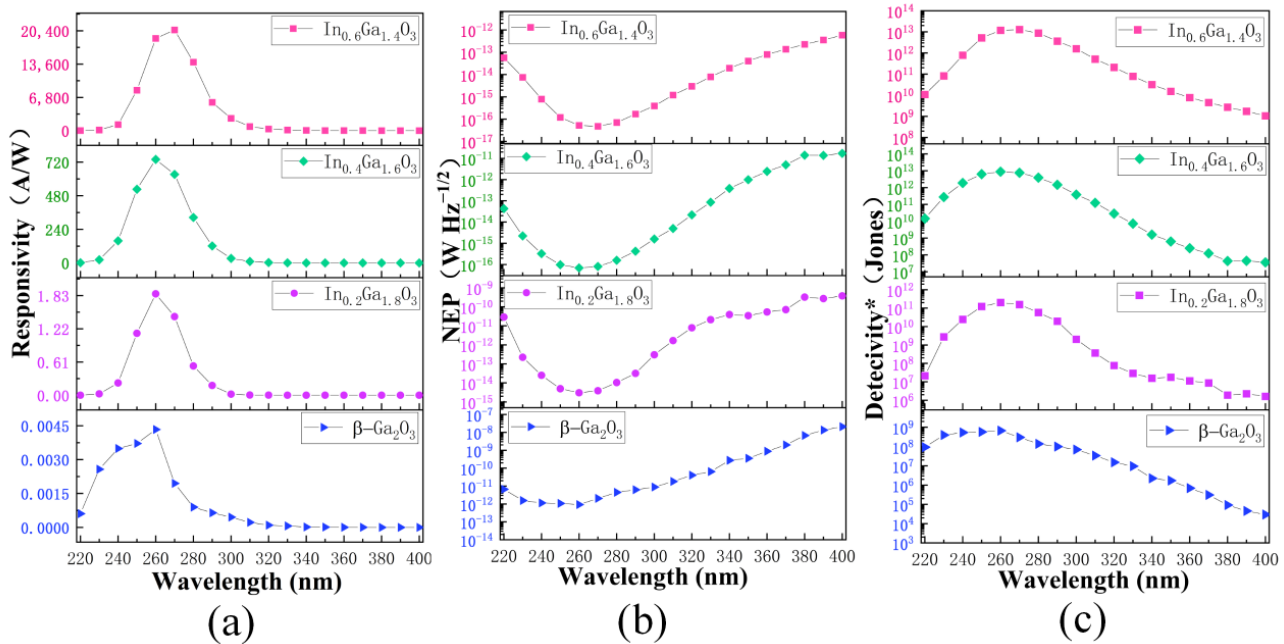


Figure 9. (a) Spectral responsivity (b) NEP (c) Spectral Detectivity* of the β -Ga₂O₃ and IGO devices at 5 V bias.

The main noise of MSM photodetectors is shot noise, so the mean-square noise current can be shot noise, as in Formula (3). NEP is relevant to a low noise signal and a large responsivity. In Figure 9b the NEP of β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃ and In_{0.4}Ga_{1.6}O₃ devices are $9.23 \times 10^{-13} \text{ W Hz}^{-1/2}$, $3.04 \times 10^{-15} \text{ W Hz}^{-1/2}$ and $6.93 \times 10^{-17} \text{ W Hz}^{-1/2}$ at 260 nm, and the NEP of In_{0.6}Ga_{1.4}O₃ device is $4.76 \times 10^{-17} \text{ W Hz}^{-1/2}$ at 270 nm. The smaller the NEP is, the better the device's performance. The detectivity* D^* represents the ability to detect weak signals from a noise environment, and is calculated by [33]:

$$D = \sqrt{A}/NEP \quad (4)$$

The D^* of β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃ and In_{0.4}Ga_{1.6}O₃ devices in Figure 9c are 6.68×10^8 Jones, 2.03×10^{11} Jones and 8.89×10^{12} Jones under 260 nm UV irradiation. The D^* of In_{0.6}Ga_{1.4}O₃ device is 1.29×10^{13} Jones under 270 nm UV irradiation.

