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Structural, Magnetic and Optical Properties of Gd and Co Co-Doped YFeO₃ Nanopowders

Meng Wang ¹,*,[†] and Ting Wang ²,[†]

- ¹ School of Mechanical and Electrical Engineering, Shenzhen Polytechnic, Shenzhen 518055, China
- ² Shenzhen Key Laboratory of Advanced Materials, Department of Materials Science and Engineering, Shenzhen Graduate School, Harbin Institute of Technology, Shenzhen 518055, China
- * Correspondence: wangmeng04@szpt.edu.cn; Tel.: +86-755-26731137
- + These authors contributed equally to this work.

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Abstract: YFeO₃, YFe_{0.95}Co_{0.05}O₃, Y_{0.95}Gd_{0.05}FeO₃ and Y_{1-x}Gd_xFe_{0.95}Co_{0.05}O₃ (x = 0.0, 0.05, 0.10, 0.15 and 0.20) nanopowders were successfully fabricated via a low-temperature solid-state reaction technique. Results obtained using X-ray diffraction (XRD), scanning electron microscopy (SEM), and Raman spectra indicate that YFeO₃ nanopowders with Gd³⁺ and Co³⁺ ions co-doping at Y and Fe-sites were fabricated at 800 °C in sizes below 50 nm, and a distorted structure was obtained. Magnetic hysteresis loop analyses illustrate that ferromagnetic behavior of YFeO₃ nanopowders can be enhanced with the addition of Gd and Co. Whereas the maximum and remnant magnetization of the powders were found to be about 5.24 and 2.6 emu/g, respectively, the optical band gap was around 2.4 eV, proving that co-doped YFeO₃ nanopowders have a strong capability to absorb visible light. Because both magnetic and optical properties of these materials are greatly improved with the addition of Gd and Co, one can expect the scope of their potential application in the magnetic and optical fields to increase.

Keywords: multiferroics; low-temperature solid-state reaction; optical properties; magnetic properties; co-doping

1. Introduction

As one of the cutting edge multiferroic materials, $AFeO_3$ (A = La, Y and Sc) materials have been the focus of industry research because of their couple orderings of ferroelectricity and anti-ferromagnetism. As a result, they have great potential application in data storage, information exchange, and 5G mobile phone systems [1]. YFeO₃ is one of the most promising applications of rare earth AFeO₃ materials [2]. YFeO₃ has been reported to feature molecular ferroelectricity at low temperatures (10–40 K), good dielectric and magnetic properties [3–6], and is also becoming one of the most widely investigated multiferroic materials. Moreover, with a narrow optical gap (1.9–2.6 eV), its potential application as an optical material, especially decomposing organics should be considered for further study [7]. However, pristine YFeO₃ is not easy to prepare as evident by problems such as the introduction of secondary phases and the hopping of charge between Fe²⁺ and Fe³⁺ [8–10], while its low magnetism characteristics [11] are possibly its main shortcoming. Studies show that an effective and efficient way [11-13] to overcome these problems, is by means of doping. Some scientists substituted Y by divalent and trivalent ions [14,15] and Fe by trivalent and high-valence ions [16]. While the use of Gd doping is rarely reported, one study showed that with Gd doping on the Y site, significant enhancement of magnetization was achieved. In addition, use of other-element doping on the Y site has the potential for improvement of the material's optical properties [15]. For Fe-site ion doping, the main improvement focus was the reduction of leakage current and the enhancement of magnetic

properties. It is known [17] that Y(Fe,Cr)O₃ showed evidence of ferromagnetic property. Ma et al. [18] synthesized Mn doped YFeO₃ by using a standard solid-state reaction method. Their findings showed that the leakage current was reduced so that ferroelectric properties are realized at room temperature.

Co doping may enhance the magnetic properties of $RFeO_3$ material because it is magnetically active. While there is no report of co-doping of Co and Gd, the purpose of this paper is the study of co-doping of Co and Gd on YFeO₃ with a particular emphasis on the microstructural, optical, and magnetic properties of the doped YFeO₃ nanopowders.

