

## Article

# The Influence of Preparation Temperature on the Different Facets of Bulk MgB<sub>2</sub> Superconductors

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**Abstract:** Two MgB<sub>2</sub> samples were prepared using the spark plasma sintering (SPS) technique at different temperatures—950 °C (S1) and 975 °C (S2)—for 2 h under 50 MPa pressure to study the influence of preparation temperature on different facets, namely those perpendicular (PeF) and parallel (PaF) to the compression direction of uniaxial pressure during the SPS of MgB<sub>2</sub> samples. We analyzed the superconducting properties of the PeF and PaF of two MgB<sub>2</sub> samples prepared at different temperatures from the curves of the critical temperature ( $T_C$ ), the curves of critical current density ( $J_C$ ), the microstructures of MgB<sub>2</sub> samples, and the crystal size from SEM. The values of the onset of the critical transition temperature,  $T_{c,onset}$ , were around 37.5 K and the transition widths were about 1 K, which indicates that the two samples exhibit good crystallinity and homogeneity. The PeF of the SPSed samples exhibited slightly higher  $J_C$  compared with that of the PaF of the SPSed samples over the whole magnetic field. The values of the pinning force related to parameters  $h_0$  and  $K_n$  of the PeF were lower than those of the PaF, except for  $K_n$  of the PeF of S1, which means that the PeF has a stronger GBP than the PaF. In low field, the most outstanding performance was S1-PeF, whose critical current density ( $J_C$ ) was 503 kA/cm<sup>2</sup> self-field at 10 K, and its crystal size was the smallest (0.24 μm) among all the tested samples, which is consistent with the theory that a smaller crystal size can improve the  $J_C$  of MgB<sub>2</sub>. However, in high field, S2-PeF had the highest  $J_C$  value, which is related to the pinning mechanism and can be explained by grain boundary pinning (GBP). With an increase in preparation temperature, S2 showed a slightly stronger anisotropy of properties. In addition, with an increase in temperature, point pinning becomes stronger to form effective pinning centers, leading to a higher  $J_C$ .

**Keywords:** MgB<sub>2</sub> bulk superconductor; spark plasma sintering; critical temperature; critical current density; pinning mechanism



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## 1. Introduction

Since Akimitsu [1] discovered the superconductivity of MgB<sub>2</sub>, it had been studied extensively by scientists from all over the world [2], and researchers have processed MgB<sub>2</sub> into tapes, wires [3], films [2,4], and bulks based on in situ and ex situ methods [5]. Compared to low-temperature superconductors, MgB<sub>2</sub> has a higher critical temperature ( $T_C$ ) and a higher critical current density ( $J_C$ ) [6–8]. However, it has a lower coherence length, lower anisotropy, and a good grain boundary connection compared with high-temperature superconductors [9]. More important is that its raw materials, magnesium and boron, are abundant on Earth [10], which means that its preparation is much less expensive than that of REBa<sub>2</sub>Cu<sub>3</sub>O<sub>7-x</sub> (RE represents a rare Earth element, such as Gd, Nd, Sm, etc.) bulk superconductors [11–13]. It is also speculated that, according to the raw materials of MgB<sub>2</sub> superconducting bulks, there are no toxic elements in their ingredients, which is relatively

safe and conducive to environmental protection [10]. All these characteristics make MgB<sub>2</sub> a promising superconducting material for NMR/MRI magnets [14], motors/generators [15], etc. Most of these applications need a strong trapped magnetic field that is proportional to the critical current density,  $J_C$ , and sample size according to the Bean model [16]. Therefore, when the sample size is determined, the improvement of  $J_C$  related to the effective pinning centers becomes the best way to obtain a stronger trapped magnetic field, and the preparation methods that can improve the  $J_C$  of MgB<sub>2</sub> samples become research hotspots. Therefore, several studies about the preparation methods for MgB<sub>2</sub> have been published, such as hot pressing [17], irradiation [18], doping/adding [19,20], and spark plasma sintering [6], which has not only caused the improvement of its preparation methods, but also a higher  $J_C$ , which is linked to properties and practical applications [21–23]. As one of these methods, the spark plasma sintering (SPS) technique can greatly suppress grain-coarsening and prepare high-density samples [8,21] successfully in a short time, even in less than two hours. The superconductor bulk MgB<sub>2</sub> prepared at a temperature of 850 °C following the ex situ method and using the SPS technique shows the best  $J_C$  in the whole field, compared with bulks prepared at different temperatures, namely 800 °C, 850 °C, 900 °C, and 1000 °C [5].

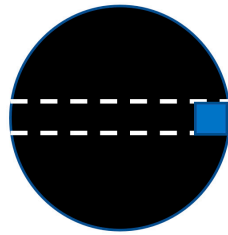
It is known that MgB<sub>2</sub> is an anisotropic material [24,25], which means its properties may be different in different directions. The SPS process is carried out under uniaxial pressure, so there are perpendicular (PeF) and parallel (PaF) facets to the compression direction of the uniaxial pressure. The effects on the different facets under various sintering durations [7] at 950 °C have been reported. In this study, two kinds of MgB<sub>2</sub> bulks are prepared at different temperatures—950 °C, marked by S1, and 975 °C marked by S2—under a uniaxial pressure of 50 MPa for 2 h using the SPS technique to analyze the properties of MgB<sub>2</sub> bulk samples. The influence of different processed temperatures on different facets of SPSed MgB<sub>2</sub> bulks is discussed from superconducting performance and microstructure points of view. Through the analysis of the flux-pinning mechanism, grain boundary pinning (GBP) and point pinning (PP) were discussed.

## 2. Experimental Section

Grade A magnesium diboride from ABCR GmbH (Karlsruhe, Germany) was used as the starting powder and then loaded into a graphite die and processed using the spark plasma sintering technique (FCT System GmbH, HD25, Rauenstein, Germany) in DC mode. A pulsed electric current (2000 A, 4 V) was passed through the sample under a dynamic vacuum ( $10^{-3}$  bar) while a 50-MPa uniaxial pressure was applied [22]. In this study, two SPSed bulk samples with disc shapes 20 mm in diameter and 3 mm in thickness were fabricated at different temperatures—950 °C (S1) and 975 °C (S2)—for 2 h under a pressure of 50 MPa.

