



Article Effect of Graphene on Ice Polymorph

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Abstract: Recently, ice with stacking disorder structure, consisting of random sequences of cubic ice (I_c) and hexagonal ice (I_h) layers, was reported to be more stable than pure I_h/I_c . Due to a much lower free energy barrier of heterogeneous nucleation, in practice, the freezing process of water is controlled by heterogeneous nucleation triggered by an external medium. Therefore, we carry out molecular dynamic simulations to explore how ice polymorphism depends on the lattice structure of the crystalline substrates on which the ice is grown, focusing on the primary source of atmospheric aerosols, carbon materials. It turns out that, during the nucleation on graphene, we find I_h is the dominant polymorph. This can be attributed to structural similarities between graphene and basal face of I_h . Our results also suggest that the substrate only affects the polymorph of ice close to the graphene surface, with the preference for I_h diminishing as the ice layer grows.

Keywords: heterogeneous nucleation; ice polymorph; stacking disorder; phase selectivity

1. Introduction

Ice nucleation is important in the physical environment and biological systems [1–3]. Thus, it is important to accurately estimate the nucleation rate of ice, which is highly related to polymorph of ice crystalline [4,5]. For a long time, hexagonal ice (I_h) was assumed to be the most stable ice phase at atmospheric pressure or below. However, this understanding has been questioned by a large number of reports by computer simulation [5–8] and experiment [9–13] in recent years. In the report of Lupi et al. [5] the stacking-disordered [14] critical ice crystallites are about 14 kJ/mol of crystallite more stable than hexagonal ice crystallites (at 230 K). In these simulations, homogeneous ice nucleation was considered, but in practice, it is almost impossible to eliminate the influence of impurities or external boundaries on ice nucleation [15,16]. Due to a much lower free energy barrier of heterogeneous nucleation. This raises the question to what extent heterogeneous ice nucleation influences the recently discovered preference for stacking-disordered ice.

As the main component of atmospheric aerosols, carbon surfaces can greatly promote heterogeneous ice nucleation [17–19]. The crystallization temperature of ice on the graphite surface has been found to be about 12 K higher than the temperature for homogeneous ice nucleation [20,21]. This stimulated both experimental and molecular dynamics (MD)



Citation: Zheng, C.; Lu, H.; Xu, Q.; Liu, T.; Patil, A.; Wu, J.; de Vries, R.; Zuilhof, H.; Zhang, Z. Effect of Graphene on Ice Polymorph. *Crystals* 2021, *11*, 1134. https://doi.org/ 10.3390/cryst11091134

Academic Editors: Haihui Zhang, Xunhui Xiong, Huihui Xiong and Yaohui Qu

Received: 18 June 2021 Accepted: 15 July 2021 Published: 18 September 2021

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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). simulations investigation of heterogeneous ice nucleation on graphene/graphite and other carbon surfaces [20,22–24].

However, the issue of how nucleation on carbon surfaces influences the preference for stacking-disordered ice has so far not been addressed. To do so, we conducted MD simulations for a series of heterogeneous ice nucleation/growth processes on different carbon surfaces and found that the carbon surfaces have a strong influence on the local ice polymorphs.

2. Methods

Modeling. As shown in Figure 1, ice nucleation was studied on four types of atomic flat carbon surfaces with different lattice structures: *Graphene* (composed by six ring carbon atom), *Oblique-Haeckelite* (*O-Haeckelite*, composed by 5-6-7 ring carbon atom), *Rectangular-Haeckelite* (*R-Haeckelite*, composed by 5–7 ring carbon atom), and *Random* (all carbon atoms distributed randomly, refer to Figure S1) [25]. The reason behind these specific surfaces is that we aimed to clarify the effect of the hexagonal structure of graphene on the polymorph of ice, compared to the two non-hexagonal *Haeckelite* structures and an amorphous 2D carbon substrate. Analogous homogeneous ice nucleation (*Homo*) simulations without any substrate were also carried out as control. The size of the 3D periodic simulation boxes (which contain 15,029 water molecules) are $15.0 \times 14.8 \times 15.0$ nm³ for the *Graphene* system, $16.3 \times 15.3 \times 15.0$ nm³ for the *O-Haeckelite* system, $17.0 \times 15.5 \times 15.0$ nm³ for the *R-Haeckelite* system, $15.0 \times 14.9 \times 15.0$ nm³ for the *Random* system, and $15.0 \times 15.0 \times 15.0$ nm³ for the *Homo* system, respectively.



