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# Optical and Thermoelectric Properties of Surface-Oxidation Sensitive Layered Zirconium Dichalcogenides $ZrS_{2-x}Se_x$ ( $x = 0, 1, 2$ ) Crystals Grown by Chemical Vapor Transport

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**Abstract:** In this work, structure, optical, and thermoelectric properties of layered  $ZrS_{2-x}Se_x$  single crystals with selenium composition of  $x = 0, 1$ , and 2 were examined. Single crystals of zirconium dichalcogenides layer compounds were grown by chemical vapor transport method using  $I_2$  as the transport agent. X-ray diffraction (XRD) and high-resolution transmission electron microscope (HRTEM) results indicated that  $ZrS_{2-x}Se_x$  ( $x = 0, 1$ , and 2) were crystallized in hexagonal  $CdI_2$  structure with one-layer trigonal (1T) stacking type. X-ray photoelectron and energy dispersive X-ray measurements revealed oxidation sensitive behavior of the chalcogenides series. Transmittance and optical absorption showed an indirect optical gap of about 1.78 eV, 1.32 eV, and 1.12 eV for the  $ZrS_{2-x}Se_x$  with  $x = 0, 1$ , and 2, respectively. From the result of thermoelectric experiment,  $ZrSe_2$  owns the highest figure-of-merit (ZT) of  $\sim 0.085$  among the surface-oxidized  $ZrS_{2-x}Se_x$  series layer crystals at 300 K. The ZT values of the  $ZrS_{2-x}Se_x$  ( $x = 0, 1$ , and 2) series also reveal increase with the increase of Se content owing to the increase of carrier concentration and mobility in the highly Se-incorporated zirconium dichalcogenides with surface states.

**Keywords:** 2D semiconductor; surface state; thermoelectric material; optical property

## 1. Introduction

Two-dimensional (2D) transition metal dichalcogenides (TMDCs) as  $MX_2$  ( $M = Mo, W, Re$  and  $X = S, Se$ ) [1–5] and III–VI layer compounds as  $NX$  ( $N = Ga, In$  and  $X = S, Se$ ) [6–8] have recently received considerable attentions on their excellent optical and electrical properties available for application in electronics and optoelectronics devices [9] because their specific characteristics of flexible, large-area, and ultra-thin, etc. These materials have a proper band gap of 1–2.5 eV (in near infrared to visible region) [10,11] that suitable for solar-energy applications. Among them, III–VI GaSe and GaS series showed feasibility for making visible light emitting devices owing to their widened and adjustable band gaps [10,11]. On the other hand,  $ZrX_2$  ( $X = S, Se$ ) are the group IVB TMDC materials that crystallized in one-layer trigonal (1T) structure of  $CdI_2$  type [12]. These materials also have band gaps ranging from 1.1 to 2 eV [13,14]. The compounds usually possess high conductivity and contains defects and surface states in its layer structure. The proper band-gap value and higher conductivity may also render  $ZrX_2$  suitable for electronic and thermoelectric device applications [15–17]. Despite their TMDC counterparts of  $MoS_2$  series undergoing an indirect-to-direct band gap transition in the monolayer form [18], the  $ZrX_2$  series TMDCs may remain in indirect band gap in all thicknesses [19]. The electronic band gaps and optical properties of the layered  $ZrX_2$  ( $X = S, Se$ ) series compounds are rarely studied up to date.

In earlier research, the first principle thermoelectric properties of zirconium dichalcogenides was reported by a number of authors [20,21]; nevertheless, experimental works on thermoelectric properties of  $ZrX_2$  ( $X = S, Se$ ) have rarely been reported. Some of the 1D nanowire and nanoscale thermoelectric materials have been claimed to have improved thermo-electronic transport property [22,23]; however, to date, experimental thermoelectric measurements on the Zr-based  $ZrS_2$ ,  $ZrSSe$ , and  $ZrSe_2$  layered TMDCs are seldom found in the literature, in spite of some thermoelectric materials [24–27] having been reported for power-generation applications. In this paper, the optical and thermoelectric properties of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) were characterized using temperature-dependent transmittance, Raman, and thermoelectric measurements from low to room temperature. Transmittance and optical absorption revealed indirect allowed transition of all series  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) layered crystals. Raman measurement of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) also sustains that the whole series 2D materials are crystallized in the 1T phase similar to the result of X-ray diffraction (XRD). The Raman mode of  $A_{1g}$  is the prominent peak in the layer plane. Because of surface oxidation, the  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series chalcogenides reveal lower resistivity. From thermoelectric measurement, the ZT value increases with the Se content is increased in the  $ZrSe_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series at room temperature.