To study the effect of preparation temperature on the performance of MgB<sub>2</sub> bulks, small specimens of rectangular shape positioned beneath the edge of MgB<sub>2</sub> bulk samples were cut into sizes of 2 mm × 2 mm × 1 mm from these bulk samples for the measurement of superconducting properties, as shown in Figure 1. The phase position and lattice parameters were confirmed using X-ray diffraction (XRD) at room temperature. Scanning electron microscopy (SEM, microscope model, company, city, country) was used to observe the microstructure of these samples, and the crystal size of the SPSed MgB<sub>2</sub> was calculated according to the SEM using Nano Measurer 1.2. The element analysis was conducted using energy-dispersive X-ray spectroscopy (EDS). The DC magnetization measurement was measured using a Quantum Design SQUID magnetometer under a magnetic field perpendicular to the tested facet for the superconducting transition and magnetic hysteresis loops of MgB<sub>2</sub> samples. The values of critical current density ( $J_C$ ) were calculated based on the extended Bean's critical state model [16,26]. The facets perpendicular to the compression direction in the SPS process were marked as perpendicular facets (PeF), and the facets parallel to the compression direction were marked as parallel facets (PaF). The two kinds of

facets were compared to study the differences in superconducting properties between the PeF and PaF.

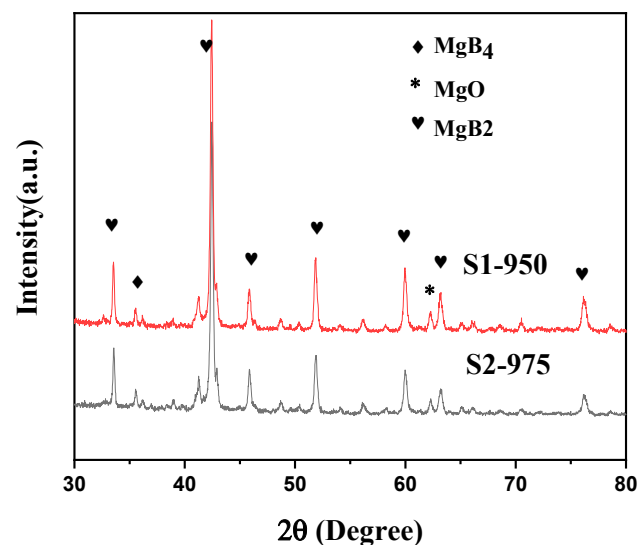


**Figure 1.** Schematic diagram of specimens cut in the position of both  $\text{MgB}_2$  bulks.

### 3. Results

#### 3.1. Superconducting Properties of the PeF and PaF of $\text{MgB}_2$ Samples at Different Preparation Temperatures ( $950\text{ }^\circ\text{C}$ and $975\text{ }^\circ\text{C}$ )

Figure 2 shows the room-temperature XRD patterns of samples prepared at 50 MPa uniaxial pressure and different temperatures,  $950\text{ }^\circ\text{C}$  (S1) and  $975\text{ }^\circ\text{C}$  (S2), using spark plasma sintering technology. It can be seen from Figure 2 that the main component of the samples is  $\text{MgB}_2$ , although there are a few second phases, such as  $\text{MgB}_4$  and  $\text{MgO}$ . The oxidation reaction appears under the oxygen condition during growth [21,27]. High temperature and pressure conditions during the process [27,28] may result in the generation of  $\text{MgB}_4$ . These kinds of particles may have an important effect on the superconducting properties of SPSed samples.



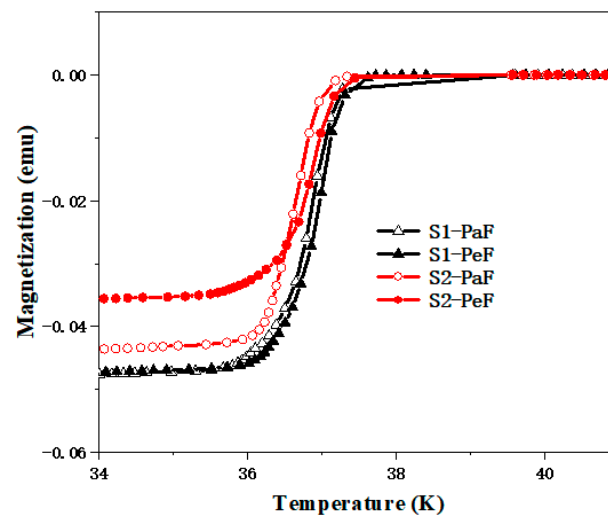
**Figure 2.** Room-temperature X-ray diffraction patterns of samples prepared using the spark plasma sintering method at  $950\text{ }^\circ\text{C}$  (S1) and  $975\text{ }^\circ\text{C}$  (S2) for 2 h under 50 MPa pressure.

The standard lattice parameter values of  $\text{MgB}_2$  without doping any other substances are  $a = 3.086\text{ \AA}$ ,  $c = 3.524\text{ \AA}$ . Table 1 shows that the lattice parameters of the samples at different temperatures,  $950\text{ }^\circ\text{C}$  (S1) and  $975\text{ }^\circ\text{C}$  (S2), are  $a = 3.0847\text{ \AA}$ ,  $3.0811\text{ \AA}$ ,  $c = 3.5232\text{ \AA}$ , and  $3.5223\text{ \AA}$ , respectively. A small a-axis lattice parameter is reflected in the substitution of carbon from the graphite mold system for boron in the crystal lattice of  $\text{MgB}_2$  during the process [29,30]. The lattice parameters of S1 are closer to the standard values of  $\text{MgB}_2$ . Therefore, we can infer that the performance of the S1 sample may be slightly better than that of the S2 sample.

**Table 1.** The lattice parameters and average grain sizes of two crystal facets (PaF and PeF) of samples prepared using SPS technology at 950 °C (S1) and 975 °C (S2) for 2 h under 50 MPa pressure.

Sample	Lattice Parameters (Å)	Average Particle Size (µm)
	a/c	PaF/PeF
S1 (950 °C)	3.0847/3.5232	0.27/0.24
S2 (975 °C)	3.0811/3.5223	0.29/0.25

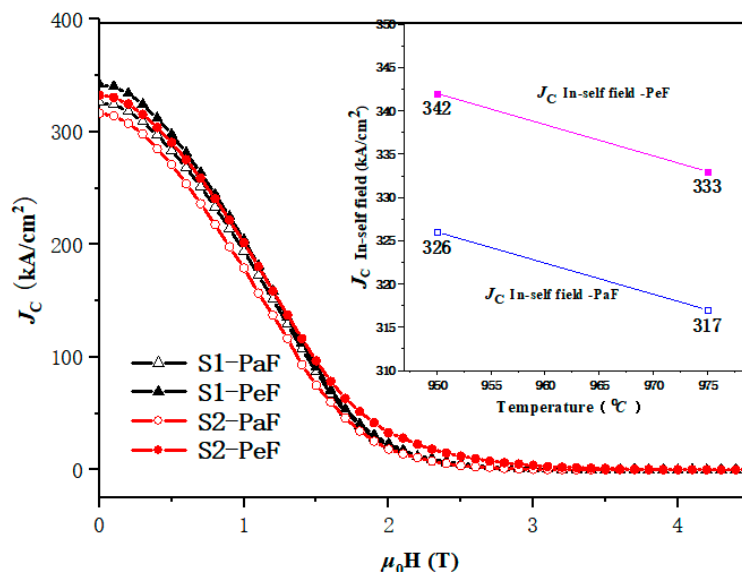
Figure 3 shows the sample temperature dependence of magnetization curves of MgB<sub>2</sub> samples prepared at different temperatures of 950 °C (S1) and 975 °C (S2) for 2 h under 50 MPa pressure. The onset of the critical transition temperature,  $T_{C,onset}$ , is at around 37.5 K and the transition width is about 1 K, which indicates the two samples exhibit good crystallinity and homogeneity. The  $T_{C,onset}$  of S1 is slightly higher than that of the S2 sample, which may be related to the good performance of the S1 sample.



