

Article

Electronic Structure and Optical Properties of $\text{Cu}_2\text{ZnSnS}_4$ under Stress Effect

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Abstract: By using the pseudopotential plane-wave method of first principles based on density functional theory, the band structure, density of states and optical properties of $\text{Cu}_2\text{ZnSnS}_4$ under isotropic stress are calculated and analyzed. The results show that $\text{Cu}_2\text{ZnSnS}_4$ is a direct band gap semiconductor under isotropic stress, the lattice is tetragonal, and the band gap of $\text{Cu}_2\text{ZnSnS}_4$ is 0.16 eV at 0 GPa. Stretching the lattice causes the bottom of the conduction band of $\text{Cu}_2\text{ZnSnS}_4$ to move toward lower energies, while the top of the valence band remains unchanged and the band gap gradually narrows. Squeezing the lattice causes the bottom of the conduction band to move toward the high-energy direction, while the top of the valence band moves downward toward the low-energy direction, and the $\text{Cu}_2\text{ZnSnS}_4$ band gap becomes larger. The static permittivity, absorption coefficient, reflectivity, refractive index, electrical conductivity, and energy loss function all decrease when the lattice is stretched, and the above optical parameters increase when the lattice is compressed. When the lattice is stretched, the optical characteristic peaks such as the dielectric function shift to the lower-energy direction, while the optical characteristic peak position shifts to the higher-energy direction when the lattice is compressed.

Keywords: $\text{Cu}_2\text{ZnSnS}_4$; stress; electronic structure; optical properties; first principles



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1. Introduction

With the increasing scarcity of non-renewable energy sources such as fossils, the energy crisis and environmental pollution problems are becoming more and more prominent. The development of new energy sources, such as solar energy and wind energy, is an effective means to solve the energy crisis and environmental pollution problems [1]. Among many new energy sources, solar energy has received great attention due to its advantages of greenness, environmental protection, and high efficiency [2–6]. Solar cells are the main means of utilizing solar energy. Solar cells include silicon-based solar cells, perovskite solar cells, organic solar cells, dye-sensitized cells, compound thin-film solar cells, etc. [7–11]. Compound thin-film solar cells are effective devices for utilizing solar energy due to their small size, light weight, and flexibility. Compound thin-film solar cells mainly include [12–14] CdTe cells, GaAs cells, $\text{Cu}(\text{In-Ga})(\text{Se,S})_2$ cells, and so on. Among them, $\text{Cu}(\text{In-Ga})(\text{Se,S})_2$ cells have high battery conversion efficiency due to their spike-like conduction band valence and they do not damage the short-circuit current [15] and are widely used. $\text{Cu}(\text{In-Ga})(\text{Se,S})_2$ cells contain rare metals In and Ga, which are limited in large-scale commercial applications. Therefore, seeking environmentally friendly solar cell materials is one of the strategies to solve the difficulty of the large-scale application of $\text{Cu}(\text{In-Ga})(\text{Se,S})_2$ cells. The band gap of the semiconductor material $\text{Cu}_2\text{ZnSnS}_4$ (CZTS) is 1.4~1.6 eV, which is strongly matched with the required band gap of the solar cell absorber material [16]. The physical and chemical properties of CZTS are very close to those of $\text{Cu}(\text{In-Ga})(\text{Se,S})_2$. The theoretical conversion efficiency of CZTS as an absorber layer of solar cells is 32.2% [17], and CZTS is rich in various elements in the crust, as well as being non-toxic and non-polluting. Therefore, CZTS is considered as a high-efficiency, low-cost, and environmentally friendly solar cell material.

At present, the experimentally reported value of the conversion efficiency of CZTS solar cells reaches 12.6% [18], but there is still a large gap between this and the theoretical calculation value, and the conversion efficiency of the cell has large room for improvement. The main factors affecting the conversion efficiency of CZTS solar cells are the low open-circuit voltage V_{oc} and low filling factor of the cell [19,20]. Since the ionic radii of Cu and Zn are very close, substitution defects of Cu and Zn are easily formed in the CZTS system. Anti-occupancy defects with Cu and Sn and Zn anti-occupancy defects [21] will introduce energy level defects to different degrees and inhibit the open-circuit voltage V_{oc} of the battery. It was found [22] that the introduction of metal ions in CZTS can play a role in the passivation of defects; for example, the introduction of a small amount of Ge can suppress intrinsic deep-level Sn_{Zn} defects. We note that the doping of monovalent metals such as silver can change the defects of Cu_2ZnSnS_4 , thus adjusting the band gap. Due to the high formation energy of Ag_{Zn} defects, when the Ag doping concentration is high, the defect concentration of the CZTS system is low, which is conducive to increasing the carrier concentration. Affected by the electronic configuration of elements, the atomic radii of Ag atoms and Cu atoms differ greatly, which causes the bending coefficient of the energy band to change. Low-concentration doping causes the energy band to change slowly, while high-concentration doping causes the energy band to change obviously [23].