## 2. Materials and Methods

High quality  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) bulk crystals were grown by chemical vapor transport (CVT) method using Iodine ( $I_2$ ) as the transport agent [28]. Prior to the growth, synthesis of the starting material powders of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) with stoichiometry were needed. For crystal growth, the starting material powder together with proper amount of transport agent ( $10 \text{ mg/cm}^3$ ) were cooled with liquid nitrogen, and then sealed in a vacuum of  $\sim 10^{-6}$  Torr inside a quartz ampoule. The growth temperature was set as  $900 \text{ }^\circ\text{C} \rightarrow 850 \text{ }^\circ\text{C}$  with a gradient of  $2.5 \text{ }^\circ\text{C/cm}$ . The growth time kept 280 h for growing single crystals. After the growth the crystal series essentially exhibited silver-black and shiny color on their fresh surface. When they are oxidized in ambient (about 8 days in air)  $ZrS_2$  becomes silver-yellow color,  $ZrSSe$  becomes silver orange, and  $ZrSe_2$  becomes silver red. The color change is maybe caused by free of chalcogen (S, Se) and which substitution by oxygen in the zirconium compounds. The crystal images are shown in Figure S1 in supporting information (SI) for contrast. With weakened van der Waals bonding between the layers, thin  $ZrX_2$  nanosheets can be mechanically exfoliated and transferred onto a  $SiO_2/Si$  substrate ( $8 \times 8 \times 0.3 \text{ mm}^3$ ) using Scotch tape for performing micro-Raman experiment. To avoid surface oxidation in air, the as-grown  $ZrX_2$  series crystals and exfoliated nanosheets (on  $SiO_2/Si$ ) were stored in an evacuated chamber before measurements. Prior to each of the experiments, the as-grown crystals can be exfoliated again to obtain fresh surface.

For X-ray diffraction (XRD) measurements, several small crystals from each composition of the  $ZrX_2$  series were finely ground and the X-ray powder patterns were taken and recorded by means of a slow moving radiation detector. The copper  $K\alpha$  radiation ( $1.542 \text{ \AA}$ ) was employed in the X-ray diffraction measurements, and a silicon standard was used for experimental calibration. The micro-Raman measurements of the  $ZrS_2$ ,  $ZrSSe$ , and  $ZrSe_2$  nanosheets were carried out in a RAMaker microscope spectrometer (ProTrustTech, Tainan, Taiwan) with a 532-nm solid state diode-pumped laser as the excitation source. An Olympus objective lens ( $50\times$  or  $100\times$  with long working distance) acted as the interconnection-coupled medium between the nanosheet sample, incident and reflected lights, and a charge-coupled-device (CCD) spectrometer equipped with one 1200 grooves/mm grating. The power of the 532-nm laser was controlled at 8 mW and the spot size was  $\sim 0.8 \text{ }\mu\text{m}$ . A Janis liquid helium open-circled cryostat equipped (Janis Co., Woburn, MA, USA) with Lakeshore 335 digital thermometer controller facilitated low-temperature and temperature-dependent measurements.

For transmission measurement of the thin  $ZrX_2$  ( $X = S, Se$ ) samples, a 150 W tungsten halogen lamp was utilized for the white light source. The white light was dispersed by a Photonics International (PTI) monochromator (Photonics International, MA, USA) equipped one 600 grooves/mm grating for providing monochromatic light. The incident light was chopped at 200 Hz, and phase-locked detection

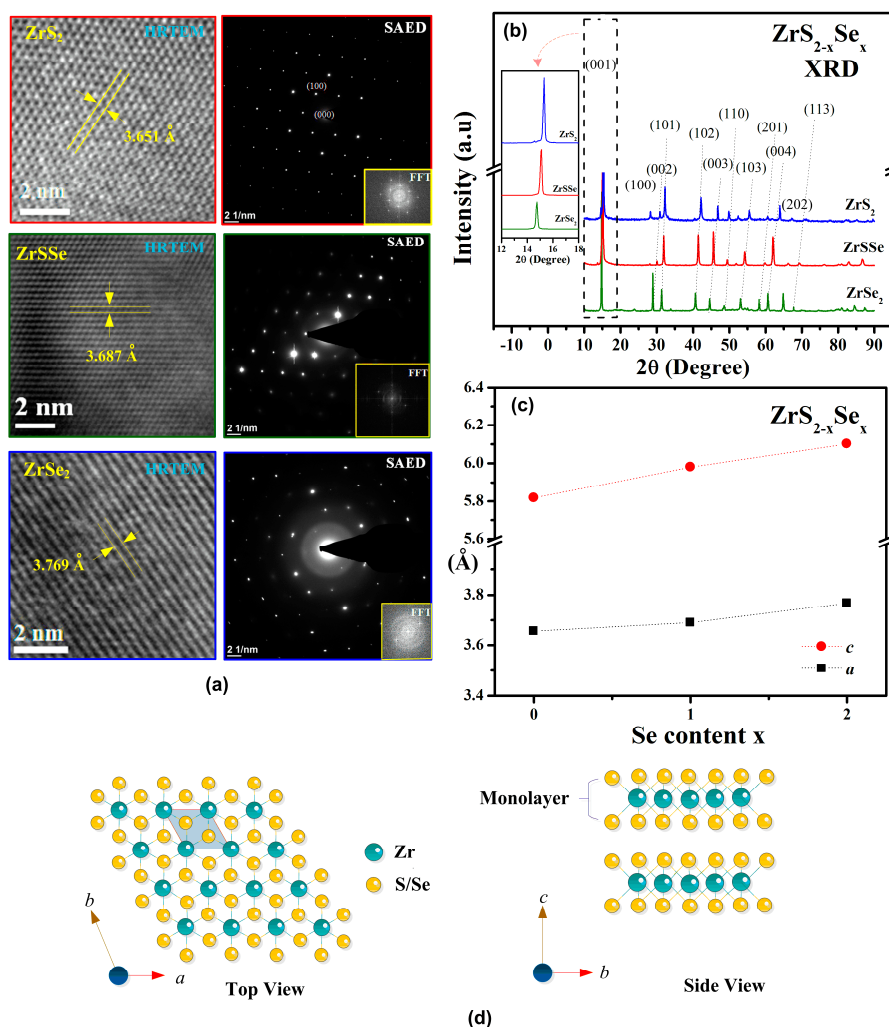
by using a lock-in amplifier was employed. The transmission light from the layer sample was detected by a Si detector ( $E > 1.1$  eV) or an InGaAs detector ( $E = 0.7$  to  $1.7$  eV) in different wavelength range.

