



Article Bottom-Up (Cu, Ag, Au)/Al₂O₃/Bi₂Te₃ Assembled Thermoelectric Heterostructures

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Abstract: The interface affects the transmission behavior of electrons and phonons, which in turn determines the performance of thermoelectric materials. In this paper, metals (Cu, Ag, Au)/Al₂O₃/ Bi_2Te_3 heterostructures have been fabricated from bottom to up to optimize the thermoelectric power factor. The introducing metals can be alloyed with Bi_2Te_3 or form interstitials or dopants to adjust the carrier concentration and mobility. In addition, the metal-semiconductor interface as well as the metal-insulator-semiconductor interface constructed by the introduced metal and Al_2O_3 would further participate in the regulation of the carrier transport process. By adjusting the metal and oxide layer, it is possible to realize the simultaneous optimization of electric conductivity and Seebeck coefficient. This work will enable the optimal and novel design of heterostructures for thermoelectric materials with further improved performance.

Keywords: thermoelectric; Bi2Te3; heterostructure; interface; metal; oxide

1. Introduction

Thermoelectric technology can convert heat to electricity freely. The thermoelectric (TE) efficiency is normally evaluated by the dimensionless figure of merit $ZT = S^2 \sigma T/\kappa$ ($\kappa = \kappa_e + \kappa_L$), where S, σ , κ , κ_e , κ_L , and T are the Seebeck coefficient, electrical conductivity, thermal conductivity, electronic thermal conductivity, lattice thermal conductivity, and absolute temperature, respectively [1]. One strategy to improve the thermoelectric performance is fabricating nanostructures such as pores [2,3] or nanograins [4,5], which could strongly scatter phonons and thereby reducing the lattice thermal conductivity. However, little room remains for further reducing κ . On the one hand, the contribution of electrons to thermal conductivity cannot be lower than the amorphous state. For Bi₂Te₃ system with excellent thermoelectric properties near room temperature, the lattice thermal conductivity has been approached to lower as ~0.3 W·m⁻¹·K⁻¹, which is the limit value predicted by Cahill model [6,7]. Further improvements in the ZT could be relied on decoupling S and σ to improve the power factor ($PF = S^2 \sigma$).

Compared with its bulk counterparts, nanostructured Bi_2Te_3 film provides promising possibilities for enhanced TE properties and potential applications in micro/nanoelectromechanical systems (MEMS/NEMS) TE device [8]. The thermoelectric performance of Bi_2Te_3 can be optimized by adjusting the (00*l*) orientation [9,10], pores [9,11], nanosheet boundary [12,13], and the intrinsic defects [14,15]. In addition, constructing a heterogeneous interface is also one of the effective methods to optimize the thermoelectric performance. The heterogeneous interface can filter low-energy carriers that contribute negative to the Seebeck coefficient through the interface barrier and thus the thermoelectric power factor will be enhanced [16]. Meanwhile, the interface could also scatter phonons to reduce



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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). the lattice thermal conductivity, and finally the *ZT* value of the material will be further improved [1]. Experimentally, metal/semi-metal or carbon had been introduced in Bi₂Te₃ to build heterostructures (such as Bi₂Te₃/Ag [12,17], Bi₂Te₃/Cu [12,18] Bi₂Te₃/Cd [19], Bi₂Te₃/Te [20], Bi₂Te₃/carbon nanotube [10,21]) especially by solution or sintering method to further enhanced the TE properties. Oxide layer was also introduced into the Ag/Sb₂Te₃ system by atomic layer deposition, and the filtering effect of the corresponding interface on low-energy carriers was discussed [6]. In this paper, group IB metals (M = Cu, Ag, Au) and Al₂O₃ interfaces were systematically introduced into Bi₂Te₃ film to modulate the carrier behavior and thus enhance the thermoelectric power factor. Incidentally, the introduced interface can also scatter phonons to reduce the thermal conductivity of the lattice, which is beneficial to improve the *ZT* value.