In addition, the introduction of stress into the material can play the role of changing the lattice constant, regulating the energy band structure, and changing the optical properties. The 4-metal alloy of Cu_2ZnSnS_4 is derived from the binary alloy of ZnS [24]; due to its good band gap matching and strong light absorption, it is an ideal solar cell absorption layer material. However, Cu_{Zn} and Sn_{Zn} defect energy levels exist when Cu_2ZnSnS_4 alloy films are actually prepared, which affects the photoelectric properties of the materials. Therefore, we control the band gap and photoelectric properties of Cu_2ZnSnS_4 alloy films by applying stress. Stressing is an effective means to regulate the heterojunction carrier dynamics [25]. As a solar cell absorbing layer material, CZTS needs to form a solar cell system together with other materials, so it is necessary to consider the transport of electrons and holes at the heterojunction interface. Adjusting the stress can change the K point energy level of the material. When the K point energy level is coupled with other energy levels, it can provide an intermediate-state energy level for electron and hole transfer, making the transfer of electrons and holes between layers easier, and improving the utilization of photogenerated carriers. In addition, adjusting the stress can change the orbital occupancy of electrons, thereby changing the photoelectric characteristics of the film. At the same time, adjusting the stress can also change the phase transition temperature of the film [26]. Therefore, it can be predicted that the introduction of stress is an effective means to adjust the photoelectric properties and phase transition temperature of CZTS. Based on the above considerations, we study the effect of stress on the electronic structure and optical properties of CZTS.

There have been extensive reports on the photoelectric properties of materials by introducing stress [27–30]. However, there are few reports on the effect of stress on the electronic structure and optical properties of CZTS. In this paper, by using the pseudopotential plane-wave method of first principles based on density functional theory, the band structure, density of states, and optical properties of Cu_2ZnSnS_4 under isotropic stress are calculated and analyzed.

2. Methods and Calculations

The calculation model selected in this paper is Cu_2ZnSnS_4 with a kesterite structure, the space group is $I\bar{4}(No.82)$, the lattice constants are $a = b = 0.54628$ nm, $c = 1.0864$ nm, and the crystal plane angle is $\alpha = \beta = \gamma = 90^\circ$ [31]. The unit cell structure of Cu_2ZnSnS_4 is shown in Figure 1.

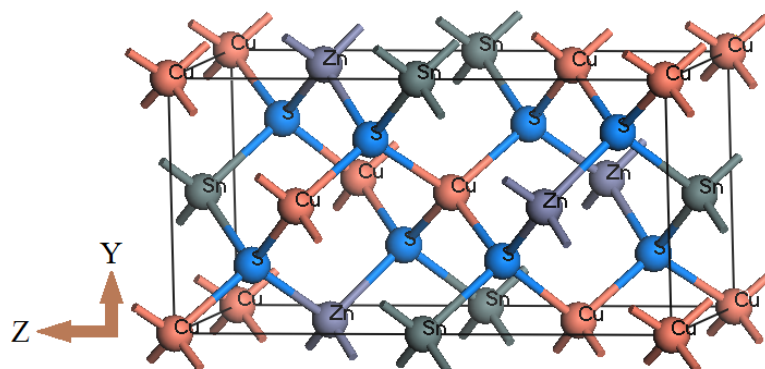


Figure 1. $\text{Cu}_2\text{ZnSnS}_4$ unit cell structure.

The first-principles-based pseudopotential method was used for the calculations; all calculations were performed using the Cambridge Serial Total Energy Package (CASTEP) [32] in Materials Studio software, and processed using the Perdew–Burke–Ernzerh(PBE) [33] functional of the generalized gradient approximation (GGA) for the exchange correlation energy between electrons. An ultrasoft pseudopotential [34] was used to deal with the interaction between ionic solids and electrons. The truncation energy was set to 400 eV, the self-consistent convergence accuracy was 1.0×10^{-7} eV/atom, and the integral of the Brillouin zone was divided by $5 \times 5 \times 2$ of Monkhorst–Pack.

3. Results and Discussion

3.1. Geometry Optimization

Table 1 shows the lattice constants of $\text{Cu}_2\text{ZnSnS}_4$ with a kesterite structure obtained after applying isotropic tensile stress and compressive stress optimization, where “–” represents tensile stress. It can be seen from Table 1 that with the increase in isotropic tensile stress, the lattice constant of CZTS increases, the volume of the unit cell increases, and the lattice of CZTS is tetragonal. With the increase in isotropic compressive stress, the lattice constant of CZTS decreases, the unit cell volume decreases, and the lattice still maintains the tetragonal system. The results calculated by us are in good agreement with those of Liu [35], which shows that the applied stress has a significant effect on the lattice constant and cell volume of $\text{Cu}_2\text{ZnSnS}_4$.

For comparison, the theoretical calculations [36] and experimental values [37] of lattice constant and cell volume are added in Table 1. It can be seen from Table 1 that when no stress is applied, the difference between the a and b values and theoretical calculation is 0.0%, and the difference between them and the laboratory values is 0.7%. The difference between the c value and theoretical calculation is 0.2%, and the difference between the c value and experimental value is 0.8%. The unit cell volume is 0.2% different from the theoretical calculation and 2.4% different from the experimental value. The results in Table 1 show that our calculation results are in good agreement with the calculation and experimental values of other research groups.