For thermoelectric measurement, a stainless-steel plate of  $4 \times 2 \times 0.01$  cm<sup>-3</sup> covered with thin insulating tape on its top surface acted as the sample holder. The insulating tape should have much lower thermal conductivity than the measured sample, to avoid any parasitic heat loss to the substrate. Gold evaporated film on the insulating tape acted as the heater ( $R \sim 9 \Omega$ ) for the layered samples. The layered zirconium dichalcogenide crystals were prepared in rectangle shape of bar-type four-point measurement configuration (size  $\sim 0.45 \times 0.12 \times 0.01$  cm<sup>3</sup>) that available for both resistivity and thermoelectric measurements. For thermoelectric measurement, one end of the bar-shape  $ZrX_2$  sample was closely attached on the Au-film heater, and a current of  $\sim 0.4$  A was sent. When the sample temperature is balanced (after seconds), the voltage difference ( $\Delta V$ ) and temperature difference ( $\Delta T$ ) are simultaneously measured from the two central probes of the bar-shape  $ZrX_2$  sample for the estimation of Seebeck coefficient and thermal conductivity. A pair of E-type thermocouples was used for sensing the temperature difference between the two central probes in the thermoelectric-measured sample. The temperature-dependent thermoelectric measurement was done from 20 to 300 K.

### 3. Results and Discussion

#### 3.1. Structure and Composition Analysis of $ZrS_{2-x}Se_x$ ( $x = 0, 1, \text{ and } 2$ ) Crystals

Figure 1a shows high-resolution-transmission-electron-microscope (HRTEM) images of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series crystals. The obvious dotted pattern in selected-area-electron-diffraction (SAED) pictures that proves all the crystals are high-quality single crystals. From the HRTEM and SAED results, the lattice constant  $a$  can be determined to be  $3.651 \text{ \AA}$  for  $ZrS_2$ ,  $3.687 \text{ \AA}$  for  $ZrSSe$ , and  $3.769 \text{ \AA}$  for  $ZrSe_2$ , respectively. Essentially, the lattice constant  $a$  will be expanded when more Se atoms are incorporated into the  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series crystals. This result is well-matched with those obtained and calculated from XRD measurement (see Table S1). The XRD results shown in Figure 1b also indicate that  $ZrSSe$  series crystals are mainly crystallized in the 1T-CdI<sub>2</sub> layer structure with verification of the preferred  $c$  plane peak indexed as (001) [29]. The lattice constant  $c$  measured from XRD can be determined to be  $5.818 \text{ \AA}$  for  $ZrS_2$ ,  $5.980 \text{ \AA}$  for  $ZrSSe$ , and  $6.101 \text{ \AA}$  for  $ZrSe_2$ , respectively. For the lattice constant of  $a$  axis, the peak indices of (101), (102), (110), (103), (201), (202), and (113) planes also show gradually increase with the selenium content is increased in the XRD pattern in Figure 1b. This result reveals lattice dilation of  $a$  axis in the  $ZrS_{2-x}Se_x$  lattice from  $x = 0$  to  $x = 1$ . The monolayer of 1T- $ZrX_2$  consists of six hexagonal chalcogen atoms that octahedrally coordinated with one center metal atom [30]. The obtained lattice constants of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) reveal that the lattice constant  $a$  possesses the similar trend as that of  $c$ , in which the values are simultaneously increased with the Se content is increased as displayed in Figure 1c. For  $ZrX_2$ , the metal atoms are located in the centre sheet and which is sandwiched by two chalcogen layers to form X-M-X monolayer with strong covalent bond within the layer plane, and adjacent layers are connected by weak van der Waals force as displayed in Figure 1d [19,31]. All the structural analysis of HRTEM and XRD in Figure 1 confirmed that the  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series crystals are crystallized in the 1T layered structure.



