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# Improvement of Fermi-Level Pinning and Contact Resistivity in Ti/Ge Contact Using Carbon Implantation

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**Abstract:** Effects of carbon implantation (C-imp) on the contact characteristics of Ti/Ge contact were investigated. The C-imp into Ti/Ge system was developed to reduce severe Fermi-level pinning (FLP) and to improve the thermal stability of Ti/Ge contact. The current density (*J*)-voltage (*V*) characteristics showed that the rectifying behavior of Ti/Ge contact into an Ohmic-like behavior with C-imp. The lowering of Schottky barrier height (SBH) indicated that the C-imp could mitigate FLP. In addition, it allows a lower specific contact resistivity ( $\rho_c$ ) at the rapid thermal annealing (RTA) temperatures in a range of 450–600 °C. A secondary ion mass spectrometry (SIMS) showed that C-imp facilitates the dopant segregation at the interface. In addition, transmission electron microscopy (TEM) and electron energy loss spectroscopy (EELS) mapping showed that after RTA at 600 °C, C-imp enhances the diffusion of Ge atoms into Ti layer at the interface of Ti/Ge. Thus, carbon implantation into Ge substrate can effectively reduce FLP and improve contact characteristics.





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# 1. Introduction

As a channel material for the next-generation field-effect transistors (FETs), Germanium (Ge) is considered a promising alternative to silicon (Si) owing to its higher carrier mobility and the process compatibility with the advanced Si microfabrication. However, the low-solid solubility and the high-diffusion coefficient of n-type dopants in Ge hinder the realization of low specific contact resistivity ( $\rho_c$ ) [1]. Moreover, Fermi-level pinning (FLP) caused by the metal-induced gap states (MIGS) at the metal/Ge interface is another problem to be solved [2–5]. FLP strongly occurs near the Ge valence band ( $E_v$ ) and forces the electron Schottky barrier height (e-SBH) above 0.5 eV irrespective of the metal workfunction [6]. Several approaches, including dopant segregation [7], dipole formation [8], and surface treatment [9] were proposed to mitigate FLP phenomena. Recently, the use of an ultra-thin insulator between the metal and Ge showed an effective reduction of FLP but the degradation of  $\rho_c$  due to a high tunneling resistance [10–13]. The formation of metal germanide can be another approach because the MIGS from metal dangling bond states in germanide can lead to an FLP reduction [14,15].

Ion implantation is another approach to achieving low  $\rho_c$  and suppressing dopantdiffusion behaviors. For example, Germanium implantation before silicidation induces surface amorphization to aid an epitaxial regrowth on the semiconductor surface [16]. Carbon implantation (C-imp) has been introduced in Ni-silicide and Ni-germinide contacts to reduce contact resistivity [17,18]. However, Ti/Ge contact with carbon implantation has been rarely reported.

Here, we investigated the effects of C-imp on the FLP reduction of a Ti/Ge contact and the related contact characteristics. Electrical characteristics were measured using the multiring-circular transmission line model (MR-CTLM) structure and Schottky barrier diode (SBD). Physical and structural properties of Ti/Ge contact with C-imp were analyzed using scanning electron microscopy (SEM), transmission electron microscopy (TEM), electron energy loss spectroscopy (EELS), and secondary ion mass spectrometry (SIMS).

#### 2. Materials and Methods

N-type Ge wafers moderately doped with phosphorus (~ $10^{18}$  cm<sup>-3</sup>) were cleaned in a 1:100 diluted HF (dHF) solution and deionized (DI) water to remove native oxide. Subsequently, C<sup>+</sup> ions were implanted into the Ge substrate at a dose of  $1 \times 10^{15}$  cm<sup>-2</sup> and an implantation energy of 10 keV. A reference sample without C-imp was also prepared. A SBD of Ti/Ge structure and a MR-CTLM structure were fabricated on the Ge substrate to characterize electrical properties. First, a 100 nm thick SiO<sub>2</sub> was deposited to isolate the contact holes using a plasma-enhanced chemical vapor deposition (PECVD). Then, the metal contact was formed using the conventional photolithography process. Sequentially, the oxide was etched using a dry etcher, and a Ti (5 nm)/TiN (5 nm) was deposited using a DC sputtering system. After a metal lift-off process, rapid thermal annealing (RTA) was performed in N<sub>2</sub> ambient for 60 s at 450–600 °C. Finally, a 100 nm thick Al was deposited as contact pad metal. The electrical measurements of current (*I*)–bias voltage (*V*) were performed using Keithley 4200-SCS. TEM images of the Ti/Ge structure without and with C-imp were obtained using a *JEOL JEM 2200FS* with an image Cs-corrector.