**Figure 3.** The  $T_C$  (temperature–magnetization) curves of the SPSed MgB<sub>2</sub> bulk samples prepared at different temperatures of 950 °C (S1) and 975 °C (S2) for 2 h under 50 MPa pressure.

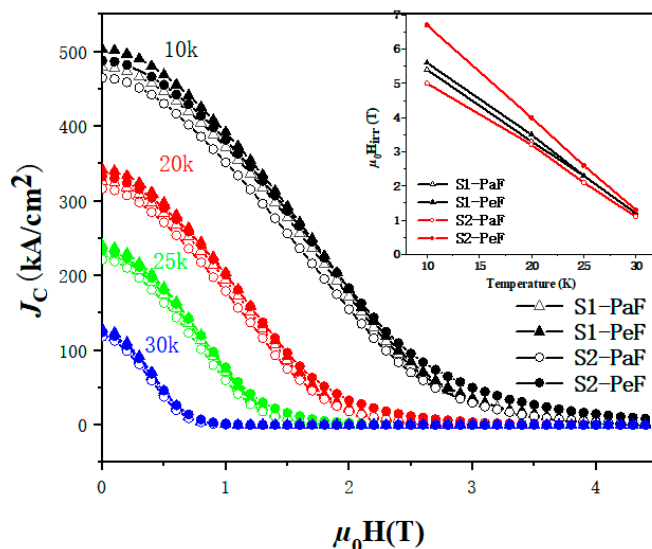
Figure 4 shows the curves of the magnetic field and  $J_C$  from different facets (PeF and PaF) measured at 20 K, and the embedded diagram shows their self-field  $J_C$  at different preparation temperatures. It can be seen that  $J_C$  decreases with the increase in the magnetic field. The PeF of the SPSed samples exhibits a slightly higher  $J_C$  compared with that of the PaF of the SPSed samples in the whole magnetic field. In the low magnetic field, the highest self-field  $J_C$  is 342 kA/cm<sup>2</sup> in the S1-PeF samples, as shown in Figure 4. The self-field  $J_C$  of the S2-PeF sample is 333 kA/cm<sup>2</sup>. The self-field of  $J_C$  in the PaF of the S1 and S2 samples is 326 kA/cm<sup>2</sup> and 317 kA/cm<sup>2</sup>, respectively. In the high magnetic field, the increase of  $J_C$  appears in the PeF of S2 samples, which is important for the future application of MgB<sub>2</sub> bulks. It is clear that the PeF and PaF of S1 processed under lower prepared temperatures possess the optimum superconducting performance in a low magnetic field; however,  $J_C$  of the PeF of the S2 sample processed under higher prepared temperatures is more advantageous in a high magnetic field, which indicates that the preparation temperatures (950 °C and 975 °C) of the samples are appropriate. According to the report [5], the superconductor bulks MgB<sub>2</sub> prepared at a temperature of 850 °C with the ex situ method using the SPS technique show the best  $J_C$  in the whole magnetic field compared with the bulks prepared at different temperatures. The value of  $J_C$  increases first and decreases then with the increase in the preparation temperature to 800 °C, 850 °C, 900 °C and 1000 °C, which further suggests that a suitable preparation temperature may exhibit a better superconducting property. Thus, it cannot be fully proved that a higher

preparation temperature is harmful to the superconducting properties of MgB<sub>2</sub> bulk in our work; further study, with a wider preparation temperature range, is required.



**Figure 4.** The  $J_C$ - $\mu_0 H$  curves of magnetic field and critical current density at 20 K of two surfaces, PeF and PaF, of SPSed MgB<sub>2</sub> bulk samples. The inset represents the critical current density measured at the self-field.

$J_C$  of the two samples at the different processed temperatures as a function of the applied field is given in Figure 5. In addition, the embedded diagram shows their irreversible field, represented by  $\mu_0 H_{irr}$ , which is defined as the field when  $J_C$  reaches 100 A/cm<sup>2</sup>, under different processed temperatures. It is clear that the  $J_C$  of the S1-PeF possesses the highest  $J_C$  in a low magnetic field, and the S2-PeF enhances the  $J_C$  in a high magnetic field, accompanied by the highest  $\mu_0 H_{irr}$  at 20 K. The significant enhancement of  $J_C$  at a temperature of 10 K appears in a high magnetic field, which will be beneficial to future applications of MgB<sub>2</sub> bulks.