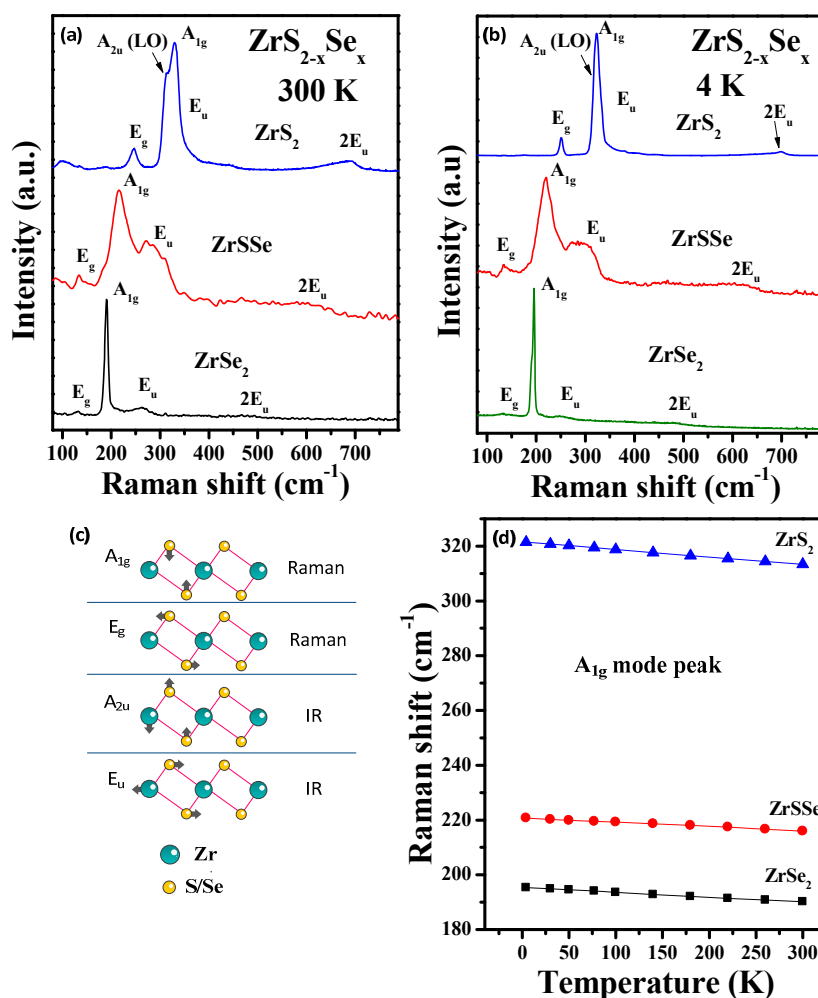
**Figure 1.** (a) High-resolution transmission electron microscope (HRTEM) results, selected-area-electron-diffraction (SAED), and FFT images of ZrS<sub>2</sub>, ZrSSe, and ZrSe<sub>2</sub>. (b) XRD pattern of ZrS<sub>2-x</sub>Se<sub>x</sub> (*x* = 0, 1, and 2). (c) Selenium doping effect on the variation of lattice constants *a* and *c*. (d) Atomic arrangements of layered ZrX<sub>2</sub> (X = S, Se) observed from top view and side view.

Energy dispersive X-ray (EDX) spectroscopy is carried out to estimate stoichiometric content of ZrS<sub>2-x</sub>Se<sub>x</sub> (*x* = 0, 1, and 2) series (Figure S2). The *x* value of chalcogen composition is lower than that of nominal composition due to the surface oxidation under ambient air condition after eight days. The ZrSe<sub>2</sub> demonstrate color change due to the surface oxidation as evident in Figure S1b,c from black shiny (fresh surface) to pink red and white color (oxidize surface). It is noticed that even with fresh surface, the ZrX<sub>2</sub> samples in Figure S2 still revealed chalcogen deficiency, the zirconium dangling bonds can easily oxidize with oxygen in the layer compounds. Thus, the surface states and defects, like S/Se vacancies or Zr interstitials, still exist in the ZrS<sub>2-x</sub>Se<sub>x</sub> to affect electrical conduction and carrier transport properties. These layered semiconductors will become degenerate and high carrier concentration, and we will evaluate the materials by temperature-dependent resistivity later.