### 3. Results

Figure 1 shows the *J*-*V* characteristics of the Ti/Ge contacts with and without C-imp at RTA temperatures in a range of 450–600 °C for 60 s in N<sub>2</sub> ambient. The Ti/Ge contact without C-imp shows a typical rectifying behavior attributed to a strong FLP near the  $E_v$ , which leads to a significantly high e-SBH and reduces the reverse current density. On the other hand, the Ti/Ge contact with C-imp shows an Ohmic-like behavior with relatively high current density under the reverse regime, indicating the alleviation of FLP.



**Figure 1.** *J*-*V* characteristics of the Ti/Ge contact (**a**) without and (**b**) with C-imp at RTA temperatures in a range of 450–600 °C for 60 s in  $N_2$  ambient.

Figure 2a shows the extracted e-SBHs of the Ti/Ge contacts without (blue box) and with (red box) C-imp after RTA at 550 °C and 600 °C, respectively, for 60 s in N<sub>2</sub> ambient. The e-SBHs were extracted from the current-temperature (*I*-*T*) curves in a range of 300–378 K. The *I*-*V* relationship of a Schottky barrier diode is represented by [19]

$$I = AA^*T^2 e^{-q \otimes_B / kT} \left( e^{qV/nkT} - 1 \right) = I_{S1} e^{-q \otimes_B / kT} \left( e^{qV/nkT} - 1 \right) = I_S \left( e^{qV/nkT} - 1 \right)$$
(1)

where  $I_s$  is the saturation current, A is the diode area,  $A^* = 4\pi q k^2 m^* / h^3 = 120 \ (m^*/m)$ A/cm<sup>2</sup>·K<sup>2</sup> Richardson's constant,  $\Phi_B$  is the barrier height, and n is the ideality factor. For  $V \gg kT/q$  Equation (1) can be written as follows:

$$\ln\left(I/T^2\right) = \ln(AA^*) - q(\varnothing_B - V/n)/kT \tag{2}$$

$$\varnothing_B = \frac{V}{n} - \frac{k}{q} \frac{d[\ln(I/T^2)]}{d(1/T)} = \frac{V}{n} - \frac{2.3k}{q} \frac{d[\log(I/T^2)]}{d(1/T)}$$
(3)





Therefore, the barrier height is calculated from the slope  $(=d[\ln(I/T^2)]/d(1/T))$ . The bandgap and electron affinity in eV of Ge at 300 K are 0.66 and 4.0 eV, respectively. The workfunction of Ti metals is about 4.3 eV. When Fermi level is pinned near  $E_v$  of Ge,  $\Phi_B$  of ~0.6 eV is calculated. If there is negligible FLP,  $\Phi_B$  of ~0.3 eV is obtained.

Without C-imp, the SBH of ~0.48 eV was obtained for both 550  $^{\circ}$ C and 600  $^{\circ}$ C RTA, indicating the occurrence of FLP. In contrast, the SBH with C-imp was significantly reduced from 0.31 eV at 550  $^{\circ}$ C to 0.27 eV at 600  $^{\circ}$ C.

Figure 2b,c show schematics of the energy band diagrams for Ti/Ge contacts. Without C-imp, Fermi-level on the Ti side is pinned with the charge neutrality level ( $E_{CNL}$ ) due to FLP [6].

Figure 3 shows a top-view SEM image of the fabricated MR-CTLM structure to extract  $\rho_c$  and the sheet resistance beneath the metal ( $R_S$ ). The current flows through multiple metal-semiconductor structures from the center region to the outer-circle region. From the *I-V* curve of MR-CTLM, the total resistance ( $R_{tot}$ ) is expressed as the sum of the effective resistance ( $R_{eff}$ ) and the parasitic resistance ( $R_{pr}$ ) as follows [20]:

$$R_{tot} = R_{eff} + R_{pr} \tag{4}$$

$$R_{eff} = \frac{R_s}{2\pi} \sum_{i=0}^{9} \left[ \ln\left(\frac{r_i + S_m}{r_i}\right) + L_t \left(\frac{1}{r_i} + \frac{1}{r_i + S_m}\right) \right]$$
(5)

$$R_{pr} = \frac{R_m}{2\pi} \left[ \sum_{i=1}^9 \ln\left(\frac{r_i - L_t}{r_i - S_s + L_t}\right) \right] \tag{6}$$

where  $r_0 \sim r_9$  are the inner radius of the serial CTLM.  $S_m$  and  $S_s$  are the spacing among metal rings and dielectric rings, respectively.  $L_t$  is the transfer length.  $S_s = 10 \ \mu\text{m}$ ,  $r_0 = 50 \ \mu\text{m}$ , and  $S_m$ , from 0.5 to 10  $\mu\text{m}$  were defined using an i-line stepper.  $\rho_c$  was calculated from the  $L_t$  (=  $\sqrt{\rho_c/R_s}$ ) which was extracted by fitting a set of  $R_t$ - $S_m$  data using Equations (4)–(6).