Figure 1. (a) Example of a simulation box. (b–d) show top view (part) of substrates *Graphene*, *O*-*Haeckelite*, and *R*-*Haeckelite*, respectively. The top view (part) of substrates *Random* is shown in Figure S1. Carbon atoms in substrate are portrayed as gray spheres. The water molecules are showed in red dots.

Simulation Details. All MD simulations were performed using large-scale atomic /molecular massively parallel simulator (LAMMPS, Version 5Jun19) package [26]. The coarse-grained mono-atomic mW water model was applied in all simulations [27]. The interactions between mW water molecules consist of a sum of pairwise interactions, while the hydrogen bonding between water molecules is represented by three-body contributions, which can nucleate liquid water to ice homogeneously and heterogeneously at a certain supercooling without any advanced sampling method. The mW model correctly displays the anomalies and structures of liquid water, ices, and amorphous solid water and the transitions between them [28]. The equations of water motion were integrated with the

velocity Verlet algorithm with a time step of 5 fs. All the simulations were conducted in the NVT ensemble. The temperature in the simulation systems was controlled by a Nosé-Hoover thermostat. The same as in previous reports [20], ice nucleation was studied through a cooling ramp of cooling rates of 1 K/ns, which was chosen after a systematic test of the cooling rates by the authors. To calculate the freezing efficiency [20], $\Delta T_f = T_f - T_f^{homo}$, 40 independent trajectories were performed for each system. The interactions between water molecules and carbon atoms are taken from the previous report by Lupi' et al. [20], in which water–carbon interaction parameters are: $\sigma_{WC} = 0.32$ nm and $\varepsilon_{WC} = 0.13$ kcal/mol, to reproduce the experimental water contact angle of water on a graphene surface (namely 86°) [29,30]. All the carbon atoms in the substrates were fixed in all the MD simulations. Water molecules with I_h/I_c structure were identified using the Chill + algorithm proposed by Nguyen et al. [31], which is available in the OVITO package [32].

3. Results and Discussion

To investigate the effect of the substrate lattice structure on the polymorph of ice, MD simulations were employed to study the ice formation process on different substrates: *Graphene*, *Oblique-Haeckelite* (*O-Haeckelite*), *Rectangular-Haeckelite* (*R-Haeckelite*), and a *Random* carbonaceous substrate. As shown in Figure 2, consistent with previous reports [20,22,33,34], after an induction period (during which, the small water clusters with ice structure formed and then disappeared quickly, due to the thermal fluctuation of system), stable ice nuclei (larger than critical size R_c of ice) formed at the water-substrate interface for the systems of *Graphene*, *O-Haeckelite*, *R-Haeckelite*, and *Random*, which should be due to the much lower heterogeneous nucleation barrier. For these heterogeneous nucleation systems, especially for the nucleation stage, the ice crystals exhibit a single-crystal-like structure with barely any grain boundary, which is due to the 1-dimensional structure match (in the direction perpendicular to the substrate surface) between a flat substrate and a flat crystalline face of ice. While for the *Homo* system, the new-formed ice exhibits a polycrystalline structure.



Figure 2. Lateral view of the ice formation procedure from top to bottom: (i) ice nucleation; (ii) ice growth; (iii) water freezing completely. Liquid water is represented by red dots. Water molecules in ice crystallites are represented by colored ball-stick model (I_h: yellow; I_c, cyan). Carbon atoms in substrate are colored gray.

Intriguingly, for the system of *Graphene*, the I_h structure occupies an absolute majority in the ice nuclei during the stage of nucleation. With the growth of ice (after about 5–6 layers of I_h formed), the I_c turns up, which is consistent with the previous report that the I_c grows on I_h embryos to form a more stable stacking disorder structure [5]. When almost all the liquid water freezes to ice, the percentage of I_c is not much different from that of I_h. While for the systems of *O*-*Haeckelite*, *R*-*Haeckelite*, *Random*, and *Homo*, in all stages of ice formation, the number of water molecules in I_c is always comparable with the number of water molecules in I_h.