2. Materials and Methods

Group IB metal M (M = Cu, Ag, Au) with a thickness of $6 \sim 10$ nm was pre-deposited on SiO₂/Si substrates by an e-beam evaporation system. Using Al₂O₃ targets, Al₂O₃ layer with a thickness of ~3 nm was first selectively sputtered on the pre-deposited metal. Then, Bi₂Te₃ films with a thickness of about 300 nm were sputtered. The base vacuum of the sputtering chamber was $5E^{-3}$ Pa and the working pressure was 0.26 Pa. The substrate was rotated at a speed of 20 rpm to ensure the uniformity of the sputtered film. The Al₂O₃ sputtered at 100 W and Bi₂Te₃ sputtered at 50 W were both in a radio frequency (RF) mode. The purity of all source materials is 99.99%. The sample preparation is illustrated in Figure 1a. For the convenience of description in the following part, pure Bi₂Te₃ is marked as B. Cu/Bi₂Te₃, Ag/Bi₂Te₃, and Au/Bi₂Te₃, and Au/Al₂O₃/Bi₂Te₃ are marked as CAB, SAB, and GAB, respectively.



Figure 1. (a) Film preparation process. SEM image of (b) Cu, (c) Ag and (d) Au surface. AFM image of (e) Ag and (f) Ag/ Al₂O₃. (g) AFM step profiler of Ag and Ag/Al₂O₃.

The surface morphology, cross section, composition, and thickness of films were characterized by atomic force microscope (AFM, Multimode Nanoscope IIIa), scanning electron microscopy with energy dispersive spectrometer (SEM-EDS, Zeiss Ultra Plus) and surface profiler (KLA-Tencor P7). The structures of the films were characterized by X-ray diffraction (XRD, D8 Advance) with a Cu K α radiation ($\lambda = 1.5406$ Å) in conventional θ –2 θ mode from 10–80° at a sweep rate of 4°/min. Electrical transport of film was obtained by Hall measurement (model: MMR) with a Van der Pauw pattern and averaged by three measurements. The contact effect and other additional effect including heating effect of current and asymmetric effect of probing electrodes can be exclude in this current

reversal method with four-probe configuration. The Seebeck coefficient was evaluated by a home-built system as reported in detail in our previous study [22].

3. Results and Discussion

3.1. Heterostructure Characterization

The pre-deposited metal presents as discontinuous particles as shown in Figure 1b–d. The height of Ag is about 10 nm, while the height of Cu and Au is about 6–7 nm as shown in Figures 1 and S1. After sputtering Al_2O_3 , the total height of the metal particles increased by ~3 nm. Interestingly, Al_2O_3 appeared as smaller particles on the surface of the pre-deposited metal particles, resulting in increased particle fluctuations, as shown in Figure 1e–g.

The Te/Bi atomic ratio of all Bi₂Te₃ films that sputtered with the same target and sputtering power was about ~1.5. The surface and cross section of Bi₂Te₃ and M/Al₂O₃/Bi₂Te₃ are given in Figure S2. Take Ag/Al₂O₃/Bi₂Te₃ as an example, the surface of each layer is given in Figure 2a–c. Regardless of whether metal/oxide was previously deposited, there was no obvious difference in the surface morphology of the Bi₂Te₃ film, the surface of which is selected as shown in Figure 2c. The bottom-up Ag/Al₂O₃/Bi₂Te₃ structure could be clearly seen from the cross section as shown in Figure 2d. The thickness of the Bi₂Te₃ film was about 300 nm, and the XRD pattern of Bi₂Te₃ indicated that the obtained Bi₂Te₃ is with good crystallinity as shown in Figure 2e.



Figure 2. SEM image of (**a**) Ag/Al₂O₃, (**b**) Ag/Al₂O₃/Bi₂Te₃, and (**c**) Bi₂Te₃. (**d**) Cross section of Ag/Al₂O₃/Bi₂Te₃. (**e**) XRD pattern of Bi₂Te₃.