2. Experimental Procedures

YFeO₃ (abbreviated as YFO), YFe_{0.95}Co_{0.05}O₃ (Co5), Y_{0.95}Gd_{0.05}FeO₃ (Gd5), Y_{0.95}Gd_{0.05}Fe_{0.95}Co_{0.05}O₃ (Gd5Co5), Y_{0.9}Gd_{0.1}Fe_{0.95}Co_{0.05}O₃ (Gd10Co5), Y_{0.85}Gd_{0.15}Fe_{0.95}Co_{0.05}O₃ (Gd15Co5), and Y_{0.8}Gd_{0.2}Fe_{0.95} $Co_{0.05}O_3$ (Gd20Co5) nanoparticles were synthesized using a low-temperature solid-state reaction technique as described in Ref. [19]. The raw reagents include $Fe(NO_3)_3 \cdot 9H_2O$, $Y(NO_3)_3 \cdot 6H_2O$, $Co(NO_3)_2 \cdot 6H_2O$, $Gd(NO_3)_3 \cdot 6H_2O$, and citric acid. Initially, $Fe(NO_3)_3 \cdot 9H_2O$, $Y(NO_3)_3 \cdot 6H_2O$ and citric acid were weighed stoichiometrically and ground in an agate mortar for half an hour, respectively. The respective powders were then mixed and ground again in an agate mortar for half an hour. A light brown viscous substance was formed during the grinding process, implying that the complex was formed. The viscous substance was heated at 120 °C for 2 h to remove free water and then a powdery composite was harvested, serving as the precursor powders. The powders were ground and subsequently calcined for 1 h in air at 800 °C to provide the nanoparticles. The crystal structures and morphology of the calcined particles were examined using XRD (D/max-RB, Rigaku, Tokyo, Japan) and SEM (S-4700, Hitachi, Tokyo, Japan). Raman spectra (InVia Reflex, Renishaw, London, UK spectrometer) were measured at room temperature with a range of 100 cm⁻¹ to 1500 cm⁻¹. Magnetic hysteresis loops were measured using PPMS (Physical Property Measurement System, DynaCool-9T, Quantum Design, Leatherhead, Surrey, UK). Optical absorbance properties of the experimental nanopowders were examined using a UV-Visible spectrophotometer (Shimadzu, UV-2600, Kyoto, Japan).

3. Results and Discussions

Figure 1a presents the XRD patterns of the tested ($0 \le x \le 0.2$) samples calcined at 800 °C. The pattern for the pristine YFeO₃ nanopowders suggests the presence of the obvious orthorhombic YFeO₃ pattern and that no minor impurity peaks were present. This outcome shows that the synthesis reaction for the orthorhombic structure was completed and the purity was high. The doped nanopowders indicate the same characteristic shape as that of the pristine $YFeO_3$ patterns with a slight shift in the peak position and phase. For Gd5, some peaks for the hexagonal YFeO₃ structure emerged. For Co5, peaks for Y_2O_3 appeared, showing that some minor impurity was introduced. This is a normal situation involving the case where instability of YFeO₃ exists and which is due to the radius difference. However, this effect totally disappears for the Co-Gd co-doped sample. Figure 1b shows the YFeO₃ peaks at $2\theta \sim 33^\circ$, shifting toward a higher 2θ angle with substitution of Co into YFeO₃. The figure also shows a shift towards a lower 2θ angle with further substitution of Gd into YFeO₃, and while keeping almost the same position of Gd and Co-doped samples due to Gd^{3+} (Gd³⁺ 0.938 Å) possessed a larger ionic radii compared with Y³⁺ (Y³⁺ 0.9 Å), and Co³⁺ (Co³⁺ 0.61 Å) possessed a smaller one compared with Fe³⁺ (Fe³⁺ 0.645 Å) [15,20]. Thus, it is evident that Gd³⁺ and Co³⁺ replaced the Y³⁺ and Fe³⁺ ions in YFeO₃, respectively. The XRD pattern shows the highest peak for the Co-doped particles and its 20 angle location is nearly identical to that of pure YFO. The stress introduced by the change in radius was reduced so that impurities disappeared.

Table 1 shows the lattice parameters with different Co and Gd concentrations. All these values increased with increase in Gd content and decreased with Co addition. This behavior demonstrates that the replacement of Y^{3+} with Gd^{3+} results in the increase of the lattice volume of YFeO₃ ($r_{Gd}^{3+} > r_Y^{3+}$), whereas the replacement of Fe³⁺ with Co³⁺ results in the decrease of YFeO₃ volume ($r_{Co}^{3+} < r_{Fe}^{3+}$).