To reveal the phase change process of ice during heterogeneous nucleation processes, the molecular numbers of I_c and I_h as a function of system temperature (namely simulation time) were extracted from the ice formation MD trajectories (shown Figure 3). The snapshots in Figure 2 share the same trajectories with the results of Figure 3 for each system. As shown in Figure 3a (*Graphene* system), with the decrease in temperature, the nucleation process was observed in the I_h ice before the steady growth of I_c . In the growth stage, the I_c exhibited a lager growth rate than I_h , leading to an equivalent molecular number of I_h and I_c at the stage of water freezing completely. While for each of the systems in Figure 3b–e, the I_c and I_h grow simultaneously as the system temperature decreases, indicating that these systems have no selectivity to ice polymorph.



Figure 3. Molecular numbers of water molecules in I_c and I_h during a cooling ramp for the systems of (a) *Graphene*, (b) *O-Haeckelite*, (c) *R-Haeckelite*, (d) *Random*, and (e) *Homo*. The dashed vertical lines indicate position of snapshots of the ice nucleation, ice growth, and water freezing completely in Figure 2. The blue borders are for panel (e).

To quantitatively reveal the polymorph of ice formed on different substrates, the cubicity [6] (i.e., the proportion of I_c) of the new-formed ices was calculated from the 40 independent MD trajectories for each system (refer to Table 1 and Table S1). As shown in Table 1, the cubicity of ice in the *Homo* system is $52.5 \pm 1.3\%$, $53.9 \pm 1.6\%$, and $58.2 \pm 0.7\%$ for the stage of ice nucleation, ice growth, and water freezing completely, respectively,

which is consistent with previous reports (about 55%) [5–7,35]. For ice formed on Graphene, the cubicity is only $28.4 \pm 2.5\%$ at the nucleation stage, which is significantly lower than that of the Homo system. The results suggest that the graphene substrate has a preferential selectivity to I_h over I_c , which we attribute to the fact that I_h (basal face) and graphene share the same hexagonal structure and have a similar lattice structure. The distance between the center of two hexagon rings on graphene surface is 2.46 Å [24,36], and the distance of water molecules in the basal face of I_h is 2.76 Å [37,38]. According to previous reports [24,39,40], in the water–graphene system, the center of a hexagon formed by carbon atoms corresponds to the adsorption energy minima positions of water molecules. Due to the similarity of lattice structure between the basal plane of $I_{\rm h}$ and graphene, slight adjustment of the position of interfacial water molecules can match the lattice structure of basal plane of I_h . The calculated mismatch [41] between the substrate and the ice is 10.9%. The ice-nucleating protein, with a mismatch of 10% to ice, was found to be able to dramatically promote the nucleation of ice [42]. In the report of Bi et al. [24], it was found that the first ice layer on graphene substrate was mainly composed of I_h, while the first ice layer on an amorphous graphene substrate (similar to the O-Haeckelite substrate in this paper) was somewhat messy. The study of Bi et al. also suggests that the similarity between the graphene and the basal face of I_h results in the selective promotion of the formation of I_h over I_c . Moreover, in the study of ice formation on AgI [43,44], it is also found that the substrate lattice structure can alter the polymorph of ice: I_h mainly observed on the hexagonal β -AgI, while the Ic mainly found on cubic γ -AgI. These reports are in good agreement to our MD simulation results, confirming that the formation of I_h can be selectively promoted by graphene surface.

Table 1. Average cubicity of each simulation system. From top to bottom are the stages of ice nucleation, ice growth, and water freezing completely, respectively.

Substrate	Graphene	O-Haeckelite	R-Haeckelite	Random	Ното
Cubicity/%	$\begin{array}{c} 28.4 \pm 2.5 \\ 34.2 \pm 2.9 \\ 53.4 \pm 1.8 \end{array}$	$\begin{array}{c} 41.2 \pm 1.9 \\ 45.2 \pm 4.5 \\ 55.3 \pm 1.5 \end{array}$	$\begin{array}{c} 42.9 \pm 2.3 \\ 44.0 \pm 3.7 \\ 56.5 \pm 1.5 \end{array}$	$\begin{array}{c} 42.6 \pm 2.8 \\ 45.5 \pm 3.9 \\ 55.0 \pm 1.7 \end{array}$	$\begin{array}{c} 52.5 \pm 1.3 \\ 53.9 \pm 1.6 \\ 58.2 \pm 0.7 \end{array}$