Table 1. $\text{Cu}_2\text{ZnSnS}_4$ lattice constants of applied tensile and compressive stresses.

Stress/GPa	$a/\text{Å}$	$b/\text{Å}$	$c/\text{Å}$	$V/\text{Å}^3$
–6	5.688	5.688	11.377	368.043
–4	5.602	5.602	11.202	351.220
–2	5.528	5.528	11.061	338.034
0	5.469	5.469	10.944	327.377
Theoretical [36]	5.469	5.469	10.921	326.647
Experimental [37]	5.427	5.427	10.854	319.676
6	5.338	5.338	10.673	304.123
12	5.240	5.240	10.480	287.732
20	5.149	5.149	10.242	271.568

3.2. Electronic Structure

3.2.1. Band Structure

Figure 2 shows the $\text{Cu}_2\text{ZnSnS}_4$ band structure obtained by applying tensile stress. It can be seen from the figure that CZTS is a typical direct band gap semiconductor, and the minimum band gap value is obtained at the highly symmetrical G point position. The band gap of CZTS is 0.16 eV when no stress is applied (0 GPa), which is consistent with the research results of Zhao [38].

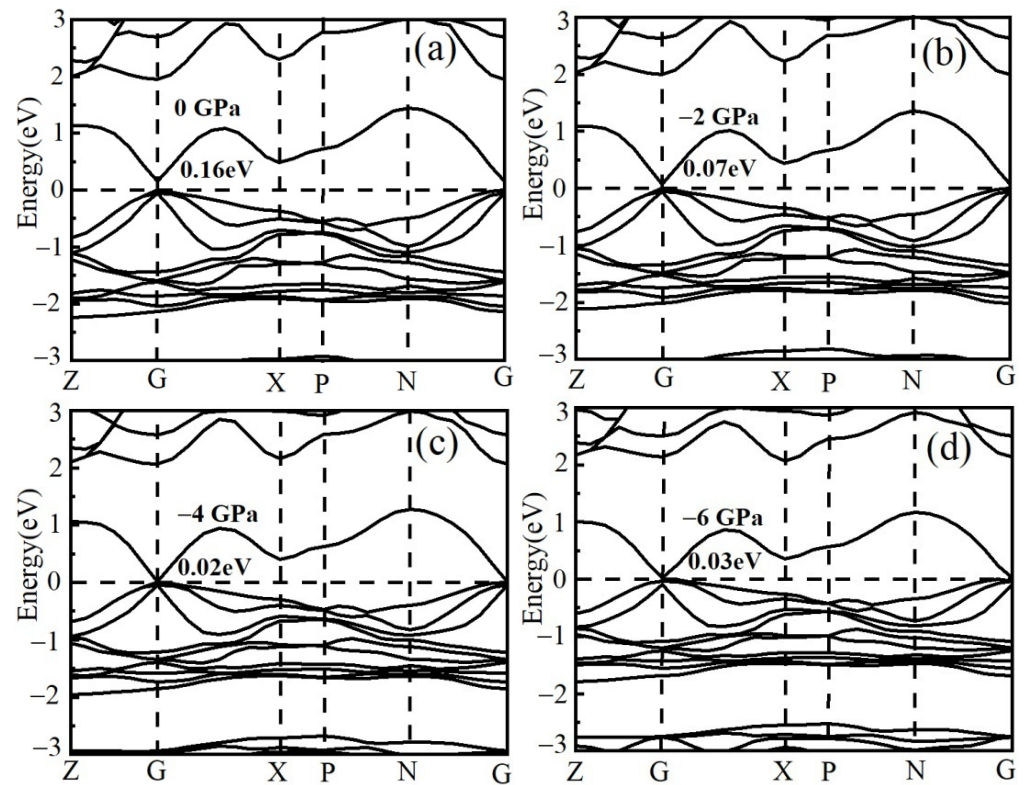


Figure 2. Band structure of $\text{Cu}_2\text{ZnSnS}_4$ under tensile stress for (a) 0 GPa band structure; (b) –2 GPa band structure; (c) –4 GPa band structure; (d) –6 GPa band structure.

With the increase in tensile stress, the bottom of the conduction band moves to the lower-energy direction, while the top of the valence band remains unchanged; the band gap decreases with the increase in tensile stress, and the band gap reaches a minimum value of 0.02 eV when the tensile stress is –4 GPa. After this, the band gap widens with the increase in tensile stress. This is because the bottom of the conduction band of CZTS moves toward the high-energy direction after the tensile stress continues to increase, while the top of the valence band moves down toward the low-energy direction, and the band gap widens.

Figure 3 shows the band structure of $\text{Cu}_2\text{ZnSnS}_4$ obtained by applying compressive stress. It can be seen from Figure 3 that CZTS is still suitable as a direct band gap semiconductor; the minimum band gap value is still obtained at the highly symmetrical G point position.

With the increase in incompressive stress, the bottom of the conduction band of CZTS moves towards high energy and the top of the valence band remains unchanged. The band gap of CZTS widens with the increase in incompressive stress. When the applied compressive stress is 20 GPa, the band gap of CZTS reaches 0.71 eV.