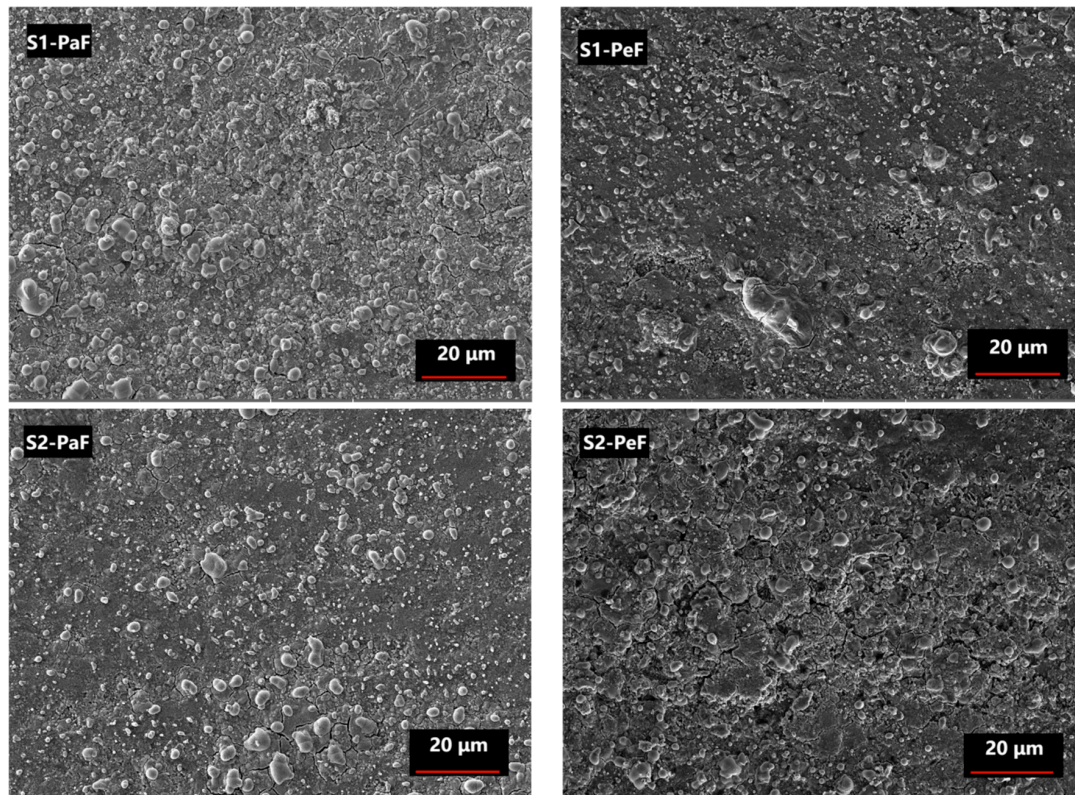


**Figure 5.** The  $J_C$ - $\mu_0 H$  curves of the magnetic field and critical current density at different temperatures (10 K, 20 K, 25 K, and 30 K) of two surface PeF and PaF, of SPSed MgB<sub>2</sub> bulk samples. The inset shows the values of  $J_C$  in an irreversible field under different temperatures.



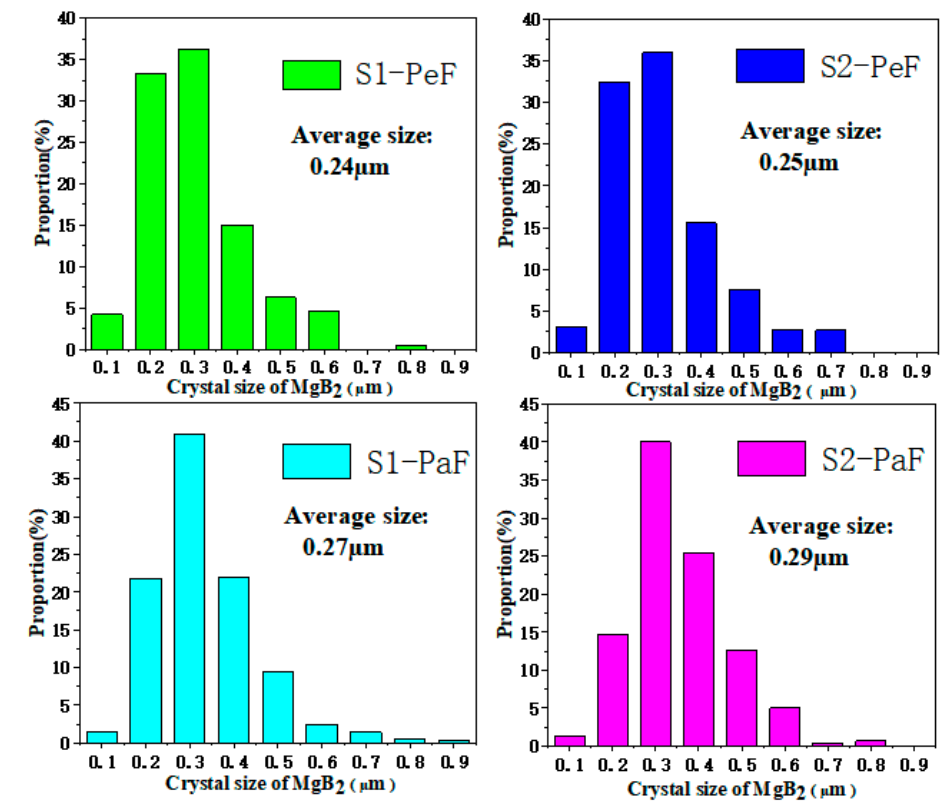
### 3.2. The Microstructure and Crystal Size of MgB<sub>2</sub> Samples at Different Preparation Temperatures (950 °C and 975 °C)

Figure 6 shows a SEM diagram of the PaF (left) and PeF (right) of the SPSed samples. It is found that some large and white round and oval particles appear in the PaF of both samples, which are oxidative metamorphic MgB<sub>2</sub> grains or MgO secondary particles [7,8]. It has been shown that the size of the effective pinning centers is similar to the coherence length of normal MgB<sub>2</sub> (about 12 nm) [31], which results in the improvement of superconducting properties. Most oxidative metamorphic MgB<sub>2</sub> grains and MgO secondary particles are too large to be effective pinning centers, which are harmful to the connectivity between normal MgB<sub>2</sub> grains and have a negative effect on  $J_C$ , as shown in Figure 4, with the decrease of  $J_C$  in the PaF of both samples. There are more grain boundaries in the PeF of both samples, which are prone to GBP to improve the superconducting properties of MgB<sub>2</sub> bulks [28,29], which is consistent with the higher  $J_C$  of the PeF in both samples, as shown in Figures 4 and 5. Moreover, the existence of more small-sized particles in the PeF of the S1 sample may be related to the highest self-field  $J_C$ , as shown in Figure 4.



**Figure 6.** SEM diagrams from the PaF (left) and the PeF (right) of the SPSed MgB<sub>2</sub> bulks prepared under 50 MPa pressure at 950 °C (S1) and 975 °C (S2), respectively.