3.2. Carrier Transport Characteristics

All Bi₂Te₃-based films exhibit n-type conductivity, and the majority of carriers in the films are electrons as characterized by Hall measurements and shown in Figure 3a–c. Compared with the pure Bi₂Te₃ film, the introduction of M (M = Cu, Ag, Au) in the film leads to a decrease in total electron concentration of the film as shown in Figure 3a.

The role of introducing group IB metals in $B_{i2}Te_3$ can be summarized as follows: (i) In one way, one M atom should produce one electron compensating a hole. This process can be achieved by forming [MBiTe₂ +Bi₃Te₄⁻]⁻ through the alloying of Bi₂Te₃ and M as shown in Figure 3g [23]. Generally, the diffusion, solid solution and alloying process of group IB metal in Bi₂Te₃ can be modulated by heat treatment [23–26]. Because the films in this study have not been annealed, the above effects, especially the alloying process (generate new phases), are weak. (ii) Another way is that the introduction of M atom could form interstitial M_i⁺ compensated by electrons. While the electronegativity (Pauling scale) of Cu, Ag, and Au atoms increases in turn (from 1.90, 1.93 to 2.54), the ability to form M_i⁺ decreases in turn [14]. Generally, since the space in the lattice is limited, atoms or cations with a smaller radius are more likely to form interstitial phases. The radius of Cu, Ag, and Au are 1.28 Å, 1.44 Å, and 1.44 Å, respectively. The radius of Cu⁺, Ag⁺, and Au⁺ are 0.98 Å, 1.10 Å, and 1.10 Å, respectively [27]. In terms of the ability to form interstitial, the electron concentration of the Cu-introduced film should be less than that in the Ag-introduced film, which is confirmed in Figure 3a. (iii) Since M (M = Cu, Ag, and Au) have fewer electrons in their outermost orbitals than Bi, doping them at the Sb-sites greatly increases the hole concentration and consequently reduces the electron concentration [19].



Figure 3. Electrical properties of Bi₂Te₃-based films at 300 K. (a) Carrier concentration (*n*). (b) Carrier mobility (μ). (c) Electrical conductivity (σ). Equilibrium band gap alignment of (d) Ag, Al₂O₃ and Bi₂Te₃, (e) Cu, Al₂O₃, and Bi₂Te₃, (f) Au, Al₂O₃, and Bi₂Te₃. Behavior of electrons at the interface of (g) Metal/Bi₂Te₃ and (h) Metal/Al₂O₃/Bi₂Te₃.

The coexistence of the above alloying, interstitial and doping process together determines the reduction of the electron concentration in the Bi_2Te_3 film after the introduction of metal. Eventually, the electron concentration of the Bi_2Te_3 , Au/Bi_2Te_3 , Ag/Bi_2Te_3 , and Cu/Bi_2Te_3 film decreases sequentially. After the Al_2O_3 layer was deposited to build a metal-insulator-semiconductor (MIS) interface, the electrons in the metal were blocked from entering Bi_2Te_3 and some carriers may tunnel through the thin oxide layer from the metal into Bi_2Te_3 as shown in Figure 3g,h. Because there were unavoidable narrow holes in such thin Al_2O_3 layer deposited by magnetron sputtering, there would still be some interfaces where the metal was in direct contact with Bi_2Te_3 . The above-mentioned combined effect made the carrier concentration of the film further decrease after the Al_2O_3 layer was introduced.

As the electron concentration decreased after metal or oxide was introduced, the scattering between electrons was weakened and the mobility of the film increased. On the other hand, the introduction of a metal-semiconductor (MS) interface or an MIS interface would lead to the reduction of the mobility which was closely related to the size and distribution of the metal particles and the surface Al_2O_3 that covers them. The interfaces and the corresponding electronic transport processes are illustrated in Figure 3d–f. The above combined effect makes the total electrical mobility and conductivity of the films other than Ag/Bi_2Te_3 increase while compared to the pure Bi_2Te_3 , which may be related to the Ohmic contact of the Ag-Bi₂Te₃ interface and Ag particle with relatively high height (thickness). In addition to filtering out low-energy electrons based on the interface barrier, an additional oxide layer was deposited to control the tunneling electron concentration. Through the above dual effects, the concentration and average energy of electrons can be adjusted.