### 3.2. Raman Spectroscopy

For micro-Raman measurement, the bulk ZrS<sub>2-x</sub>Se<sub>x</sub> (*x* = 0, 1, and 2) samples are exfoliated into thin layers (*t* < 40 nm, measured by atomic force microscopy) on a commercialized SiO<sub>2</sub>/Si substrate (thickness: SiO<sub>2</sub>~200 nm and Si~300 μm). The ZrX<sub>2</sub> thin layers were characterized from 4 to 300 K. Figure 2a shows the Raman spectra of multilayer ZrS<sub>2-x</sub>Se<sub>x</sub> (*x* = 0, 1, and 2) at room temperature,

and the low temperature Raman spectra at 4 K is also included in Figure 2b, for comparison. The temperature-dependent Raman spectra of  $\text{ZrS}_2$ ,  $\text{ZrSSe}$ , and  $\text{ZrSe}_2$  are shown in Figure S3. The in-plane ( $E_g$ ) and out-plane ( $A_{1g}$ ) Raman modes can be detected in these  $\text{ZrX}_2$  related layer crystals. At 300 K,  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1$ , and 2) exhibit much broadened peaks compared to those measured at 4 K owing to thermal effect on lattice vibration. An  $A_{2u}(\text{LO})$  peak near  $A_{1g}$  of  $\text{ZrS}_2$  is detected which may be caused by IR mode [32] or by non-harmonic effect of  $A_{1g}$  induced by acoustic phonon [33]. Figure 2c shows the representative scheme of vibrated eigen-modes of  $\text{ZrX}_2$  [32].  $\text{ZrS}_2$  reveals three peaks positioned at  $246 \pm 2$  ( $E_g$ ),  $316 \pm 2$  ( $A_{2u}$ ) and  $332 \pm 2$  ( $A_{1g}$ ), as well as IR mode  $2E_u$  at approximately  $680 \pm 2$   $\text{cm}^{-1}$  at 300 K. Raman spectrum of  $\text{ZrSSe}$  reveals strong  $A_{1g}$  mode at  $\sim 214$   $\text{cm}^{-1}$  and weak  $E_g$  mode at  $133 \pm 2$   $\text{cm}^{-1}$ . The longitudinal optical mode  $E_u$  appears at  $286 \pm 2$   $\text{cm}^{-1}$ . For  $\text{ZrSe}_2$ , there are three peaks positioned at  $126 \pm 2$  ( $E_g$ ),  $191 \pm 2$  ( $A_{1g}$ ), and  $262 \pm 2$  ( $E_u$ ) at 300 K. The Raman modes of the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1$ , and 2) series show frequency-reduction behavior with the increase of Se content. The heavier selenium atom causes the vibration peak shifted to lower wavenumber for identification of Se doping effect. Figure 2d shows the Raman shift of  $A_{1g}$  mode as the function of temperature from 4 to 300 K. The variation mode follows a linear dependence of energy reduction when the temperature is increased. This behavior may be attributed to the lattice dilation of  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1$ , and 2) at room temperature to decrease lattice vibration frequency [34]. The temperature-dependent Raman spectra of  $\text{ZrS}_{2-x}\text{Se}_x$  series are also shown in Figure S3 for detailed comparison of the modes' shift with temperature.

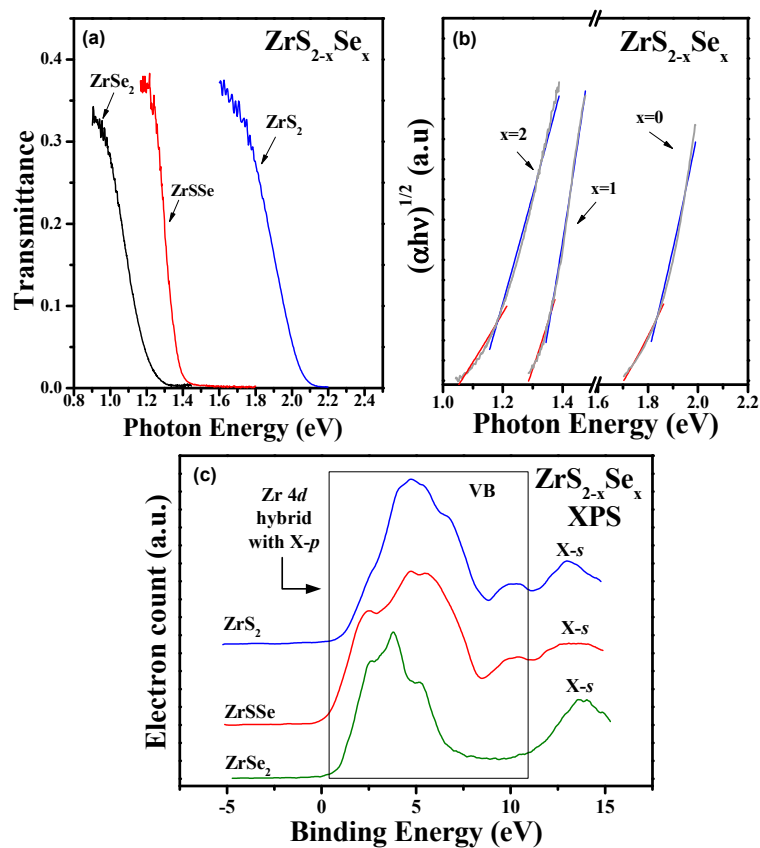


**Figure 2.** Raman spectra of  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1$ , and 2) at (a) 300 K and (b) 4 K. (c) The illustration of atomic movements of vibrational modes. (d) Raman frequency of  $A_{1g}$  mode as a function of temperature in  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1$ , and 2).