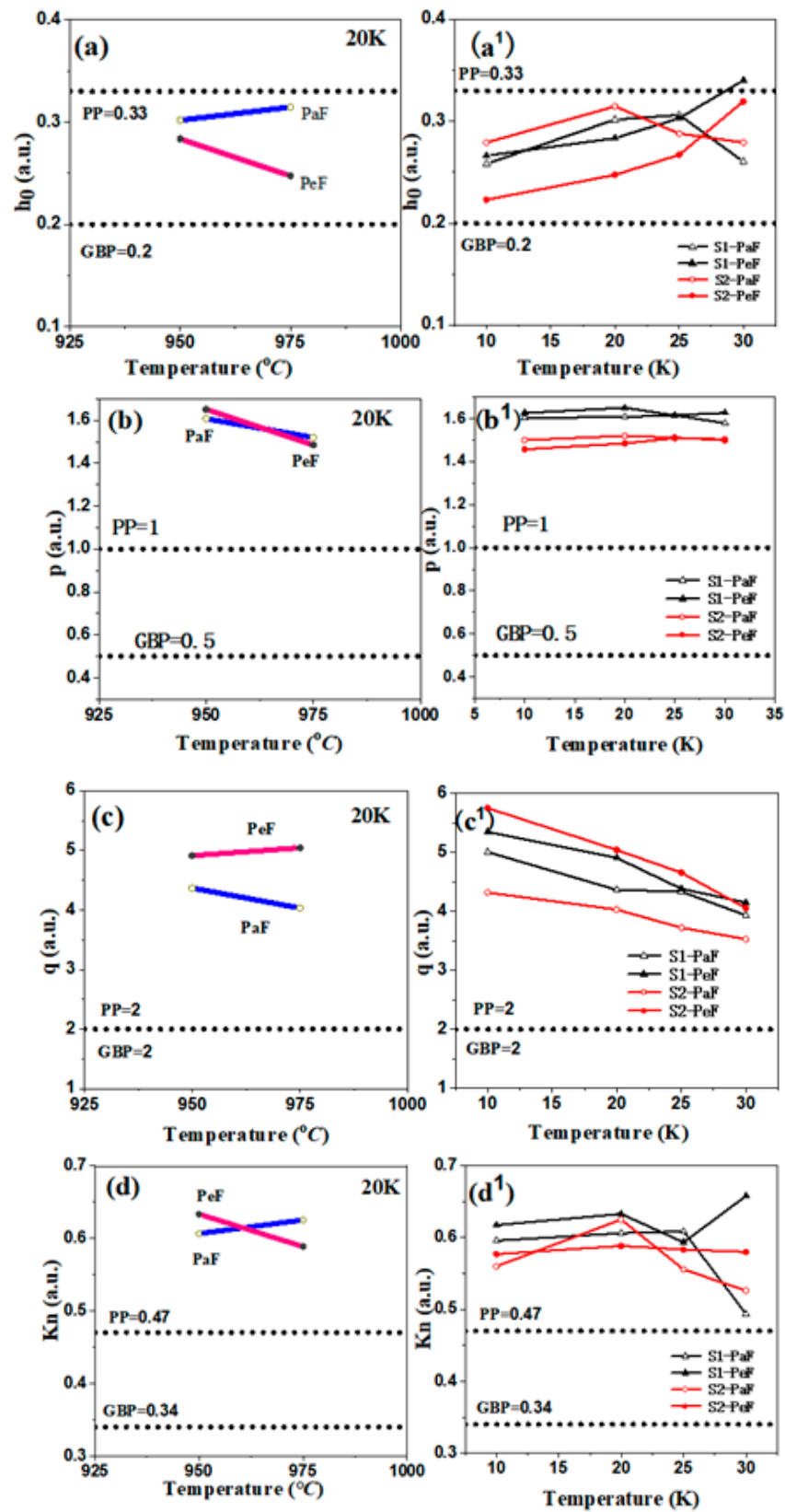
Figure 7 shows the statistical results of the MgB<sub>2</sub> crystal size in the SEM pictures (Figure 6) based on Nano Measurer 1.2 software. The average sizes of samples are listed in Table 1. The horizontal axis is the range of the crystallite size, and the vertical axis is the percentage of the number of MgB<sub>2</sub> grains in a certain size range based on the total statistical MgB<sub>2</sub> grains in each specimen. It is easy to see that S1-PeF has the smallest average size (0.24 μm), and S2-PaF has the largest average size (0.29 μm). A smaller crystal size means more MgB<sub>2</sub> particles and grain boundaries per unit volume, which can increase the number of effective pinning centers and further improve the  $J_C$  of MgB<sub>2</sub> [32]. This is consistent with the trend of  $J_C$  in Figures 3 and 4, which confirms that the  $J_C$  is greatly influenced by the crystal size of the SPSed MgB<sub>2</sub> bulks.



**Figure 7.** The crystal size distribution and average size of SPSed  $\text{MgB}_2$  bulks measured by Nano Measurer 1.2 software.

### 3.3. The Flux-Pinning Mechanism

To study the flux-pinning mechanism, two scaling procedures should be used to identify the dominant pinning mechanism from the peak position [33,34]. According to Dew-Hughes [33], one of them is the universal law  $F_p = Ah^p(1-h)^q$  (where  $A$  is constant, and  $F_p$  is the volume pinning force. Another equation is  $F_p = \mu_0 H \times J_C$  and  $h$  is a reduced field; it can be presented by the equation  $h = H/H_{irr}$ ). The law was chosen to fit the data of  $F_p$  and  $H/H_{irr}$ . The parameter  $h_0$  is the field when  $F_p$  reaches its maximum, and the exponents  $p$  and  $q$  are fitted from the data of  $F_p$  and  $H/H_{irr}$ , which analyzes the flux-pinning behavior more deeply and easily. Figure 8 shows the curves of the parameters related to the pinning force:  $h_0$  (a),  $p$  (b), and  $q$  (c) as a function of the preparation temperature and  $h_0$  ( $a_1$ ),  $p$  ( $b_1$ ), and  $q$  ( $c_1$ ) as a function of test temperature. The theoretical values of the GBP and PP for reference are represented by the dotted lines. For isotropic materials, GBP is identified by  $h_0 = 0.2$ ,  $p = 0.5$ , and  $q = 2$ , while PP is identified by  $h_0 = 0.33$ ,  $p = 1$ , and  $q = 2$ . However, this scaling procedure has its limitation for  $\text{MgB}_2$  materials, because it does not apply to untextured, anisotropic materials. Therefore, another scaling procedure, whose parameter  $K_n$  is presented by  $K_n = h_0/h_n$  (where  $h_n$  refers to the field when  $F_p$  is halved), was proposed by Eisterer [34], to analyze the flux-pinning mechanism to achieve a more accurate flux-pinning mechanism behavior. As a modified scaling procedure, it can greatly reduce the influence of previously unknown parameters of anisotropy. The  $K_n$  values of 0.34 and 0.47 are defined for GBP and PP, respectively, and the plots of  $K_n$  (d) and ( $d_1$ ) as functions of the preparation temperature and the test temperature are shown in Figure 8. These two procedures mentioned above will be combined to discuss the pinning mechanism of the samples.



**Figure 8.** The plots of the parameters related to the pinning force:  $h_0$  (a),  $p$  (b),  $q$  (c), and  $K_n$  (d) as functions of preparation temperature at 20 K;  $h_0$  (a<sup>1</sup>),  $p$  (b<sup>1</sup>),  $q$  (c<sup>1</sup>), and  $K_n$  (d<sup>1</sup>) as functions of test temperature. The theoretical values of grain boundary pinning (GBP) and point pinning (PP) for reference are represented by dotted lines.