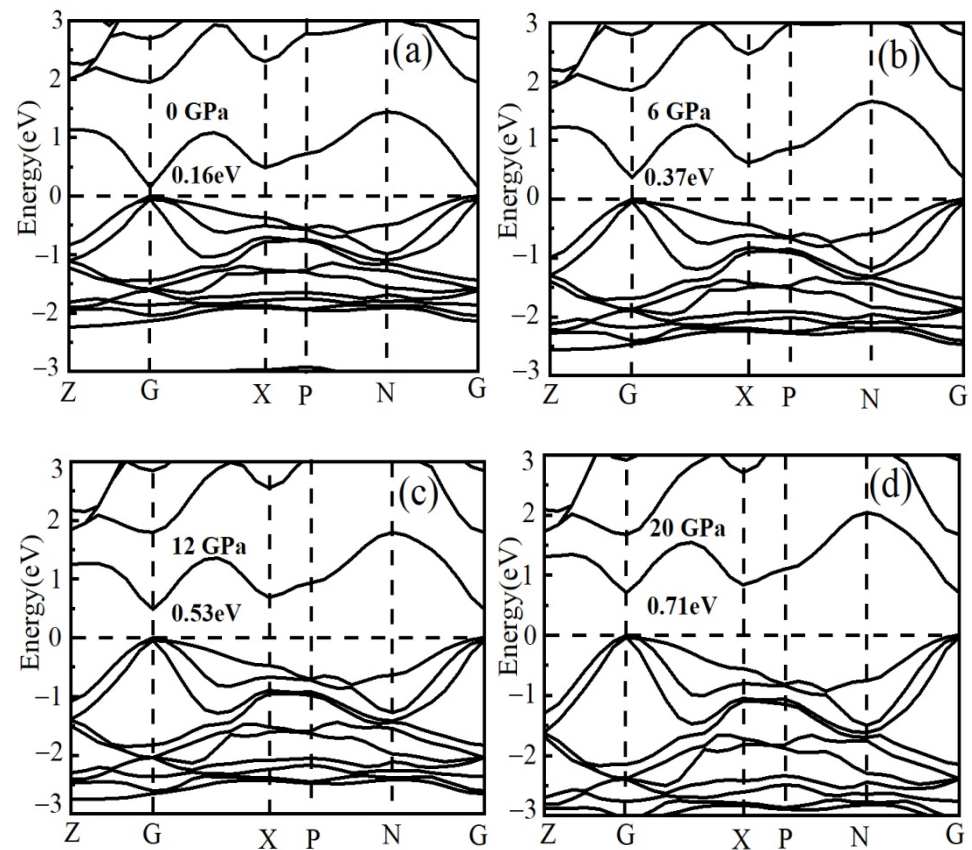


Figure 3. Band structure of $\text{Cu}_2\text{ZnSnS}_4$ under compressive stress for (a) 0 GPa band structure; (b) 6 GPa band structure; (c) 12 GPa band structure; and (d) 20 GPa band structure.

3.2.2. Density of Electronic States

The total density of states and the density of states of each atomic partial wave under the action of tensile stress are shown in Figure 4. The main contributions to the CZTS density of states are the Cu 3d configuration, Zn 3d configuration, Sn 5s and 5p configuration, and S 3s and 3p configuration, and the 3p configuration of Cu and Zn also has a small contribution.

It can be seen from Figure 4 that the lower valence band region of $-14.5\sim-12.3$ eV is mainly contributed by the 3s state of S and a small amount of 5p state electrons of Sn. The mid-valence band region of $-8.4\sim-5.5$ eV is mainly contributed by the 3d state of Zn, the 5s state of Sn, and a small amount of the 3p state of S and 4s state of Zn. The upper valence band region of $-5.5\sim 0$ eV is mainly contributed by the 3d state of Cu, a small amount of the Sn 5p state, and the 3p state of S. The conduction band part is mainly contributed by the 5s and 5p states of Sn and a small amount of the 3p state of S.

With the increase of tensile stress, the peak position of the density of states in the valence band of CZTS shifts to the higher-energy direction. The -13 eV peak is shifted to the high-energy band by 0.45 eV, the -6.88 eV peak is shifted to the high-energy band by 0.27 eV, the -3.69 eV peak is shifted to the high-energy band by 0.8 eV, and the -1.73 eV peak is shifted to the high-energy band by 0.42 eV. The peak position of the conduction band is shifted to the lower-energy direction, and the shift amount is essentially the same at around 0.13 eV.

The $\text{Cu}_2\text{ZnSnS}_4$ density of states and the partial wave density of states under the action of compressive stress are shown in Figure 5.