### 3.3. Optical Band Gap and Valence-Band Structure

To further characterize the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series, transmittance measurement and X-ray photoelectron spectroscopy (XPS) are carried out. The transmittance measurement is very useful to evaluate optical absorption near band-edge portion [35]. The transmittance spectrum needs to be measured at near-normal incidence, and the absorption coefficient should be determined from measured transmittance by examining the spectral dependence of reflectance [35]. The transmittance spectra of  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) crystals measured at 300 K were shown in Figure 3a. The thickness of the crystal was controlled to be  $\sim 6 \mu\text{m}$  to  $\sim 10 \mu\text{m}$  measured by atomic force microscopy. Essentially, bulk  $\text{ZrX}_2$  series crystals were mechanically exfoliated to  $\mu\text{m}$  scale to obtain fresh surface using Scotch tape. It is noticed that even a little oxide ( $\text{ZrO}_2$ ) still existed on the surface of  $\text{ZrS}_2$ ; the band gap of  $\text{ZrO}_2$  ( $\sim 5 \text{ eV}$ ) is usually larger than  $\text{ZrS}_2$ , similar to that of  $\text{Ga}_2\text{O}_3$  [36] >  $\text{Ga}_2\text{S}_3$  [37] and  $\text{In}_2\text{O}_3$  [38] >  $\text{In}_2\text{S}_3$  [39]. There is no light-shielding effect from the surface oxide on  $\text{ZrX}_2$  when doing transmittance. The related band-edge value of  $\text{ZrX}_2$  could be obtained by analyzing the photon ( $E_g$ ) and phonon ( $E_p$ ) energies expressed by  $\alpha h\nu = (h\nu - E_g + E_p)^2 / [\exp(E_p/kT) - 1] + (h\nu - E_g - E_p)^2 / [1 - \exp(-E_p/kT)]$  for assisting an indirect transition of absorption and emission of a phonon in the optical process. The plot of  $(\alpha h\nu)^{1/2}$  versus  $h\nu$  for the analysis of indirect gaps are shown in Figure 3b. The indirect band edge  $E_g$  of each crystal is 1.78 eV for  $\text{ZrS}_2$ , 1.32 eV for  $\text{ZrSSe}$ , and 1.12 eV for  $\text{ZrSe}_2$  at 300 K. These values are in agreement with previous result reported for  $\text{ZrS}_2$  and  $\text{ZrSSe}$  [40]. The obtained phonon energies from the linear fits of  $(\alpha h\nu)^{1/2}$  versus  $h\nu$  [41] are  $E_p = 40 \pm 10 \text{ meV}$  for all the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series. The variation of indirect band edge indicated the doping effect of selenium content in pure  $\text{ZrS}_2$  would cause the band-edge position shifting to lower energy with increasing Se content. This situation resembled the gradual lattice constant change with respect to the increasing of selenium composition obtained by XRD and HRTEM measurements. The XPS spectra in Figure 3c show the main valence-band structure of  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$  with fresh surface) is coming from Zr 4d hybrid with chalcogen p orbital. According to the electron configuration, transition metal Zr d electrons used to have bonding with chalcogen S/Se, which forms  $\text{Zr}^{4+}$  ( $d^0s^0$ ) configuration connected to  $2 \cdot (\text{S/Se})^{2-}$  ( $s^2p^6$ ) [42]. For octahedral coordination, the d orbitals will split into two sets of energies due to the trigonal distortion and ligand field. The higher energy set comes from  $d_z^2$  symmetry and lower energy set comes from  $d_{xz}, d_{yz}, d_{xy}$  symmetry [43]. For group IVB TMDC materials, especially for the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series, the  $d_z^2$  set is vacant and forms the lowest conduction band. The complete set of XPS spectra were also measured and displayed in Figure S4 for comparison. The energies of Zr 3d, S 2p, and Se 3d orbitals show redshift as the Se content is increased in the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series. It also verifies a little oxygen-incorporation effect in the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series layer crystals.