For IGO photodetectors with a 5V bias applied, Figure 10 shows the rise time (τ_r) and decay time (τ_d) of the photocurrent. The response times (from 10% to 90% of the total value) for β -Ga₂O₃, In_{0.2}Ga_{1.8}O₃, In_{0.4}Ga_{1.6}O₃, and In_{0.6}Ga_{1.4}O₃ are 0.57 s, 3.58 s, 6.23 s, and 9.42 s, respectively. The recovery times (from 90% to 10% of the total value) are 0.07 s, 0.76 s, 4.18 s, and 18.63 s. The devices exhibit such time response characteristics mainly due to the conduction mechanism of oxygen vacancies. The undoped β -Ga₂O₃ thin film device has high crystalline quality and fewer oxygen vacancies, allowing electron-hole pairs to recombine rapidly. In the In-doped thin film devices, the presence of an additional electron layer of the In element compared to the Ga element results in a larger ionic radius,

which hinders the crystal growth process of $\beta\text{-Ga}_2\text{O}_3$, leading to an amorphous appearance and the generation of more oxygen vacancy defects [34,35]. Therefore, the response and recovery times of the devices increase significantly in thin film devices with different In doping concentrations. The higher the In doping concentration, the more oxygen vacancies are generated, and the longer the response and recovery times of the thin film devices are.