The continuous evolution of lattice parameters with increasing Gd and Co concentrations also shows the presence of a successful substitution of Y^{3+} and Fe^{3+} ions by Gd^{3+} and Co^{3+} ions.



Figure 1. (a) XRD patterns of as-prepared nanoparticles calcined at 800 °C; (b) the magnified patterns of peaks at $2\theta \sim 32^{\circ}$.

Sample	a (Å)	b (Å)	c (Å)	Volume (Å ³)
Pure YFO	5.58079	5.27096	7.58866	223.23
Co5	5.57756	5.26958	7.58258	222.86
Gd5	5.58195	5.27209	7.5924	223.43
Gd5Co5	5.58018	5.27106	7.58943	223.23
Gd10Co5	5.58255	5.27279	7.59356	223.52
Gd15Co5	5.58303	5.27411	7.59364	223.60
Gd20Co5	5.58523	5.27791	7.59615	223.92

Table 1. Lattice parameters and volume of the samples.

SEM images of YFeO₃ nanopowders are shown in Figure 2. These demonstrate that the particle sizes of the test nanopowders were homogeneous with some minor agglomeration. The particle sizes of the test samples of YFO, Co5, Gd5, Co5Gd5, Co5Gd10, Co5Gd15 and Co5Gd20 were approximately 150, 120, 95, 80, 75, 45 and 40 nm, respectively. From the results obtained, it is apparent that the Co and Gd substitution decreased the grain size, while particle size decreased significantly with Co and Gd co-doping. As is well known, whenever the diffusion rate is low, rare earth ions can inhibit grain growth of RFeO₃ [11]. In addition, Gd³⁺ ions (radius of 0.983 Å) are larger than Y³⁺ ions (radius of 0.9 Å). The mismatch of ion sizes introduces defects in the lattice, leading to particle refinement. In addition, Co³⁺ can suppress hopping of Fe²⁺ and Fe³⁺, depriving the oxygen vacancy and subsequently preventing growth of the grains [21]. Some agglomeration of particles appears in Co and Gd co-doping samples, and may be attributed to the large surface area to volume ratio of the nanoparticles. The latter is to be expected whenever the low-temperature solid-state reaction method

is used. Similar effects of grain refinement induced by rare earth and magnetic element co-doping has been reported in other references [14,17,21,22].



Figure 2. SEM micrographs of as-prepared nanoparticles: (a) YFO; (b) Co5; (c) Gd5; (d) Gd5Co5; (e) Gd10Co5; (f) Gd15Co5; (g) Gd20Co5.

Figure 3 shows the Raman patterns obtained from the YFeO₃ powders. As is well known, for the case involving the Pnma structure of YFeO₃, only modes of A_{1g} , B_{1g} , B_{2g} , and B_{3g} are active. Furthermore, one can say with a high degree of certainty that the first and second A_{1g} modes are associated with Y–O bonds, and the A_{1g} and B_{g} modes are related to Fe-O bonds at higher

frequencies [23]. The 221 cm⁻¹ A_{1g} mode is represented as an FeO₆ octahedral structure [9,10,24]. The vibrating modes of Co5 were similar to those of pure YFO, while the intensity was reduced slightly and the peak at 610 cm⁻¹ increased as expected owing to the minor change of Fe–O bonds resulting from the doping of the Co element. For the case involving substitution of Gd³⁺ for Y³⁺, the parameters of lattice (grain size) increased, leading to a shift in the Raman bands to a lower wave-number. Raman bands were found to broaden with increase of Gd³⁺ and Co³⁺ (in YFeO₃) because of the disorder introduced by two different cations. The Gd–O and Y–O bonds, with their inherent different strengths, can influence their vibration frequencies. In general, one can state that these results are consistent with the XRD patterns. Thus, it can be concluded that Co and Gd co-doping has a substantial effect on the modification of the YFeO₃ structure.



Figure 3. (a) Raman spectra of as-prepared nanoparticles. (b) The magnified patterns of 221 cm⁻¹.