It is obvious from Figure 8a,d that the values of the pinning force related to parameters  $h_0$  and  $K_n$  of the PeF are lower than those of the PaF, except for  $K_n$  of the PeF of S1, which means that the PeF has a stronger GBP than the PaF. Based on the previous report [7,35], the smaller crystallite size causes more grain boundaries per unit volume, which means a stronger GBP, as shown in the microstructure of Figure 6 and the smaller crystallite size of Figure 7, accompanied by a better superconducting performance of the PeF in both samples. The values of the pinning force-related parameters  $h_0$  and  $K_n$  of the PeF of S2 decrease with the increase in the preparation temperature of the  $MgB_2$  bulk samples, which means that S2 has a stronger GBP than S1. We can speculate that the crystallite size of S2 is smaller than the crystallite size of S1. In addition, according to this deduction,  $J_C$  of S2 should be higher than that of S1 in a magnetic field, but the trend of  $J_C$  is opposite to this rule, as shown in Figures 4 and 5. This represents that GBP is not the only factor affecting  $J_C$  in the PeF of S1. There must be other factors that influence it—maybe the secondary phase particles in the suitable size, such as MgO or  $MgB_4$ , etc. The right picture shows the chemical composition of the area circled in red in the left picture, as shown in Figure 9. From the EDS spectra, we can calculate the chemical formula and then verify the existence of MgO particles, which is confirmed by the XRD analysis of Figure 2. MgO particles can be effective pinning centers. In addition, both PP and GBP work on  $J_C$  at the same time to result in the improvement of  $J_C$  in the PeF of the S1 sample under the lower preparation temperature, as shown in Figures 4 and 5. The discrepancy of S2 under higher preparation temperature in the parameters  $h_0$  and  $K_n$  is higher than that of S1, which is related to the increasing anisotropy of the SPSed samples, which accounts for the difference in properties between the PeF and PaF of S2. This is not conducive to practical application.

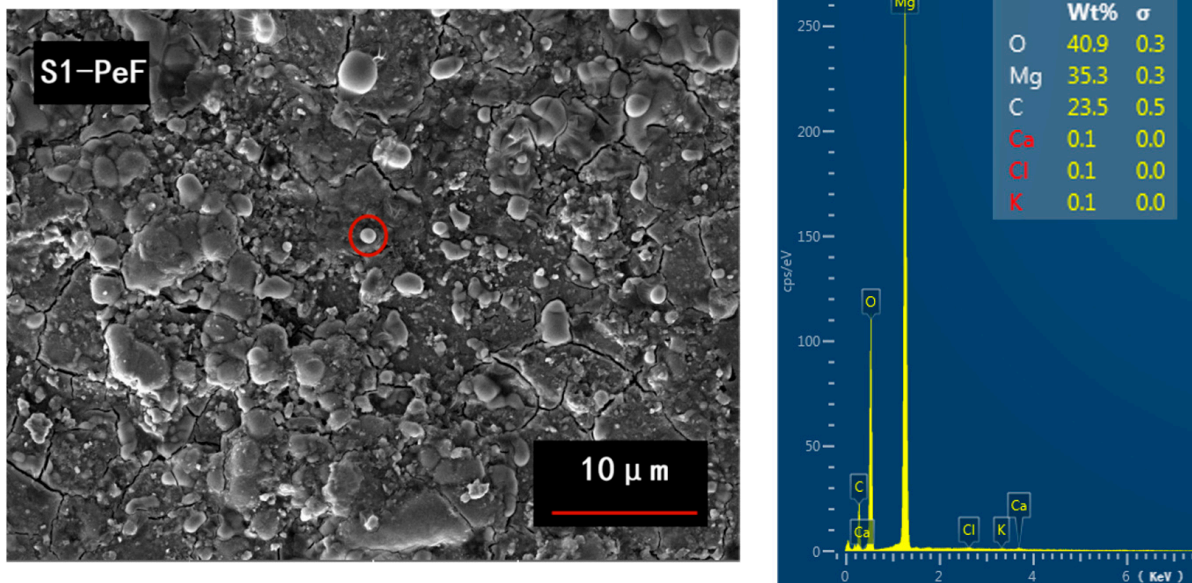


Figure 9. The EDS spectra (right) of the area circled in red (left).

The relationship between the pinning mechanism and the critical current of  $MgB_2$  bulks has been studied, and further investigation of the relationship between pinning force-related parameters and  $J_C$  in high magnetic fields has been studied as well. It is inferred that more grain boundaries can enhance the  $H_{C2}$  [33,36], which can be presented by  $H_{irr}$  [30,36], so we can predict that stronger GBP can lead to higher  $H_{irr}$  and then improve  $J_C$  in high magnetic fields [7]. This coincides with the result, as shown in Figures 4 and 5, which has the best performance in a high magnetic field in  $J_C$  of S2-PeF.

It is evident that the values of  $h_0$  and  $K_n$  in the PeF of both samples, as shown in Figure 8a<sup>1</sup>, become larger with the increase in the test temperature, from 10 K to 30 K, which means the PP grows stronger. The stronger PP indicates that much more effective

pinning centers exist, and these effective pinning centers can bring higher  $J_C$ . This is consistent with our observation of the MgO second phases in Figure 2, which can be effective pinning centers. This theory is consistent with the performance of  $J_C$ , as shown in Figure 5. We can notice that the values of  $h_0$  and  $K_n$  of the PaF in both samples from Figure 8a<sup>1</sup>,d<sup>1</sup> are not a monotonous trend, but rather a fluctuation, which is not consistent with the conclusion that PP as the main pinning becomes stronger as the test temperature increases, which means that the values of  $h_0$  and  $K_n$  show an overall increasing trend with increasing test temperature [19,36,37]. Maybe some evidence can be found from some studies with small temperature intervals [21,37]. The change in anisotropy or percolation with temperature may cause this temperature dependence of the pinning mechanism [7].