3.3. Thermoelectric Properties

The Seebeck coefficient is closely related to the carrier concentration and its energy distribution, and the contribution of low-energy carriers to the Seebeck coefficient is negative [28]. The measured Seebeck coefficient of Bi_2Te_3 -based film is presented in Figure 4a. It can be seen from Figure 4a that the Seebeck coefficients of the film with Cu and Au are lower than the pure Bi_2Te_3 , which may be related to the relatively high interface barrier that filters out some high-energy electrons. The deposited Al_2O_3 layer allowed

tunneling electrons to pass through, which would reduce the concentration of high-energy electrons in the film, and ultimately reduce the Seebeck coefficient.



Figure 4. Thermoelectric properties of Bi₂Te₃-based films at 300 K. (a) Seebeck coefficient (*S*). (b) Seebeck coefficient as a function of carrier concentration (Pisarenko plot). Dashed lines represent theoretical fits for $m^* = 0.3 m_0$, 2.5 m_0 , 3 m_0 , 4 m_0 and 7 m_0 based on Equations (1)–(3). m_0 is mass of free electron. (c) Power factor ($S^2\sigma$).

For a material with a single parabolic band, the Seebeck coefficient, DOS effective mass of carriers (m^*) and Lorenz number (L) can be written as [22,29]

$$S = \frac{K_B}{e} \left(\frac{(5/2 + \lambda)F_{3/2 + \lambda}(\eta)}{(3/2 + \lambda)F_{1/2 + \lambda}(\eta)} - \eta \right)$$
(1)

$$m^* = \frac{h^2}{2K_b T} \left(\frac{n}{4\pi F_{1/2}(\eta)}\right)^{2/3}$$
(2)

$$F_i(\eta) = \int_0^\infty \frac{x^i dx}{1 + \exp(x - \eta)} \tag{3}$$

$$L = \left(\frac{k_B}{e}\right)^2 \left\{ \frac{(\lambda + 7/2)F_{\lambda + 5/2}(\eta)}{(\lambda + 7/2)F_{\lambda + 1/2}(\eta)} - \frac{(\lambda + 5/2)F_{\lambda + 3/2}(\eta)^2}{(\lambda + 3/2)F_{\lambda + 1/2}(\eta)^2} \right\}$$
(4)

$$\kappa_{\rm e} = LT\sigma \tag{5}$$

where *e*, *m*^{*}, *h*, *k*_B, *L*, and *T* are the electronic charge, effective mass, Planck's constant, Boltzmann constant, Lorenz number, and absolute temperature, respectively. η is the reduced Fermi energy (E_F/k_BT), F_i is the fermi integral of the order of x. Acoustic phonon scattering ($\gamma = -1/2$) has been assumed as the main carrier scattering mechanism at room temperature. The corresponding calculation results at 300 K are listed in Table 1. In order to compare the effect of interface on the Seebeck coefficient of Bi₂Te₃ at room temperature, Pisarenko plot is applied as shown in Figure 4b.

Table 1. Thermoelectric parameters of Bi_2Te_3 -based films at 300 K.