Figure 4 shows the magnetic properties of the nanopowders at room temperature. Values of maximum magnetization (M_m), remnant magnetization (M_r), and the coercive field (H_c) of the test samples are listed in Table 2. As is commonly known, YFeO₃ is antiferromagnetic and features weak magnetic properties. In the case of pristine YFeO₃, the hysteresis loop is typical weak magnetization in antiferromagnetic type. The magnetization parameters for the pristine YFeO₃ sample M_m , M_r , and H_c had magnitudes of 3.50 emu/g, 0.89 emu/g, and 161 Oe. For the case involving the doped samples, a large open region was seen at the center of the hysteresis loops suggesting ferromagnetic behavior. Even when exposed to a 60 kOe outer magnetic field, the loops were not saturated. For Co5, M_m , M_r were slightly improved and H_c was reduced. Their corresponding values were about 4.56 emu/g, 0.95 emu/g and 120 Oe, respectively. For the case of co-doping of Co and Gd, the net magnetization decreased slightly at first, and then began to increase, reaching a high value of 5.24 emu/g for the Co5Gd20 particles. These values are summarized in Table 2. The value of M_r reached a maximum ($M_r = 1.66 \text{ emu/g}$) for the Co5Gd5 nanoparticles. The reasons for observed improvement of the magnetization may be summarized as follows:

(1) The effect of nanoparticles. Uncompensated surface spins of Fe^{3+} ions are created when particle size is small. This situation leads to a strong magnetic enhancement [25,26].

(2) The distorted structure is affected by the doping effect. A FeO₆ octahedron structure consists of one Fe³⁺ ion and six O^{2-} ions, while each Fe³⁺ magnetic moment is not precisely parallel to the neighboring ones, forming a small angle [25,26], which causes weak ferromagnetism in the antiferromagnetic YFeO₃. Substituting an Fe³⁺ ion with a larger radius Co³⁺ ion will reduce the canted angle of the FeO₆ octahedra and release the distortion, and combined with a refined powder size, a higher magnetization is achieved. [27].

(3) The Gd³⁺ ion with a large magnetic moment ($\mu_{eff} = 8.0 \ \mu_B$) is magnetically active. When Y³⁺ is replaced by Gd³⁺, Y–O–Gd chains instead of Gd–O–Gd chains are formed, further improving the magnetization of the Co and Gd co-doped nanopowders [15].

In summary, the three reasons noted above serve to bring about the enhancement of magnetization. Others who conducted similar research made similar findings. Khalifa et al. [16] reported that M_m and M_r were about 0.8 emu/g and 0.1 emu/g for Ti-doped YFeO₃ nanoparticles prepared using an improved sol-gel technique. Shi et al. [17] synthesized YFe_{0.5}Cr_{0.5}O₃ nanoparticles by the sol-gel method, and showed that M_m and M_r were approximately 2.0 and 0.5 emu/g, respectively. As reported by Shi, the maximum M_m and M_r values for Y_{0.95}Ho_{0.05}Fe_{0.5}Cr_{0.5}O₃ nanoparticles prepared using a low-temperature citric acid assisted sol-gel technology were about 4.5 and 1.2 emu/g, respectively [17]. Yuan et al. [15] synthesized Y_{0.9}Gd_{0.1}FeO₃ using a solid-state reaction method, and obtained values of M_m , M_r , and H_c of 2.5 emu/g, 1.0 emu/g and 30,000 Oe, respectively. Still, it is worth noting that the property values in the present study are comparable with, or better than the results reported by others.



Figure 4. Cont.



Figure 4. Magnetic hysteresis loops of as-prepared nanoparticles: (**a**) YFO; (**b**) Co5; (**c**) Gd5; (**d**) Gd5Co5; (**e**) Gd10Co5; (**f**) Gd15Co5; (**g**) Gd20Co5.

Sample	M_m (emu/g)	M_r (emu/g)	H_c (kOe)	Photon Energy (eV)
Pure YFO	3.49	0.88	160	2.42
Co5	4.10	1.47	3700	2.34
Gd5	4.56	0.84	170	2.33
Gd5Co5	4.40	1.66	6800	2.32
Gd10Co5	4.68	1.76	2800	2.32
Gd15Co5	5.05	1.62	4800	2.29
Gd20Co5	5.24	1.60	4200	2.24

Table 2. List of the magnetic and optical parameters for YFeO₃ powders.