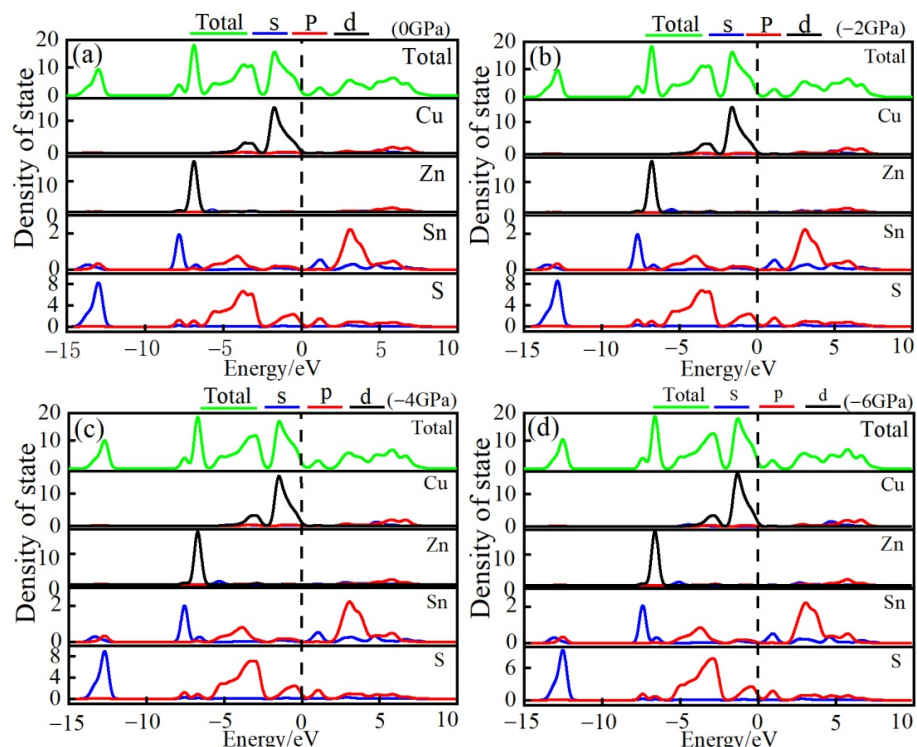


Figure 4. $\text{Cu}_2\text{ZnSnS}_4$ density of states under tensile stress for (a) 0 GPa density of states and fractional density of states; (b) -2 GPa density of states and fractional density of states; (c) -4 GPa density of states and fractional density of states; (d) -6 GPa density of states and partial wave density of states.

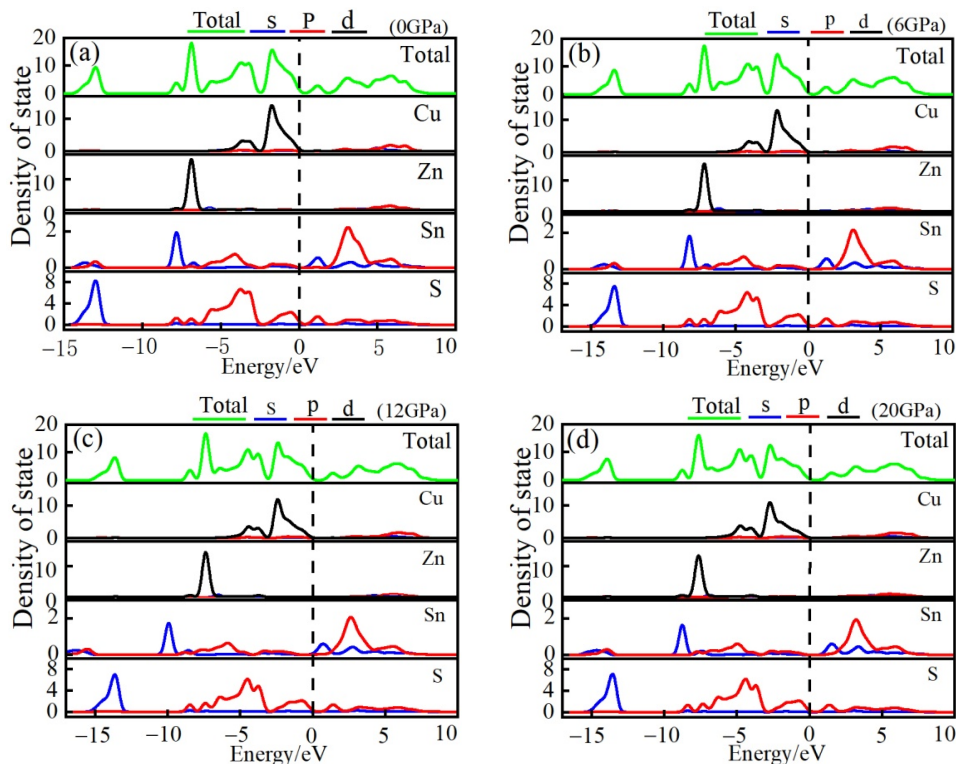


Figure 5. $\text{Cu}_2\text{ZnSnS}_4$ density of states under compressive stress for (a) 0 GPa density of states and fractional density of states; (b) 6 GPa density of states and fractional density of states; (c) 12 GPa density of states and fractional density of states; (d) 20 GPa density of states and partial wave density of states.

It can be seen from Figure 5 that the electronic configurations contributing to the valence and conduction bands of CZTS are essentially the same as those under tensile stress. The peak positions of the density of states in the valence band of CZTS shifted to the lower-energy direction with the increase in compressive stress, with an average offset of 0.64 eV, and the peak positions of the density of states in the conduction band shifted to the high-energy direction with an average offset of 0.17 eV.