Sample	В	СВ	CAB	SB	SAB	GB	GAB
m^*/m_0	6.38	2.56	0.26	3.80	0.30	3.92	3.12
$L (V^2 \cdot K^{-2})$	$2.04 imes10^{-8}$	$2.08 imes10^{-8}$	$2.12 imes 10^{-8}$	$2.01 imes 10^{-8}$	$2.05 imes10^{-8}$	$2.09 imes10^{-8}$	$2.14 imes10^{-8}$
$\kappa_{\rm e} (W \cdot m^{-1} \cdot K^{-1})$	0.37	0.40	1.00	0.18	0.56	0.47	0.72

After the metal was introduced into Bi_2Te_3 film, the m^* of the film was reduced, and it was further reduced after the oxide layer was introduced as shown in Table 1. However, when compared with Cu-introduced and Ag-introduced films, the effective mass change in the Au-introduced film was small, which was closely related to the carrier behavior as analyzed before. The power factor of Bi_2Te_3 -based film at 300 K is given in Figure 4c. Compared with pure Bi_2Te_3 , the power factor of the film introduced with metal and oxide

is increased, which benefits from the effect of metals and the MIS interface. As a result, the power factor of the Cu/Al₂O₃/Bi₂Te₃ film is 6.5 μ W·cm⁻¹·K⁻², which is 1.75 times that of a single Bi₂Te₃ film. In addition, the corresponding electronic thermal conductivity calculated according to Equation (5) is given in Table 1.

Although the introduction of Cu/Al_2O_3 leads to an increase in the electronic thermal conductivity of the film, the smaller Cu particles result in a higher interface density (Figure 1b), which could scatter phonons more strongly and make the lattice thermal conductivity lower [30]. In addition to scattering phonons at the interface to reduce the thermal conductivity, it is particularly important to use the interface to optimize the thermoelectric power factor to improve the overall performance of the film.

4. Conclusions

In summary, metal (Cu, Ag, Au)/Al₂O₃/Bi₂Te₃ heterostructures have been fabricated to improve the thermoelectric properties. By constructing the metal-semiconductor and metal-insulator-semiconductor interface, using the filter carrier effect or controlling the tunneling electrons, it is possible to coordinately control the carrier concentration and mobility, and finally optimize the thermoelectric power factor. The power factor of the Cu/Al₂O₃/Bi₂Te₃ heterostructure film is 6.5 μ W·cm⁻¹·K⁻², which is 1.75 times that of a single Bi₂Te₃ film. Future work will focus on optimizing the metal particle size, oxide layer thickness, and annealing process to further optimize the thermoelectric performance. Related works are still ongoing.

Supplementary Materials: The following are available online at https://www.mdpi.com/article/ 10.3390/mi12050480/s1, Figure S1: AFM step profiler of (a) Cu and (d) Au. Figure S2: Surface of (a) Bi₂Te₃, (b) Cu/Al₂O₃/Bi₂Te₃, (c) Ag/Al₂O₃/Bi₂Te₃, (d) Au/Al₂O₃/Bi₂Te₃. Cross section of (e) Bi₂Te₃, (f) Cu/Al₂O₃/Bi₂Te₃, (g) Ag/Al₂O₃/Bi₂Te₃, (h) Au/Al₂O₃/Bi₂Te₃.

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Conflicts of Interest: The authors declare no conflict of interest.