YFeO₃ possesses a narrow optical band gap (1.6–2.4 eV) and has been used in light-electric energy conversion applications. From the UV-visible absorption spectra (Figure 5), the optical energy band gap (E_g) of the seven test samples can be calculated using a Tauc function expressed in Equation (1) [28]:

$$(\alpha hv)^n = A(hv - E_g) \tag{1}$$

The E_g values shown in Figure 5 illustrate strong visible light absorption, indicating their promising decomposition application (See Table 2). YFO has the maximum E_g value (2.42 eV). It is decreased to 2.15 eV for Gd20Co5 nanoparticles. Clearly, the energy band gap becomes smaller with refined particle size of the YFeO₃ powders. Thus, the reduction of the energy gap with co-doping is ascribed to the smaller particle size and lattice distortion reduction generated by the addition of Co and Gd. Further, according to Reference [23], reduced particle size leads to a narrow energy gap. Nonuniform microstrains caused by lattice distortions can impact energy levels, affecting the energy band gap [28,29]. YFeO₃ possesses the band gap of 2p O, and 4d Y atoms. All these states are partially filled [28].

The Gd atom is with the 4f state which is also not fully filled orbits. Thus f-d hybridization of the 4d Y and 3d Fe atoms causes light absorption. In addition, the same d states of Fe and Co overlap with the 2p states O atom, causing a narrow energy gap [15,30]. Moreover, Gd³⁺ and Co³⁺ ions partially substituting Y³⁺ and Fe³⁺ ions refine the as-synthesized nanoparticles size and increase chemical pressure, thereby resulting in a smaller value of E_g [30]. Zhang et al. [28] obtained E_g values of 1.94, 2.43 and 2.30 eV, respectively, for the hexagonal, orthorhombic, and YFeO₃ containing a mixture of the two phases. Shen et al. [31] reported that the E_g value obtained from a first-principles calculation for YFeO₃ ceramics was 2.58 eV. Wu et al. [23] prepared YFeO₃ in which, hexagonal and orthorhombic phases co-existed; this material had a band gap of 2.41 eV. For YFeO₃ ceramics prepared by the conventional solid-state method, the energy band gap was found to be 2.58 eV [7]. Liu et al. [32] measured the optical properties of hexagonal-YFeO₃/ α -Fe₂O₃ heterojunction composite nanowire and obtained E_g values of approximately 2.15 eV. Once again, it is worth noting that our E_g values compare favorably with these results. Therefore, one can reasonably conclude that the YFeO₃ nanopowders used in this study can be used in the decomposition of organic compounds.



Figure 5. Cont.



Figure 5. $(\alpha h\nu)^2$ as a function of photon energy for as-prepared nanoparticles: (**a**) YFO; (**b**) Co5; (**c**) Gd5; (**d**) Gd5Co5; (**e**) Gd10Co5; (**f**) Gd15Co5; (**g**) Gd20Co5.

4. Conclusions

Co and Gd co-doped YFeO₃ nanopowders were fabricated using a low-temperature solid-state reaction method. Data obtained using XRD and Raman analyses show that, with Gd and Co substitution, YFeO₃ nanoparticles exhibit a distortion refined structure. The particle size for pristine YFeO₃ nanoparticles is about 100-200 nm, and those for Co5, Gd5, Co5Gd5, Co5Gd10, Co5Gd15 and Co5Gd20 are approximately 120, 95, 80, 75, 45 and 40 nm, respectively. The maximum magnetization and remnant magnetization for the co-doped YFeO₃ powders are about 5.49 and 2.20 emu/g, respectively, when exposed to a magnetic field of 60 kOe. The energy band gap of YFeO₃ nanopowders was reduced from 2.41 to 2.23 using the co-doping method, thereby indicating their potential in decomposition applications. Because the co-doping method used in the study was found to be easy to control for fabricating the YFeO₃ nanopowder samples, it is proposed that the method be adopted for use in the applied magnetic and optical fields.