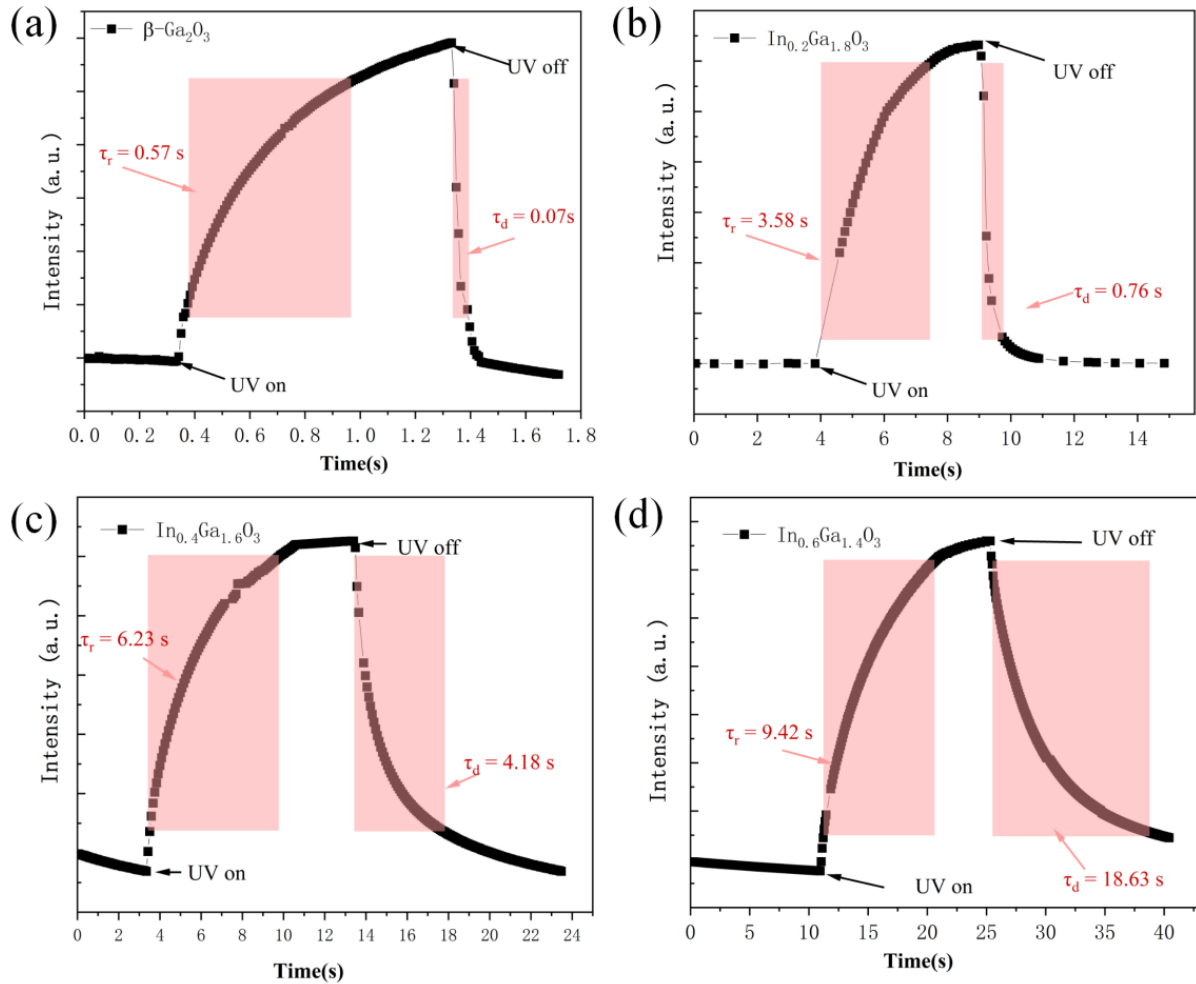


Figure 10. Time response characteristics of (a) $\beta\text{-Ga}_2\text{O}_3$; (b) $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$; (c) $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ and (d) $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$ devices.

Table 1 is a summary of the parameters of $\beta\text{-Ga}_2\text{O}_3$, $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$, $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, and $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$ MSM UV photodetectors.

Table 1. A summary of the parameters of $\beta\text{-Ga}_2\text{O}_3$, $\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$, $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, and $\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$ photodetectors.

	$\beta\text{-Ga}_2\text{O}_3$	$\text{In}_{0.2}\text{Ga}_{1.8}\text{O}_3$	$\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$	$\text{In}_{0.6}\text{Ga}_{1.4}\text{O}_3$
E_g (eV)	4.84	4.46	4.28	4.18
I_{light} (A)	2.0×10^{-10}	5.32×10^{-8}	2.34×10^{-5}	7.59×10^{-4}
I_{dark} (A)	5.1×10^{-11}	1.0×10^{-10}	8.2×10^{-9}	3.1×10^{-6}
I_{light} to I_{dark} ratio	4	532	2853	253
R (AW^{-1})	0.00433@260 nm	1.86@260 nm	739.2@260 nm	20,579@270 nm
NEP ($\text{W Hz}^{-1/2}$)	9.23×10^{-13} @260 nm	3.04×10^{-15} @260 nm	6.93×10^{-17} @260 nm	4.76×10^{-17} @270 nm
D^* (Jones)	6.68×10^8 @260 nm	2.03×10^{11} @260 nm	8.89×10^{12} @260 nm	1.29×10^{13} @270 nm
τ_r (s)	0.57	3.58	6.23	9.42
τ_d (s)	0.07	0.76	4.18	18.63