References

- 1. Shi, X.L.; Zou, J.; Chen, Z.G. Advanced thermoelectric design: From materials and structures to devices. *Chem. Rev.* 2020, 120, 7399–7515. [CrossRef] [PubMed]
- Xu, B.; Feng, T.; Agne, M.T.; Zhou, L.; Ruan, X.; Snyder, G.J.; Wu, Y. Highly porous thermoelectric nanocomposites with low thermal conductivity and high figure of merit from large-scale solution-synthesized Bi₂Te_{2.5}Se_{0.5} hollow nanostructures. *Angew. Chem. Int. Ed.* 2017, *56*, 3546–3551. [CrossRef] [PubMed]
- 3. Giulia, P. Thermoelectric materials: The power of pores. Nat. Rev. Mater. 2017, 2, 17006. [CrossRef]
- 4. Lee, S.; Lee, K.H.; Kim, Y.M.; Kim, H.S.; Snyder, G.J.; Baik, S.; Kim, S.W. Simple and efficient synthesis of nanograin structured single phase filled skutterudite for high thermoelectric performance. *Acta Mater.* **2018**, *142*, 8–17. [CrossRef]
- 5. Xie, K.; Gupta, M.C. Thermoelectric properties of SiGe thin films prepared by laser sintering of nanograin powders. *J. Alloys Compd.* **2020**, *820*, 153182. [CrossRef]
- Zhang, Y.; Bahk, J.H.; Lee, J.; Birkel, C.S.; Snedaker, M.L.; Liu, D.; Zeng, H.; Moskovits, M.; Shakouri, A.; Stucky, G.D. Hot carrier filtering in solution processed heterostructures: A paradigm for improving thermoelectric efficiency. *Adv. Mater.* 2014, 26, 2755–2761. [CrossRef]
- 7. Mao, J.; Chen, G.; Ren, Z. Thermoelectric cooling materials. Nat. Mater. 2020. [CrossRef]
- 8. Mu, E.; Wu, Z.; Wu, Z.; Chen, X.; Liu, Y.; Fu, X.; Hu, Z. A novel self-powering ultrathin TEG device based on micro/nano emitter for radiative cooling. *Nano Energy* **2019**, *55*, 494–500. [CrossRef]
- 9. Wu, Z.; Mu, E.; Che, Z.; Liu, Y.; Sun, F.; Wang, X.; Hu, Z. Nanoporous (00*l*)-oriented Bi₂Te₃ nanoplate film for improved thermoelectric performance. *J. Alloys Compd.* **2020**, *828*, 154239. [CrossRef]