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References

- Dho, J.; Blamire, M. Competing functionality in multiferroic YMnO₃. *Appl. Phys. Lett.* 2005, *87*, 2504. [CrossRef]
- 2. Wang, L.J.; Feng, S.M.; Zhu, J.L.; Yu, R.C.; Jin, C.Q.; Yu, W.; Wang, X.H.; Li, L.T. Ferroelectricity of multiferroic hexagonal TmMnO₃ ceramics synthesized under high pressure. *Appl. Phys. Lett.* **2007**, *91*, 172502. [CrossRef]
- Cherry, M.; Islam, M.S.; Catlow, C.R.A. Oxygen ion migration in perovskite-type oxides. *J. Solid State Chem.* 1995, 118, 125–132. [CrossRef]
- 4. Minh, N.Q. Ceramic fuel cells. J. Am. Ceram. Soc. 1993, 76, 563-588. [CrossRef]
- Sharon, M.; Prasad, B.M. Preparation and photocharacterization of YFeO₃ semiconductor. *Electrochim. Acta* 1985, 30, 331–334. [CrossRef]
- Traversa, E.; Nunziante, P.; Sangaletti, L.; Allieri, B.; Depero, L.E.; Aono, H.; Sadaoka, Y. Synthesis and structural characterization of trimetallic perovskite-type rare-earth orthoferrites, La_xSm_{1-x}FeO₃. *J. Am. Ceram. Soc.* 2000, *83*, 1087–1092. [CrossRef]
- Butler, M.A.; Ginley, D.S.; Eibschutz, M. Photoelectrolysis with YFeO₃ electrodes. J. Appl. Phys. 1977, 48, 3070–3072. [CrossRef]
- Mathur, S.; Veith, M.; Rapalaviciute, R.; Shen, H.; Goya, G.F.; Filho, W.L.M.; Berquo, T.S. Molecule derived synthesis of nanocrystalline YFeO₃ and investigations on its weak ferromagnetic behavior. *Chem. Mater.* 2004, *16*, 1906–1913. [CrossRef]
- 9. Schmool, D.S.; Keller, N.; Guyot, M.; Krishnan, R.; Tessier, M. Magnetic and magneto-optic properties of orthoferrite thin films grown by pulsed-laser deposition. *J. Appl. Phys.* **1999**, *86*, 5712–5717. [CrossRef]
- Shang, M.; Wang, C.; Chen, Y.; Sun, F.; Yuan, H. The multiferroic epitaxial thin film YFeO₃. *Mater. Lett.* 2016, 175, 23–26. [CrossRef]
- 11. Akbashev, A.; Semisalova, A.; Perov, N.; Kaul, A. Weak ferromagnetism in hexagonal orthoferrites RFeO₃ (R = Lu, Er-Tb). *Appl. Phys. Lett.* **2011**, *99*, 122502. [CrossRef]
- 12. Hemberger, J.; Schrettle, F.; Pimenov, A.; Lunkenheimer, P.; Ivanov, V.Y.; Mukhin, A.A.; Balbashov, A.M.; Loidl, A. Multiferroic phases of Eu_{1-x}Y_xMnO₃. *Phys. Rev. B Condens. Matter* **2007**, *75*, 5118. [CrossRef]
- Zhang, W.; Fang, C.; Yin, W.; Zeng, Y. One-step synthesis of yttrium orthoferrite nanocrystals via sol-gel auto-combustion and their structural and magnetic characteristics. *Mater. Chem. Phys.* 2013, 137, 877–883. [CrossRef]
- Van Tac, D.; Mittova, V.; Mittova, I.Y. Influence of lanthanum content and annealing temperature on the size and magnetic properties of sol–gel derived Y_{1-x}La_xFeO₃ nanocrystals. *Inorg. Mater.* 2011, 47, 521–526. [CrossRef]
- 15. Yuan, X.; Sun, Y.; Xu, M. Effect of Gd substitution on the structure and magnetic properties of YFeO₃ ceramics. *J. Solid State Chem.* **2012**, *196*, 362–366. [CrossRef]
- Khalifa, N.O.; Widatallah, H.M.; Gismelseed, A.M. Magnetic and Mössbauer studies of pure and Ti-doped YFeO₃ nanocrystalline particles prepared by mechanical milling and subsequent sintering. *Hyperfine Interact.* 2016, 237, 1–11. [CrossRef]
- Shi, L.R.; Xia, Z.C.; Huang, S.; Xiao, G.L.; Jin, Z.; Wei, M.; Chen, B.R.; Shang, C.; Cheng, H.; Ouyang, Z.W. The magnetization reversal and high temperature dielectric response in Y_{1-x}H_xFe_{0.5}Cr_{0.5}O₃. *Ceram. Int.* 2016, 42, 10808–10812. [CrossRef]
- 18. Ma, Y.; Wu, Y.J.; Lin, Y.Q.; Chen, X.M. Microstructures and multiferroic properties of YFe_{1-x}Mn_xO₃ ceramics prepared by spark plasma sintering. *J. Mater. Sci. Mater. Electron.* **2010**, *21*, 838–843. [CrossRef]
- 19. Wang, M.; Wang, T.; Song, S.; Tan, M. Structure-controllable synthesis of multiferroic YFeo₃ nanopowders and their optical and magnetic properties. *Materials* **2017**, *10*, 626. [CrossRef]
- 20. Shannon, R.D. Revised effective ionic radii and systematic studies of interatomic distances in halides and chalcogenides. *Acta Crystallogr. A* **1976**, *32*, 751–767. [CrossRef]
- 21. Chakrabarti, K.; Das, K.; Sarkar, B.; Ghosh, S.; De, S.K.; Sinha, G.; Lahtinen, J. Enhanced magnetic and dielectric properties of Eu and Co co-doped BiFeO₃ nanoparticles. *Appl. Phys. Lett.* **2012**, *101*, 759. [CrossRef]
- 22. Yang, L.P.; Zhang, A.M.; Dai, Y.; Wu, X.S. The effect of Dy–Fe co-doping on the structural and magnetic properties of h-YMnO₃. *J. Mater. Sci. Mater. Electron.* **2017**, *28*, 8872–8877. [CrossRef]
- 23. Wu, L.; Yu, J.C.; Zhang, L.; Wang, X.; Li, S. Selective self-propagating combustion synthesis of hexagonal and orthorhombic nanocrystalline yttrium iron oxide. *J. Solid State Chem.* **2004**, 177, 3666–3674. [CrossRef]