Figure 3. Top-view SEM image of the fabricated MR-CTLM structure.

Figure 4 shows the extracted  $\rho_c$  values versus RTA temperature.  $\rho_c$  was obtained using a MR-CTLM test structure [20]. A relatively high  $\rho_c$  value seems mainly because of the low activation of a substrate doping concentration of ~1 × 10<sup>18</sup> cm<sup>-3</sup> [21,22]. After RTA annealing at 600 °C, the  $\rho_c$  values of the Ti/Ge with and without C imp were 1.3 × 10<sup>-5</sup> and 8.4 × 10<sup>-4</sup>  $\Omega$ ·cm<sup>2</sup>, respectively. Owing to the FLP effect, the Ti/Ge contact without C-imp shows  $\rho_c$  values higher than 1.0 × 10<sup>-4</sup>  $\Omega$ ·cm<sup>2</sup>.

To further analyze the effect of C-imp on the Ti/Ge composition, TEM and SIMS were conducted. The decrease of  $\rho_c$  is mainly attributed to the dopant segregation in the Ti/Ge interface [23]. In particular, for the Ti/Ge contact with C-imp after RTA at 600 °C, a further reduction of  $\rho_c$  is observed. These results can be expected by TiGe<sub>x</sub> formation. The low resistive C54-TiGe<sub>x</sub> is formed at a temperature above 600 °C [24], which mitigates FLP and improves the contact resistivity [14,15].

Figure 5a,b show SIMS profiles for Ti/Ge contacts without and with C-imp, respectively. At the Ti/Ge interface with C-imp, the peak P concentration increases from  $1.6 \times 10^{18}$  cm<sup>-3</sup> to  $3.6 \times 10^{18}$  cm<sup>-3</sup>, attributed to the dopant segregation facilitated by carbon [18]. This dopant segregation can increase the tunneling current by reducing the depletion thickness at the interface and lowering the contact resistivity.

To directly observe the microstructure of Ti/Ge contact, the cross-sectional TEM images and the corresponding EELS were analyzed. The samples were prepared after RTA at 600 °C for 60 s in N<sub>2</sub> ambient, as shown in Figure 6. In EELS maps, a bright region represents the area that the element of interest is abundant. With C-imp, Ge element is considerably observed in the Ti layer (red box in Figure 6b). The diffused Ge reacts with Ti and forms the Ti-germanide during the RTA process, which is beneficial to reduce the contact resistivity [14,15]. These results show that the C-imp is a promising approach to lower the contact resistivity in Ti/Ge contact by inducing the dopant segregation and Ge diffusion into the Ti layer.



**Figure 4.**  $\rho_c$  of the Ti/Ge contacts without (blue curve) and with (red curve) C-imp as a function of RTA temperatures ranging from 450 to 600 °C.



**Figure 5.** SIMS profiles for Ti/Ge contacts (**a**) without and (**b**) with C-imp after RTA at 600 °C. With C-imp, a dopant (phosphorous) segregation at the Ti/Ge interface is clearly observed.



**Figure 6.** Cross-sectional TEM images and corresponding electron energy loss spectroscopy (EELS) mapping images for Ge and Ti in the Ti/Ge contacts (**a**) without and (**b**) with C-imp after RTA at 600 °C.

## 4. Conclusions

We investigated the electrical and material characteristics of a Ti/Ge contact with C-imp. The current-voltage behavior shows that the carbon implantation changes the Ti/Ge rectifying behavior into an Ohmic-like behavior above RTA at 450 °C. The extracted Schottky barrier height was also decreased due to the mitigation of Fermi-level pinning. The specific contact resistivity of the Ti/Ge contact with C-imp was significantly reduced by approximately two orders of magnitude. Transmission electron microscopy and secondary ion mass spectrometry showed that carbon element at the Ti/Ge interface facilitates the dopant segregation and induces the diffusion of Ge into Ti layer. Therefore, the carbon implantation is promising to improve the Ti/Ge contact properties for high-performance Ge-FET applications.

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