Figure 11 shows the energy band diagram to explain the overall mechanism. Now we discuss the main cases of the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ MSM structure. Before contact in Figure 11a,b, the work function of Au is 5.23 eV, and the work functions of $\beta\text{-Ga}_2\text{O}_3$ and $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ are 4.11 eV and 5.12 eV, respectively, which are both higher than those of Au, and thus forming Schottky contacts. The electron affinities of $\beta\text{-Ga}_2\text{O}_3$ and $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ are 4.11 eV and 4.9 eV, and the two bandgaps are 4.9 eV and 4.3 eV. Both $\beta\text{-Ga}_2\text{O}_3$ and In_2O_3 are n-type semiconductors, and the bandgap of In_2O_3 is 3.67 eV. With In doping into $\beta\text{-Ga}_2\text{O}_3$, the bandgap becomes small and the Fermi level changes from 4.11 eV to 5.12 eV [36,37]. Furthermore, with increasingly higher In doping concentrations, a larger number of oxygen vacancies capture holes, leading to an augmentation of electron transition density [38]. These interface oxygen vacancies contribute to the interface state, which leads to a decrease in the Schottky barrier. After contact with no bias in Figure 11c, there are bulk oxygen vacancies and interface oxygen vacancies in $\text{Au}/\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3/\text{Au}$, and the spikes form when both Fermi levels are unified. When a certain bias is applied on the MSM structure, there is always a reverse Schottky diode. Electrons will cross the barrier under bias, in the way of hot electron emission and tunneling, which forms the dark current, as is shown in Figure 11d. The Fermi level of Au in two reverse Schottky junctions will differ with electrons flowing under bias. Under UV light and a certain bias in Figure 11e, photogenerated electron-hole pairs will be separated by the built-in electrical field of Au and $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$, which contributes to the accumulation of carriers and thus improves the photocurrent [39]. As the In doping concentration increases from 10% to 20%, the amorphous state of the film increases. The significant increase in photocurrent is mainly attributed to the increase in the density of electron transitions within the film, while the enhancement of the light-to-dark current ratio is attributed to the pronounced increase in photocurrent [40,41]. With the In doping content rising from 20% to 30%, the film transforms from an amorphous state to the In_2O_3 phase. The significant increase in dark current of the device is mainly attributed to the increase in electron mobility (bulk oxygen vacancies by In doping) and the narrowing of the Schottky barrier (interface state by interface oxygen vacancies) due to the formation of the In_2O_3 phase [42]. The decrease in the light-to-dark current ratio of the device is attributed to the significant increase in dark current. Since the growth rates of photocurrent and dark current are inconsistent, the device with an In doping concentration of 20% has the highest photocurrent-to-dark current suppression ratio.

Table 2 is the performances compared with other GaO-alloy MSM UV photodetectors, which differ among various element compositions and proportions. The $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ MSM UV photodetector in this work has a relatively high light current, photo to dark current ratio and responsivity, which indicates the $\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$ material having a promising future in UV detection.

Table 2. Performances compared with different MSM ultraviolet photodetectors.

Material	V_{bias} (V)	I_1 (A)	I_1/I_d	Responsivity (A/W)	Ref.
InGaO	5	1.9×10^{-9}	82.6	6.9×10^{-5} @270 nm	[43]
$(\text{In}_x\text{Ga}_{1-x})_2\text{O}_3$	5	3.1×10^{-5}	1.3×10^3	27.7@255 nm	[25]
$(\text{Mg}_x\text{Ga}_{1-x})_2\text{O}_3$	5	1.4×10^{-5}	$\sim 10^5$	8.9@254 nm	[44]
$\text{Mg}_{0.18}\text{Zn}_{0.82}\text{O}$	13	1.2×10^{-6}	2	0.2@322 nm	[3]
$\text{In}_{0.9}\text{Ga}_{0.1}\text{O}$	10	5.0×10^{-6}	$\sim 10^5$	0.31@300 nm	[23]
$\text{Ga}_{2-x}\text{Sn}_x\text{O}_3$	50	8.7×10^{-7}	1.4×10^2	9.6×10^{-2} @254 nm	[45]
$\text{In}_{0.4}\text{Ga}_{1.6}\text{O}_3$	5	2.3×10^{-5}	2.8×10^3	739.2@260 nm	this work