We believe that when the lattice is subjected to tensile stress, the interatomic distance increases, the coupling degree of valence electrons decreases, the electrostatic repulsion decreases, the attractive force between valence electrons near the Fermi level and electrons in the conduction band increases, and some valence electrons break free. The tendency for nuclear confinement to occupy higher energy levels becomes more pronounced with increasing tensile stress.

Since valence electrons occupy higher energy levels, the energy required for their transition to the conduction band decreases and the band gap narrows. On the contrary, when the lattice is subjected to compressive stress, the atomic spacing decreases, the coupling degree of valence electrons increases, and the electrostatic repulsion increases, which hinders valence electrons from occupying high-energy states and pushes valence electrons to move to lower-energy orbitals. Therefore, the energy required for the transition of valence electrons to the conduction band increases, and the band gap becomes wider.

On the other hand, when the material is stressed, the Coulomb force changes due to the change in atomic spacing. These changes lead to an increase or decrease in exciton binding energy [39], which affects the dielectric shielding effect and ultimately affects the electron transport, making the band gap wider or narrower.

3.3. Optical Properties

3.3.1. Complex Dielectric Function

The dielectric function $\epsilon(\omega) = \epsilon_1(\omega) + i\epsilon_2(\omega)$ is a complex number. The imaginary part ϵ_2 indicates that the polarization of the molecules inside the material cannot keep up with the change in the external electric field; it represents the loss. The real part ϵ_1 indicates the ability of the material to bind charges. The dielectric function reflects the band structure of the solid and its spectral information [40], and the dielectric peak of the dielectric function is mainly determined by the band structure and density of states. Figure 6 is a graph showing the change in the CZTS real part ϵ_1 and imaginary part ϵ_2 with photon energy under stress.

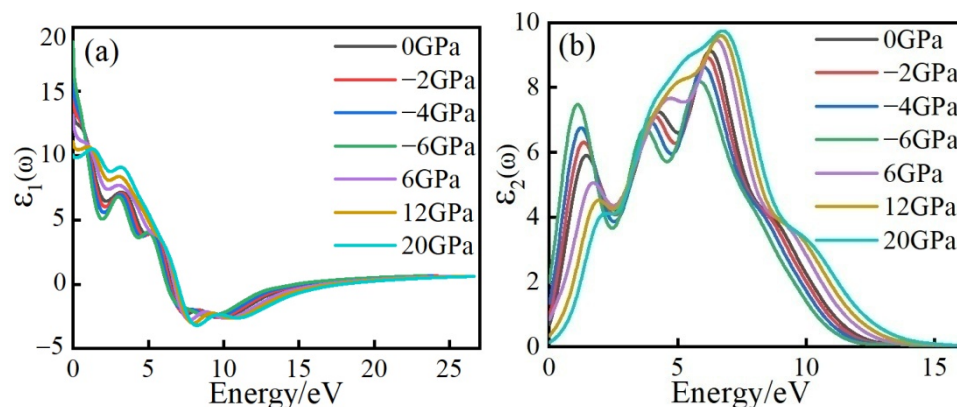


Figure 6. The complex dielectric function of Cu₂ZnSnS₄ under stress for (a) the real part of the complex dielectric function; (b) the imaginary part of the complex dielectric function.

It can be seen from Figure 6a that when the tensile stress is 0, -2, -4, and -6 GPa, the corresponding static dielectric constants are 13.74, 14.98, 16.48, and 18.88. When the compressive stress is 6, 12, and 20 GPa, the corresponding static dielectric constants are 13.09, 11.09, and 10.04. It can be seen from the figure that the static permittivity of CZTS

increases when the lattice is subjected to tensile stress, while the static permittivity of CZTS decreases when it is subjected to compressive stress. At 0 GPa, the imaginary part of the complex dielectric function of CZTS has four distinct characteristic peaks at 1.44, 4.23, 6.26, and 8.58 eV. Combining the density of states in Figure 4a, it can be seen that the dielectric peak at 1.44 eV has mainly transitioned from Cu 3d orbital electrons to Sn 5s orbital electrons, and the dielectric peak at 4.23 eV has mainly transitioned from Cu 3d orbital electrons to Sn 5p orbital electrons, at 6.26 eV. The dielectric peak at 8.58 eV mainly transitions from Cu 3d orbital electrons to Sn 5p orbital electrons, and the dielectric peak at 8.58 eV mainly transitions from Zn 3d orbital electrons to Sn 5s orbital electrons. When the energy is lower than 3.78 eV, the dielectric peak shifts to the high-energy direction with the increase in stress, and the peak decreases with the increase in stress. When the energy is greater than 3.78 eV, the dielectric peak also shifts to the high-energy direction with the increase in stress, but, at this time, the peak increases with the increase in stress. As the stress increases, the CZTS band gap increases, so the dielectric peak shifts toward higher energies. The above conclusions are consistent with the results of the literature [41]; this is in agreement with the relation between the band gap energy and the dielectric response, which implies that the larger band gap energy results in a small dielectric constant.