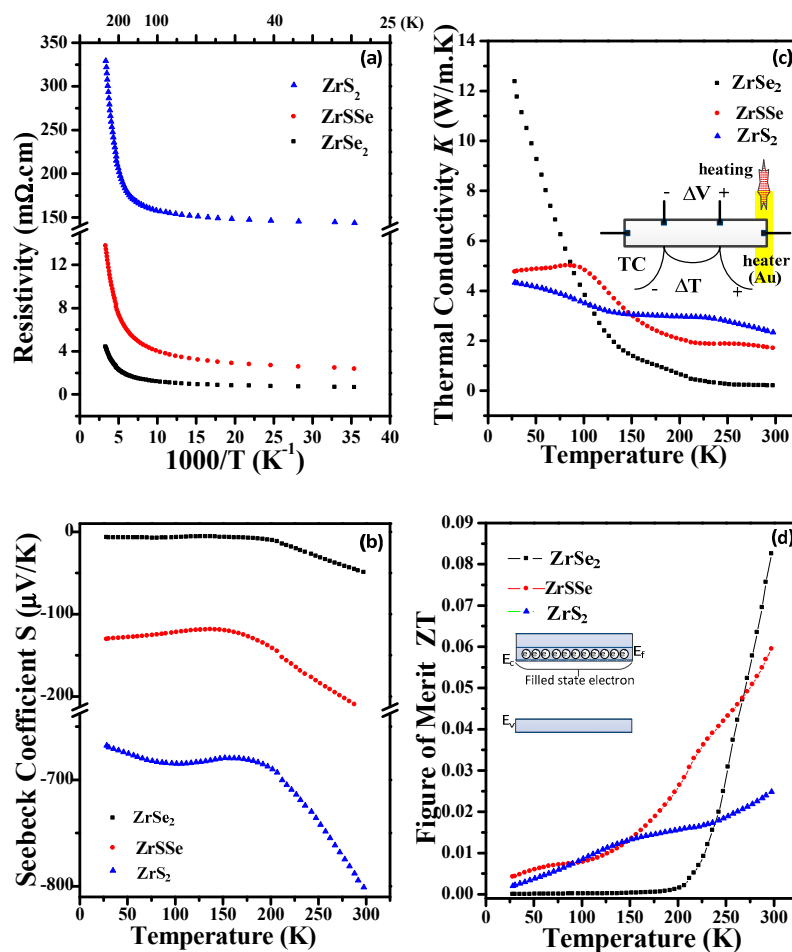


**Figure 3.** (a) Transmittance spectra of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) at  $300$  K. (b) Plot of  $(\alpha h\nu)^{1/2}$  versus  $h\nu$  for the analysis of indirect gaps of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) at  $300$  K. (c) Valence-band XPS spectra of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) compounds.

#### 3.4. Transport and Thermoelectric Properties

For the characterization of electrical and thermoelectric properties of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) series layer crystals, temperature-dependent resistivity and thermoelectric measurements from  $20$  to  $300$  K were carried out. Since many researchers have claimed that  $ZrX_2$  ( $X = S, Se$ ) series displayed good thermoelectric properties, we measured the experimental result and compared with previous theoretical results [44,45] herein. The layered samples of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) crystals with dimension about  $0.45 \times 0.12 \times 0.01$  cm<sup>3</sup> were prepared. Because the existence of defects, surface states, surface oxide (i.e.,  $ZrX_{2-x}O_x$ ,  $x = 0$  to  $2$ ), etc., room-temperature resistivity and Hall-effect measurements of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) were also done to identify their electrical property. The results are listed in Table S2, for comparison. The carrier type of all the  $ZrX_2$  series was n type, and carrier concentrations were determined to be  $4.05 \times 10^{17}$  cm<sup>-3</sup> for  $ZrS_2$ ,  $1.6 \times 10^{18}$  cm<sup>-3</sup> for  $ZrSSe$ , and  $2.9 \times 10^{18}$  cm<sup>-3</sup> for  $ZrSe_2$ , respectively. The high carrier concentration also renders lower resistivity of  $0.25$   $\Omega$ -cm for  $ZrS_2$ ,  $0.0211$   $\Omega$ -cm for  $ZrSSe$ , and  $0.0058$   $\Omega$ -cm for  $ZrSe_2$ , respectively. Figure 4a shows temperature dependent resistivity of  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) series as a function of reciprocal temperature. The series shows metallic carrier-conduction behavior due to the increasing resistivity as the temperature increased. The lower resistivity of the  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) series may come from heavily doping from defects, like chalcogen vacancies or native impurities, i.e., surface oxidation states in air (or from environment water vapor) [46], which cause the material to behave like a degenerate semiconductor, as the band scheme shown in the inset of Figure 4d. Dissimilar to the other 2D materials, like  $MoS_2$  and  $MoSe_2$ ,  $ZrS_{2-x}Se_x$  ( $x = 0, 1,$  and  $2$ ) series exhibits relatively low resistivity owing to the surface states. The sign of Seebeck coefficient ( $S = -\Delta V/\Delta T$ ) indicates all three  $ZrX_2$  compounds are n-type semiconductor as shown in Figure 4b. The carrier type is similar to that determined by Hall-effect