- 24. Ma, Y.; Chen, X.M.; Lin, Y.Q. Relaxorlike dielectric behavior and weak ferromagnetism in YFeO₃ ceramics. *J. Appl. Phys.* **2008**, *103*, 124111. [CrossRef]
- 25. Moriya, T. New mechanism of anisotropic superexchange interaction. *Phys. Rev. Lett.* **1960**, *4*, 228–230. [CrossRef]
- 26. Dzyaloshinsky, I. A thermodynamic theory of weak ferromagnetism of antiferromagnetics. *J. Phys. Chem. Solids* **1958**, *4*, 241–255. [CrossRef]
- 27. Arlt, G.; Hennings, D.; With, G.D. Dielectric properties of fine-grained barium titanate ceramics. *J. Appl. Phys.* **1985**, *58*, 1619–1625. [CrossRef]
- 28. Zhang, Y.; Yang, J.; Xu, J.; Gao, Q.; Hong, Z. Controllable synthesis of hexagonal and orthorhombic YFeO₃ and their visible-light photocatalytic activities. *Mater. Lett.* **2012**, *81*, 1–4. [CrossRef]
- 29. Mocherla, P.S.V.; Karthik, C.; Ubic, R.; Rao, M.S.R.; Sudakar, C. Tunable bandgap in BiFeO₃ nanoparticles: The role of microstrain and oxygen defects. *Appl. Phys. Lett.* **2013**, *103*, 2463–2485. [CrossRef]
- Zhang, Q.; Zhu, X.; Xu, Y.; Gao, H.; Xiao, Y.; Liang, D.; Zhu, J.; Zhu, J.; Xiao, D. Effect of La³⁺ substitution on the phase transitions, microstructure and electrical properties of Bi_{1-x}La_xFeO₃ ceramics. *J. Alloys Compd.* 2013, 546, 57–62. [CrossRef]
- Shen, T.; Hu, C.; Yang, W.L.; Liu, H.C.; Wei, X.L. Theoretical investigation of magnetic, electronic and optical properties of orthorhombic YFeO₃: A first-principle study. *Mater. Sci. Semicond. Process.* 2015, 34, 114–120. [CrossRef]
- Liu, J.; He, F.; Chen, L.; Qin, X.; Zhao, N.; Huang, Y.; Peng, Y. Novel hexagonal-YFeO₃/α-Fe₂O₃ heterojunction composite nanowires with enhanced visible light photocatalytic activity. *Mater. Lett.* 2016, 165, 263–266. [CrossRef]



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