#### 4. Conclusions

In this study, we prepared two MgB<sub>2</sub> bulk samples, with a disc shape of 20 mm in diameter and 3 mm in thickness, at different temperatures of 950 °C (marked S1) and 975 °C (marked S2) for 2 h under 50 MPa pressure using the spark plasma sintering technique. The properties of the samples were analyzed according to the results of magnetization measurement and SEM diagrams. Critical temperature ( $T_C$ ), critical current density ( $J_C$ ), microstructure, and the pinning mechanism of the perpendicular facets (PeF) and parallel facets (PaF) to the compression direction of MgB<sub>2</sub> bulk samples are all included. The onset of critical temperature is around 37.5 K and the transition width is about 1 K, which indicates all samples exhibit good crystallinity and homogeneity. In terms of  $T_C$ , the onset transition temperature,  $T_{C,onset}$ , of S1 is slightly higher than that of the S2 sample, which may be related to the good performance of the S1 sample. In terms of  $J_C$ , the critical current density of all MgB<sub>2</sub> bulk samples decreases with the increase of the applied field. In a low magnetic field, S1-PeF has the highest  $J_C$  (342 kA/cm<sup>2</sup>), followed by S2-PeF (333 kA/cm<sup>2</sup>), S1-PaF (326 kA/cm<sup>2</sup>), and S2-PaF (317 kA/cm<sup>2</sup>) in self-field at 20 K, respectively. The crystal size of the samples is 0.24 μm (S1-PeF), 0.25 μm (S2-PeF), 0.27 μm (S1-PaF), and 0.29 μm (S2-PaF), respectively. Thus, the values of  $J_C$  and the crystal size of samples obey the rule that a smaller crystal size can improve the  $J_C$  of MgB<sub>2</sub> samples. However, in a high magnetic field, the S2-PeF has the highest  $J_C$ , which contradicts the rule above. It illustrates that other factors can affect  $J_C$ . According to the pinning mechanism, we find the factor is GBP, where stronger GBP can improve the  $J_C$  in a high magnetic field. In terms of the pinning mechanism, according to the analysis of the trend of the parameters  $h_0$ ,  $p$ ,  $q$ , and  $K_n$ , the PeF has a stronger GBP than the PaF, which is related to the smaller crystallite size, resulting in more grain boundaries per unit volume. Therefore, PeF shows better superconducting performance with an increase in the preparation temperature. The existence of a slightly higher discrepancy of properties occurs in the S2 sample. In addition, with the increase in temperatures, PP grows stronger, bringing higher  $J_C$ . It is clear that the PeF and PaF of S1 processed under a lower prepared temperature possess the optimum superconducting performance in a low magnetic field; however,  $J_C$  of the PeF of the S2 sample processed under a higher prepared temperature is more advantageous in a high magnetic field. Therefore, further suitable preparation temperatures to fabricate MgB<sub>2</sub> superconductor bulks using the SPS technique will appear in future experiments.

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## References

1. Nagamatsu, J.; Nakagawa, N.; Muranaka, T.; Zenitani, Y.; Akimitsu, J. Superconductivity at 39 K in magnesium diboride. *Nature* **2001**, *410*, 63. [\[CrossRef\]](#)
2. Guo, C.; Wang, H.Z.; Cai, X.W.; Luo, W.H.; Huang, Z.G.; Zhang, Y.; Feng, Q.R.; Gan, Z.Z. High performance superconducting joint for MgB<sub>2</sub> films. *Phys. C* **2021**, *584*, 1353863. [\[CrossRef\]](#)
3. Jin, S.; Mavoori, H.; Bower, C.; Van Dover, R.B. High critical currents in iron-clad superconducting MgB<sub>2</sub> wires. *Nature* **2001**, *411*, 563. [\[CrossRef\]](#)
4. Eom, C.B.; Lee, M.K.; Choi, J.H. High critical current density and enhanced irreversibility field in superconducting MgB<sub>2</sub> thin films. *Nature* **2001**, *411*, 558. [\[CrossRef\]](#) [\[PubMed\]](#)
5. Dadiel, J.L.; Naik, S.P.K.; Peczowski, P.; Sugiyama, J.; Ogino, H.; Sakai, N.; Kazuya, Y.; Warski, T.; Wojcik, A.; Oka, T.; et al. Synthesis of Dense MgB<sub>2</sub> Superconductor via In-Situ and Ex-Situ Spark Plasma Sintering Method. *Materials* **2021**, *14*, 7395. [\[CrossRef\]](#)
6. Takahashi, Y.; Naito, T.; Fujishiro, H. Trapped Field Properties of MgB<sub>2</sub> Bulks Prepared Via an in-situ Infiltration-Reaction Process Using Refined Boron Powders. *IEEE Trans. Appl. Supercond.* **2022**, *32*, 6801005. [\[CrossRef\]](#)
7. Zhang, J.Y.; Zhang, Y.F.; Lou, Z.W.; Zhang, P.H.; Li, C.Y.; Yuan, J.W.; Peng, L.; Ma, Y.X.; Noudem, J.G.; Izumi, M. The discrepancies in different facets of MgB<sub>2</sub> bulk superconductors prepared under various sintering durations by spark plasma sintering. *Supercond. Sci. Technol.* **2021**, *34*, 045011. [\[CrossRef\]](#)
8. Xing, Y.; Bernstein, P.; Miryala, M.; Noudem, J.G. High Critical Current Density of Nanostructured MgB<sub>2</sub> Bulk Superconductor Densified by Spark Plasma Sintering. *Nanomaterials* **2022**, *12*, 2583. [\[CrossRef\]](#) [\[PubMed\]](#)
9. Buzea, C.; Yamashita, T. Review of superconducting properties of MgB<sub>2</sub>. *Supercond. Sci. Technol.* **2001**, *14*, R115–R146. [\[CrossRef\]](#)
10. Noudem, J.G.; Dupont, L.; Gozzelino, L.; Bernstein, P. Superconducting Properties of MgB<sub>2</sub> Bulk Shaped by Spark Plasma Sintering. *Mater. Today-Process.* **2016**, *3*, 545–549. [\[CrossRef\]](#)
11. Zhang, Y.F.; Lou, Z.W.; Zhang, P.H.; Li, C.Y.; Zhang, J.Y.; Zhang, X.J. Flux Pinning Properties of Single-Grain Bulk GdBCO Superconductors Processed by Different Thicknesses of Y123 Liquid Source. *Micromachines* **2022**, *13*, 701. [\[CrossRef\]](#) [\[PubMed\]](#)
12. Zhu, Y.H.; Zmorayová, K.; He, J.Y.; Zhang, Y.X.; Diko, P.; Zhou, D.F.; Yao, X. An in situ self-assembly strategy for exact-(110)-plane-controlled crystallization of high-performance YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7–δ</sub> single. *Ceram. Int.* **2022**, *48*, 22196–22205. [\[CrossRef\]](#)
13. Mishima, F.; Terada, T.; Akiyama, Y.; Nishijima, S. High gradient superconducting magnetic separation for iron removal from the glass polishing waste. *IEEE Trans. Appl. Supercond.* **2011**, *21*, 2059–2062. [\[CrossRef\]](#)
14. Moseley, D.A.; Wilkinson, D.P.; Mousavi, T.; Dennis, A.R.; Speller, S.; Durrell, J.H. A new MgB<sub>2</sub> bulk ring fabrication technique for use in magnetic shielding or bench-top NMR system. *Supercond. Sci. Technol.* **2022**, *35*, 085003. [\[CrossRef\]](#)
15. Kajikawa, K.; Nakamura, T.; Kobayashi, H.; Wakuda, T.; Tanaka, K. Development of Stator Windings for Fully Superconducting Motor with MgB<sub>2</sub> Wires. *IEEE Trans. Appl. Supercond.* **2013**, *23*, 2240033. [\[CrossRef\]](#)
16. Bean, C.P. Magnetization of hard superconductors. *Phys. Rev. Lett.* **1962**, *8*, 250–253. [\[CrossRef\]](#)
17. Hassler, W.; Scheiter, J.; Hadrich, P.; Kauffmann-Weiss, S.; HolzaPaFel, B.; Oomen, M.; Nielsch, K. Properties of ex-situ MgB<sub>2</sub> bulk samples prepared by uniaxial hot pressing and spark plasma sintering. *Phys. C Supercond. Appl.* **2018**, *551*, 48–54. [\[CrossRef\]](#)
18. Erdem, O.; Yanmaz, E. Effect of laser irradiation on activation energy, irreversibility field and upper critical magnetic field of bulk MgB<sub>2</sub> superconductor. *J. Mater. Sci.* **2016**, *27*, 6502–6510. [\[CrossRef\]](#)
19. Koblischka, M.R.; Koblischka-Veneva, A.; Miryala, M.; Murakami, M. Magnetic Characterization of Bulk C-Added MgB<sub>2</sub>. *J. IEEE Trans. Appl. Supercond.* **2019**, *29*, 6800104. [\[CrossRef\]](#)
20. Erdem, O.; Guner, S.; Celik, S.; Kucukomeroglu, T. Superconducting and levitation force characterisation of pyrene added MgB<sub>2</sub> bulk superconductors. *Cryogenics* **2020**, *112*, 103205. [\[CrossRef\]](#)
21. Badica, P.; Burdusel, M.; Popa, S.; Pasuk, I.; Ivan, I.; Borodianska, H.; Vasylykiv, O.; Kuncser, A.; Ionescu, A.M.; Miu, L.; et al. Reactive spark plasma sintering of MgB<sub>2</sub> in nitrogen atmosphere for the enhancement of the high-field critical current density. *Supercond. Sci. Tech.* **2016**, *29*, 105020. [\[CrossRef\]](#)
22. Aldica, G.; Batalu, D.; Popa, S.; Ivan, I.; Nita, P.; Sakka, Y.; Vasylykiv, O.; Miu, L.; Pasuk, I.; Badica, P. Spark plasma sintering of MgB<sub>2</sub> in the two-temperature route. *Phys. C Supercond. Appl.* **2012**, *477*, 43–50. [\[CrossRef\]](#)
23. Jirsa, M.; Rames, M.; Koblischka, M.R.; Koblischka-Veneva, A.; Berger, K.; Douine, B. Douine. Relaxation and pinning in spark-plasma sintered MgB<sub>2</sub> superconductor. *Supercond. Sci. Technol.* **2016**, *29*, 1. [\[CrossRef\]](#)
24. Noudem, J.G.; Bernstein, P.; Dupont, L.; Martin, F.G.R.; Sotelo, G.G.; Dias, D.H.N.; de Andrade, R.; Muralidhar, M.; Murakami, M. Spark plasma sintering of bulk MgB<sub>2</sub> and levitation force measurements. *Supercond. Sci. Technol.* **2020**, *33*, 024001. [\[CrossRef\]](#)
25. Sandu, V.; Ivan, I. On the scaling of pinning force in ceramic MgB<sub>2</sub>. *Int. Conf. Low Temp. Phys.* **2012**, *400*, 22102–22105. [\[CrossRef\]](#)
26. Chen, D.X.; Goldfarb, R.B. Kim model for magnetization of type-II superconductors. *J. Appl. Phys.* **1989**, *66*, 2489–2500. [\[CrossRef\]](#)
27. Guo, Y.; Zhang, W.; Yang, D.; Yao, R.-L. Decomposition and oxidation of magnesium diboride. *J. Am. Ceram. Soc.* **2012**, *95*, 754–759. [\[CrossRef\]](#)