3.3.2. Absorption and Reflection Spectra

The absorption spectrum of $\text{Cu}_2\text{ZnSnS}_4$ is shown in Figure 7a. It can be seen from the figure that the absorption spectrum of CZTS is mainly divided into three parts: the visible light region of 0.5–4.8 eV, the ultraviolet light absorption region of 4.8–14.2 eV, and the absorption range of more than 14.2 eV in the high-energy absorption area. In the range of 0.5–9.9 eV, the absorption coefficient gradually increased with the increase in incident light energy, and the light absorption decreased sharply when the energy was greater than 9.9 eV, while CZTS almost no longer absorbed the spectrum when the energy was greater than 16.19 eV. It can be seen from the figure that the absorption coefficient of CZTS in the visible light band is greater than 10^4 cm^{-1} [42], and the light absorption is good.

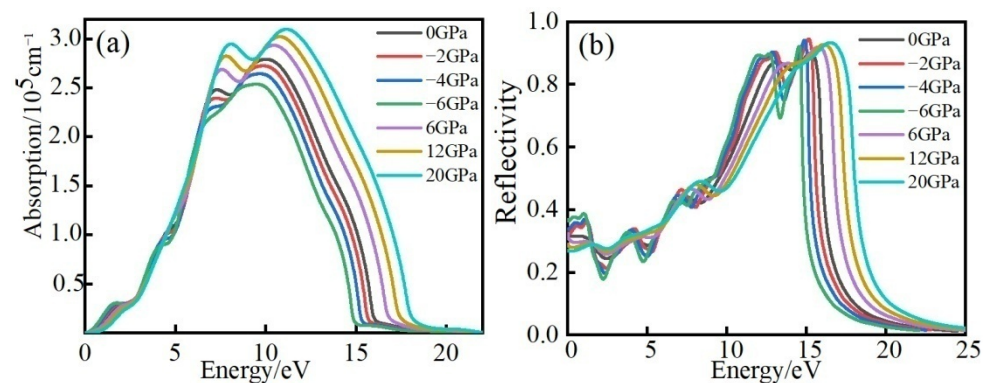


Figure 7. Absorption and reflection spectra for (a) the absorption spectrum and (b) the reflection spectrum.

With the increase in compressive stress, the absorption peak of CZTS shifted to the high-energy direction, and the peak value gradually increased. With the increase in tensile stress, the absorption peak of CZTS shifted to the lower-energy direction, and the peak value gradually decreased. The results obtained in this paper are very consistent with the research results of Kahlaoui [43], and the phenomena could be attributed to the decrease in the electronic transition energy with the strain, which results in more photons being absorbed.

In this work, it is believed that when the lattice is subjected to tensile stress, the atomic spacing becomes larger, and the light wave of the same energy is "short-wave" relative to the lattice, the lattice scattering effect is enhanced, and the light absorption is weakened. In contrast, when the lattice is subjected to compressive stress, the atomic spacing decreases,

and the light wave with the same energy is "long-wave" relative to the lattice, the lattice scattering effect is weakened, and the light absorption is enhanced.

Figure 7b shows the reflectance spectrum of CZTS. The reflectance in the visible light region of 0.5~4.8 eV is lower than 30%, indicating that CZTS has good light absorption in the visible light region. With the increase in incident light energy, the reflectivity gradually increased and reached a peak value of 90% near 15.30 eV, and high reflectivity appeared in the ultraviolet and high-energy regions. As the compressive stress increases, the reflection peak shifts towards higher energy. However, as the tensile stress increases, the reflection peak shifts toward the lower-energy direction.

3.3.3. Complex Refractive Index

Figure 8 shows the $\text{Cu}_2\text{ZnSnS}_4$ complex refractive index; (a) is the refractive index, and (b) is the extinction coefficient.

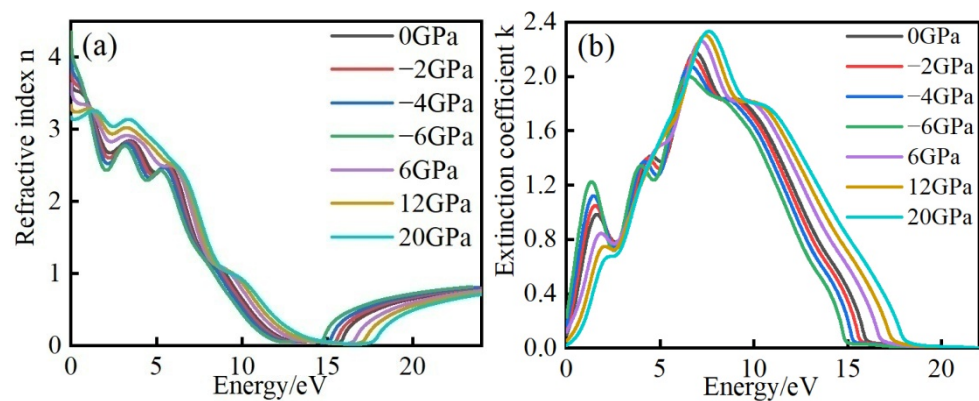


Figure 8. Complex refractive index for (a) the refractive index and (b) the extinction coefficient.