measurement. The resistivity values of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) in Table S2 reveal reduction with the increase of Se content owing to the increased mobility and carrier concentration obtained by Hall measurement. This property may result in reduction of Seebeck coefficient of the high Se compound as displayed in Figure 4b. The Seebeck coefficient of  $ZrSe_2$  is about  $-50 \mu\text{V/K}$  at 300 K. The magnitude of  $S$  value is close to a previous  $ZrS_{3-3x}Se_{3x}$  compound [47]. The efficiency of a thermoelectric material is determined by a dimensionless figure of merit ( $ZT$ ) defined as  $ZT = S^2 \cdot T / (\rho \cdot \kappa)$ , where  $\rho$ ,  $T$ , and  $\kappa$ , are electrical resistivity, absolute temperature, and thermal conductivity. Thermal conductivity  $\kappa$  can be calculated from the heating power and temperature difference  $\Delta T$  from the sample geometry as  $P = \kappa \cdot (w \cdot t / l) \cdot \Delta T$ , where  $w$  is width,  $t$  is thickness, and  $l$  is length, respectively. The  $ZT$  is an index of efficient thermoelectric material, but  $ZT$  value is inversely proportional to thermal conductivity  $\kappa$  [48,49]. Figure 4d shows the  $ZT$  values of the  $ZrS_{2-x}Se_x$  series from 20 to 300 K. The  $ZT$  values of all samples increase when the temperature is increased, while thermal conductivity  $\kappa$  decreases when the temperature is raised (see Figure 4c). At 300 K, the highest  $ZT$  value of 0.085 was achieved by  $ZrSe_2$ , followed by  $ZrSSe$  of  $ZT = 0.06$  and  $ZrS_2$  of  $ZT = 0.02$ , respectively. When the Se content increases, the mobility and electron concentration are also increased to promote thermoelectric performance in the  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series. The carrier concentration of  $ZrSe_2$  is close to  $\sim 10^{19} \text{ cm}^{-3}$  in Table S2. The concentration value was claimed to be an optimum carrier density for achieving optimal  $ZT$  in the thermoelectric materials [50].  $ZrSe_2$  is therefore shows the best thermoelectric performance in the  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) series layered TMDCs.



**Figure 4.** (a) Resistivity as function of reciprocal temperature. (b) Seebeck coefficient of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ) as function of temperature. (c) Thermal conductivity as function of temperature of  $ZrS_{2-x}Se_x$  ( $x = 0, 1, \text{ and } 2$ ). The inset shows the experimental configuration of thermoelectric measurement. (d) The dependences of  $ZT$  values (figure of merit) versus temperature.



#### 4. Conclusions

In conclusion,  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) crystals were successfully grown by CVT method. The structural, optical, and thermoelectric properties of the compounds were characterized. The HRTEM image, SAED, and XRD results confirmed the 1T-CdI<sub>2</sub> structure. The  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series revealed an indirect band gap positioned at 1.78 eV, 1.32 eV, and 1.12 eV, respectively. The energy gap varies accordingly with Se composition change, which indicates the similar electronic and band structures of the layered  $\text{ZrX}_2$  ( $X = \text{S, Se}$ ). A few layers of Raman spectra of  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) were observed from 4 to 300 K. The frequencies of vibration modes decreased with the increase of Se content due to the heavier atomic weight of Se. All the layered  $\text{ZrS}_{2-x}\text{Se}_x$  reveal chalcogen deficiency and existing native oxide ( $\text{ZrX}_{2-x}\text{O}_x$ ) and surface states on their surface to render lower resistivity compared to the other 2D wide-gap TMDCs. Temperature-dependent resistivity shows degenerate semiconductor behavior of all  $\text{ZrS}_{2-x}\text{Se}_x$  layer compounds, with  $\text{ZrSe}_2$  possessing the lowest resistivity. Along with Hall-effect measurement,  $\text{ZrSe}_2$  also shows higher mobility and higher electron concentration to promote carrier transport and thermoelectric performance in the layered  $\text{ZrS}_{2-x}\text{Se}_x$ . At 300 K, the ZT is increased, and  $\kappa$  is decreased, with the increase of the Se content. An experimental value of  $\text{ZT} \approx 0.085$  is obtained for  $\text{ZrSe}_2$  owing its enhanced mobility and optimal carrier concentration in the  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ) series layered crystals.

**Supplementary Materials:** The following are available online at <http://www.mdpi.com/2073-4352/10/4/327/s1>, Figure S1: Crystal morphology of *c*-plane layered  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ). Figure S2: Energy Dispersive X-ray (EDX) analysis of  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ). Figure S3: Temperature-dependent Raman spectra of  $\text{ZrS}_2$ ,  $\text{ZrSe}_2$ , and  $\text{ZrSSe}$ . Figure S4: The XPS spectra of (a) Zr 3d<sub>3/2</sub> and Zr 3d<sub>5/2</sub>, (b) S 2p<sub>3/2</sub> and S 2p<sub>1/2</sub>, and (c) Se 3d<sub>5/2</sub> and Se 3d<sub>3/2</sub> in  $\text{ZrS}_{2-x}\text{Se}_x$  ( $x = 0, 1, \text{ and } 2$ ). Table S1: The comparison of lattice constants of  $\text{ZrS}_{2-x}\text{Se}_x$  obtained by XRD and TEM measurements. Table S2: Resistivity and Hall measurement results at 300 K.

**Author Contributions:** T.M.H. and C.-H.H. synthesized the compounds and grew the crystals. C.-H.H. conceives the research idea and designs & implements the thermoelectric-measurement system. T.M.H. did the optical and electrical measurements under the supervision of C.-H.H.; T.M.H. and C.-H.H. wrote the paper. All authors have read and agreed to the published version of the manuscript.

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