28. Kim, S.; Stone, D.S.; Cho, J.-I.; Jeong, C.-Y.; Kang, C.-S.; Bae, J.-C. Phase stability determination of the Mg-B binary system using the CALPHAD method and ab initio calculations. *J. Alloy. Compd.* **2009**, *470*, 85–89. [[CrossRef](#)]
29. Aldica, G.; Burdusel, M.; Popa, S.; Enculescu, M.; Pasuk, I.; Badica, P. The influence of heating rate on superconducting characteristics of MgB<sub>2</sub> obtained by spark plasma sintering technique. *Phys. C Supercond. Appl.* **2015**, *519*, 184–189. [[CrossRef](#)]
30. Badica, P.; Burdusel, M.; Popa, S.; Hayasaka, Y.; Ionescu, M.A.; Aldica, G. Addition of Sb<sub>2</sub>O<sub>5</sub> into MgB<sub>2</sub> Superconductor Obtained by Spark Plasma Sintering. *Supercond. Novel Magn.* **2017**, *30*, 2073–2080. [[CrossRef](#)]
31. Shimada, Y.; Hata, S.; Ikeda, K.-I.; Nakashima, H.; Matsumura, S.; Tanaka, H.; Yamamoto, A.; Shimoyama, J.-I.; Kishio, K. Microstructural Characteristics of Ball-Milled Self-Sintered Ex Situ MgB<sub>2</sub> Bulks. *IEEE Trans. Appl. Supercond.* **2015**, *25*, 1–5. [[CrossRef](#)]
32. Shah, M.S.; Shahabuddin, M.; Parakkandy, J.M.; Qaid, S.; Alzayed, N.S. Enhanced critical current density in undoped MgB<sub>2</sub> prepared by in situ/ex situ combination technique. *Solid State Commun.* **2015**, *218*, 31–34. [[CrossRef](#)]
33. Dew-Hughes, D. Flux pinning mechanisms in type II superconductors. *Philos. Mag.* **1974**, *30*, 293305. [[CrossRef](#)]
34. Eisterer, M. Calculation of the volume pinning force in MgB<sub>2</sub> superconductors. *Phys. Rev. B* **2008**, *77*, 144524. [[CrossRef](#)]
35. Ma, Z.; Liu, Y.; Han, Y.; Zhao, Q.; Gao, Z. Variation of the enhancement mechanism in the critical current density of Cu-doped MgB<sub>2</sub> samples sintered at different temperatures(Article). *J. Appl. Phys.* **2008**, *104*, 063917. [[CrossRef](#)]
36. Batalu, D.; Aldica, G.; Popa, S.; Kuncser, A.; Mihalache, V.; Badica, P. GeO<sub>2</sub>-added MgB<sub>2</sub> superconductor obtained by Spark Plasma Sintering. *Solid State Sci.* **2015**, *48*, 23–30. [[CrossRef](#)]
37. Aldica, G.; Popa, S.; Enculescu, M.; Pasuk, I.; Ionescu, A.M.; Badica, P. Dwell time influence on spark plasma-sintered MgB<sub>2</sub>. *J. Supercond. Novel Magn.* **2018**, *31*, 317–325. [[CrossRef](#)]

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