To further investigate the effect of substrate lattice structure on heterogeneous ice nucleation, freezing efficiency of the substrates, based on the method of Lupi et al. [20], were calculated (shown in Figure 4). It follows that the calculated freezing efficiencies for *Graphene*, *O*-Haeckelite, *R*-Haeckelite, and *Random* substrates are 12.7 ± 0.6 K, 11.5 ± 1.2 K, 12.5 ± 1.0 K, and 11.1 ± 1.3 K, respectively, which are consistent with the previous reports that the crystallization temperature of ice on the graphite surface is 12 ± 3 K higher than the temperature of homogeneous ice nucleation [20,21]. The freezing efficiency of these different substrates has no significant difference, indicating that, although the substrates exhibit different polymorph selectivity of ice, the heterogeneous nucleation promotion effects of the substrates on ice are similar. Moreover, it should be pointed out that the heterogeneous nucleation promotion effect changes with the system temperature [24]. To calculate the freezing efficiency, the system temperature is steadily decreased. The freezing processes are driven by a very high supercooling, which can reduce the heterogeneous nucleation promotion effect. For instance, the heterogeneous ice nucleation rates on different substrates exhibit significant differences only when the system temperature increased to as high as 235 K [24]. Thus, for a systematic study, e.g., heterogeneous nucleation barrier and heterogeneous nucleation rate, it is necessary to gain more accurate results by taking into account the influence of temperature. Another thing that needs to be specified is that, in our results, the calculated homogeneous nucleation temperature of ice is 1.6 K higher than that in the previous report [20], which can be attributed to the difference in the number of water molecules between this work and the previous report (refer to Figure S2). According to classic nucleation theory (CNT), the greater the number

of water molecules in the system, the greater the nucleation probability of ice formation is, which could lead to this 1.6 K discrepancy.



Figure 4. Freezing efficiency of substrates in Graphene, O-Haeckelite, R-Haeckelite, and Random systems.

4. Conclusions

MD simulations of heterogeneous ice formation on a series of carbon surfaces show the significant impact of the substrate lattice structure on the polymorph of ice that is formed, but only during the nucleation stage. Specifically, graphene substrates have a preferential selectivity toward the formation of I_h over I_c during the nucleation stage, due to the similarity between the lattice structure of graphene and I_h . After the nucleation stage the cubicity of new-formed ice increases up to about 53%, due to the higher stability of stacking disordered ice. When subsequently the water freezes completely, the cubicity of ice ranges from 53% to 58%, in good agreement with homogeneous ice formation and previously reported results. This study enhances our understanding of the surface selectivity mechanism on certain ice polymorphs and provides a unique perspective in the field of crystal growth.

Supplementary Materials: The following are available online at https://www.mdpi.com/article/ 10.3390/cryst11091134/s1. Table S1: Cubicity of the nucleation stage in each simulation trajectory. Figure S1: Surface morphology of 2-dimensional carbon substrate with all the carbon atoms randomly distributed, the Random system. Figure S2: Freezing temperature in *Homo-5241* (5241 water molecules) and *Homo-15029* (15029 water molecules) systems.

Author Contributions: Conceptualization, C.Z., J.W. and Z.Z.; methodology, C.Z., T.L. and Z.Z.; software, C.Z. and Q.X.; investigation, C.Z. and H.L.; data curation, C.Z., H.L., A.P. and Z.Z.; writing—original draft preparation, C.Z., H.L. and Z.Z.; writing—review and editing, C.Z., H.L., R.d.V., H.Z. and Z.Z.; visualization, C.Z., H.L. and Z.Z.; supervision, Z.Z.; project administration, Z.Z.; funding acquisition, J.W. and Z.Z. C.Z. and H.L. contributed equally to this manuscript. All authors have read and agreed to the published version of the manuscript.

Funding: This research was funded by the National Natural Science Foundation of China (Grant Nos. 11904300, 11772278 and 11502221), the Jiangxi Provincial Outstanding Young Talents Program (Grant No. 20192BCBL23029), the Fundamental Research Funds for the Central Universities (Xiamen University: Grant Nos. 20720180014, 20720180018 and 20720180066).

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: All data generated or analyzed during this study are included in this published article and its additional files.

Acknowledgments: Y. Yu and Z. Xu from Information and Network Center of Xiamen University are acknowledged for the help with the high-performance computer clusters.

Conflicts of Interest: The authors declare no conflict of interest.

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