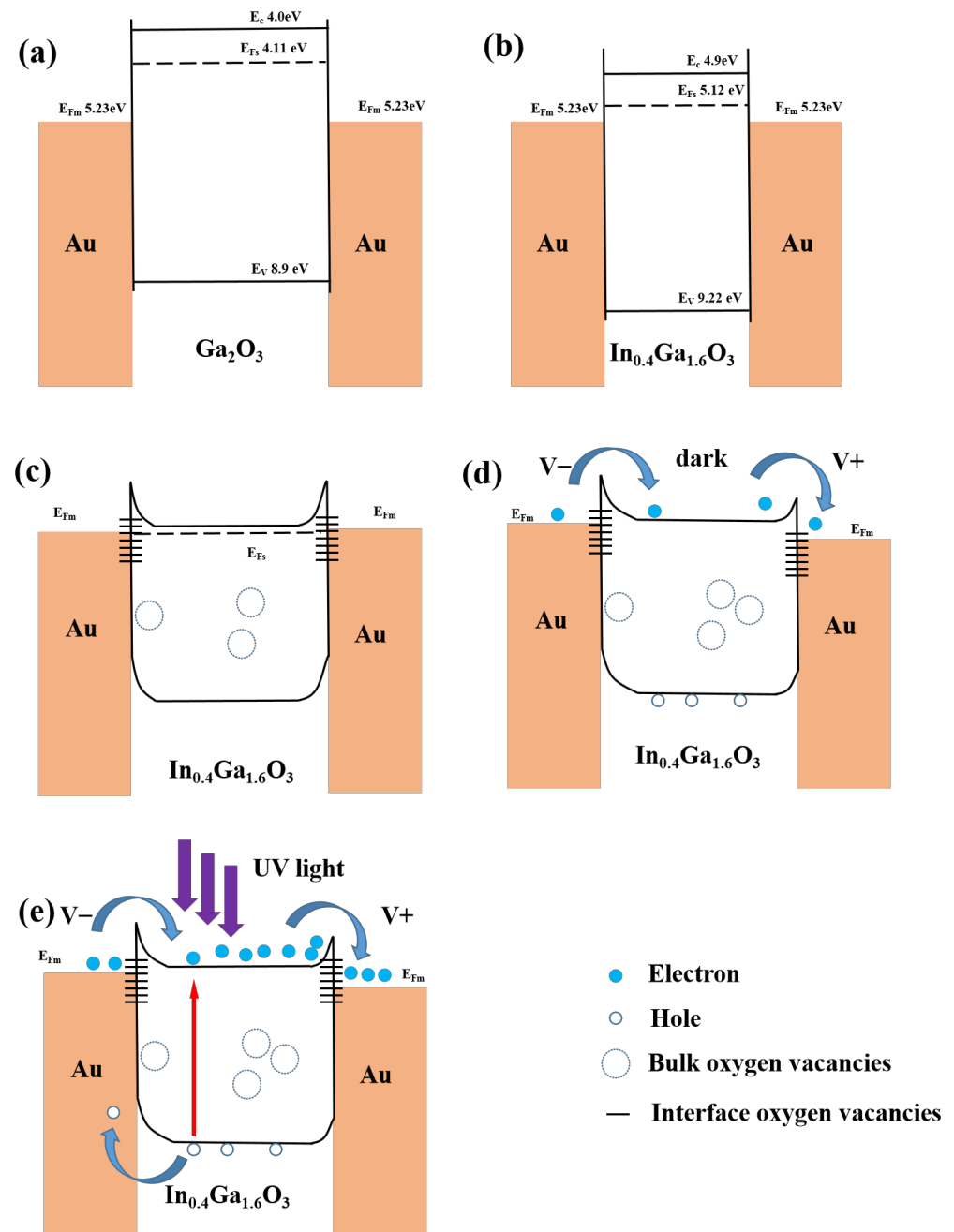


Figure 11. Schematic energy band diagrams of MSM structure. (a) Au/Ga₂O₃/Au before contact; (b) Au/In_{0.4}Ga_{1.6}O₃/Au before contact; (c) Au/In_{0.4}Ga_{1.6}O₃/Au after contact without bias; (d) Au/In_{0.4}Ga_{1.6}O₃/Au with bias in dark; (e) Au/In_{0.4}Ga_{1.6}O₃/Au with bias under UV light.

4. Conclusions

In summary, the effects of indium (In) content on IGO thin-film photodetectors were investigated. As the In doping concentration increased, the thin film transitioned from the β -Ga₂O₃ phase to amorphous phase, and then to the In₂O₃ phase. A device that exhibits high responsivity (739.2 A/W) and a high photocurrent-to-dark-current suppression ratio (2.8×10^3) was prepared by utilizing the In_{0.4}Ga_{1.6}O₃ thin film, which is in the intermediate phase transition process. Furthermore, the influence of Schottky barriers and oxygen vacancies on photocurrent and dark current was explored. Devices with superior performance and greater application potential can be developed by improving traditional single-layer devices through elemental doping.

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