- 10. Jin, Q.; Jiang, S.; Zhao, Y.; Wang, D.; Qiu, J.; Tang, D.M.; Tan, J.; Sun, D.M.; Hou, P.X.; Chen, X.Q. Flexible layer-structured Bi₂Te₃ thermoelectric on a carbon nanotube scaffold. *Nat. Mater.* **2019**, *18*, 62–68. [CrossRef]
- 11. Hu, X.; Hu, J.; Feng, B.; Pan, Z.; Liu, P.; Zhang, Y.; Li, R.; He, Z.; Li, G.; Li, Y. Artificial porous structure: An effective method to improve thermoelectric performance of Bi₂Te₃ based alloys. *J. Solid State Chem.* **2020**, *282*, 121060. [CrossRef]
- 12. Dun, C.; Hewitt, C.A.; Li, Q.; Guo, Y.; Jiang, Q.; Xu, J.; Marcus, G.; Schall, D.C.; Carroll, D.L. Self-assembled heterostructures: Selective growth of metallic nanoparticles on V₂–VI₃ nanoplates. *Adv. Mater.* **2017**, *29*, 1702968. [CrossRef]
- 13. Wu, Z.; Mu, E.; Wang, Z.; Chen, X.; Wu, Z.; Liu, Y.; Hu, Z. Bi₂Te₃ nanoplates' selective growth morphology on different interfaces for enhancing thermoelectric properties. *Cryst. Growth Des.* **2019**, *19*, 3639–3646. [CrossRef]
- 14. Zhu, T.; Hu, L.; Zhao, X.; He, J. New insights into intrinsic point defects in V₂VI₃ thermoelectric materials. *Adv. Sci.* **2016**, *3*, 1600004. [CrossRef]
- Zhang, M.; Liu, W.; Zhang, C.; Xie, S.; Li, Z.; Hua, F.; Luo, J.; Wang, Z.; Wang, W.; Yan, F. Identifying the manipulation of individual atomic-scale defects for boosting thermoelectric performances in artificially controlled Bi₂Te₃ Films. *ACS Nano* 2021, 15, 5706–5714. [CrossRef]
- 16. Gayner, C.; Amouyal, Y. Energy filtering of charge carriers: Current trends, challenges, and prospects for thermoelectric materials. *Adv. Funct. Mater.* **2020**, *30*, 1901789. [CrossRef]
- Zhang, Q.; Ai, X.; Wang, L.; Chang, Y.; Luo, W.; Jiang, W.; Chen, L. Improved thermoelectric performance of silver nanoparticlesdispersed Bi₂Te₃ composites deriving from hierarchical two-phased heterostructure. *Adv. Funct. Mater.* 2015, 25, 966–976. [CrossRef]
- Wu, H.J.; Yen, W.T. High thermoelectric performance in Cu-doped Bi₂Te₃ with carrier-type transition. *Acta Mater.* 2018, 157, 33–41. [CrossRef]
- 19. Hao, F.; Qiu, P.; Tang, Y.; Bai, S.; Xing, T.; Chu, H.-S.; Zhang, Q.; Lu, P.; Zhang, T.; Ren, D. High efficiency Bi₂Te₃-based materials and devices for thermoelectric power generation between 100 and 300 °C. *Energ. Environ. Sci.* **2016**, *9*, 3120–3127. [CrossRef]
- Choi, H.; Jeong, K.; Chae, J.; Park, H.; Baeck, J.; Kim, T.H.; Song, J.Y.; Park, J.; Jeong, K.H.; Cho, M.-H. Enhancement in thermoelectric properties of Te-embedded Bi₂Te₃ by preferential phonon scattering in heterostructure interface. *Nano Energy* 2018, 47, 374–384. [CrossRef]
- 21. Zhao, Y.; Li, Y.; Qiao, J.; Jiang, S.; Mao, P.; Qiu, J.; Kang, S.; Tan, J.; Tai, K.; Liu, C. Decoupling phonon and carrier scattering at carbon nanotube/Bi₂Te₃ interfaces for improved thermoelectric performance. *Carbon* **2020**, *170*, 191–198. [CrossRef]
- 22. Wu, Z.; Chen, X.; Mu, E.; Liu, Y.; Che, Z.; Dun, C.; Sun, F.; Wang, X.; Zhang, Y.; Hu, Z. Lattice strain enhances thermoelectric properties in Sb₂Te₃/Te heterostructure. *Adv. Electron. Mater.* **2020**, *6*, 1900735. [CrossRef]
- Bludska, J.; Jakubec, I.; Drašar, Č.; Lošťák, P.; Horak, J. Structural defects in Cu-doped Bi₂Te₃ single crystals. *Philos. Mag.* 2007, 87, 325–335. [CrossRef]
- 24. Keys, J.D.; Dutton, H.M. Diffusion and solid solubility of gold in single-crystal bismuth telluride. *J. Appl. Phys.* **1963**, *34*, 1830–1831. [CrossRef]
- 25. Keys, J.D.; Dutton, H.M. Diffusion and solid solubility of silver in single-crystal bismuth telluride. *J. Phys. Chem. Solids* **1963**, *24*, 563–571. [CrossRef]
- 26. Chen, S.; Hutabalian, Y.; Hu, Z.; Chen, H.; Shih, H.; Wang, W. Liquation phenomena in Sn/Bi₂Te₃, In/Bi₂Te₃ and Cu/Bi₂Te₃ couples. *Acta Mater.* **2020**, *196*, 418–429. [CrossRef]
- 27. Heyrovska, R. The Golden ratio, ionic and atomic radii and bond lengths. Mol. Phys. 2005, 103, 877–882. [CrossRef]
- 28. Minnich, A.; Dresselhaus, M.S.; Ren, Z.; Chen, G. Bulk nanostructured thermoelectric materials: Current research and future prospects. *Energ. Environ. Sci.* 2009, 2, 466–479. [CrossRef]
- 29. Ioffe, I.F. Semiconductor Thermoelements and Thermoelectric Cooling; Infosearch Ltd.: London, UK, 1957.
- 30. Yang, J.; Xi, L.; Qiu, W.; Wu, L.; Shi, X.; Chen, L.; Yang, J.; Zhang, W.; Uher, C.; Singh, D.J. On the tuning of electrical and thermal transport in thermoelectrics: An integrated theory–experiment perspective. *NPJ Comput. Mater.* **2016**, *2*, 1–17. [CrossRef]