It can be seen from Figure 8 that when the energy is greater than 16 eV, the refractive index is essentially constant, and the extinction coefficient is around 0, indicating that the absorption of CZTS is weak at high frequencies, which is consistent with the conclusion of the absorption spectrum. It can also be seen from the figure that the refractive index decreases with increasing tensile stress and increases with increasing compressive stress. This is mainly because the unit cell volume becomes larger when tensile stress is applied, the density becomes worse, and the refractive index decreases. When compressive stress is applied, the volume of the unit cell decreases, the density of CZTS becomes better, and therefore the refractive index increases.

3.3.4. Complex Conductivity

It can be seen from Figure 9a that when the energy is less than 0.5 eV and greater than 12.6 eV, the real part of the complex conductivity is around 0—that is, there is almost no dissipation. It can be seen from the figure that the peak of the real part of the conductivity appears in the energy range of 1.5~8.8 eV, and the peak of the imaginary part appears in the energy range of 1.5~11.35 eV, which corresponds to the absorption spectrum. The real part of the conductivity reaches the maximum value near 6.4 eV. Combined with the analysis of the density of states map, it can be seen that these interband transition sources are related to the transition of Cu 3d orbital electrons to Sn 5p orbital electrons. With the increase in tensile stress, the peak value of the real part of conductivity shifts to the direction of low energy, and with the increase in compressive stress, the peak value of the real part of conductivity shifts to the direction of high energy.

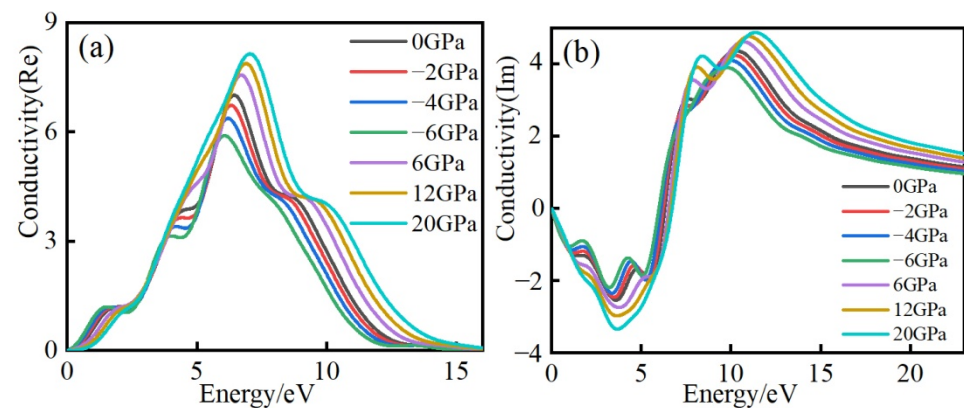


Figure 9. Complex conductivity for (a) real part and (b) imaginary part.

3.3.5. The Energy Loss Function

The energy loss function describes the energy loss of electrons when passing through a homogeneous medium. It can be seen from Figure 10 that the energy loss of $\text{Cu}_2\text{ZnSnS}_4$ reaches a maximum value of 57.22 near 15.89 eV at 0 GPa. When the stress is -6 GPa, the maximum energy loss is 149.23, and the minimum energy loss is 47.37 when the stress is 20 GPa.

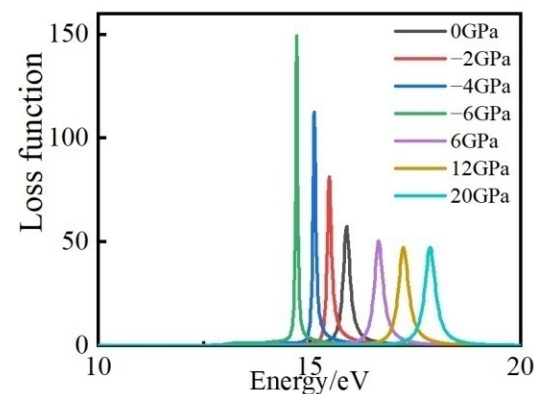


Figure 10. The energy loss function.

With the increase in tensile stress, the peak position of the energy loss function shifts to the lower-energy direction, and the peak value increases significantly with the increase in tensile stress. With the increase in compressive stress, the peak position of the energy loss function shifts to the high-energy direction, and the peak value decreases slowly with the increase in compressive stress. This is mainly because the CZTS band gap changes significantly when the lattice is stretched, while the CZTS band gap changes slowly when the lattice is compressed. Therefore, applying stress can tune the peak position and peak value of the CZTS energy loss function.

4. Conclusions

By using the pseudopotential plane-wave method of first principles based on density functional theory, the band structure, density of states, and optical properties of $\text{Cu}_2\text{ZnSnS}_4$ under isotropic stress are calculated and analyzed. The CZTS band gap is reduced by stretching the lattice under isotropic stress, and increased by compressing the lattice. The band gap is 0.16 eV when no stress is applied. When the material is stressed, the atomic spacing becomes larger or smaller, which changes the Coulomb force. These changes lead to an increase or decrease in exciton binding energy, which ultimately affects the electron transport and makes the band gap wider or narrower. The valence band of CZTS is mainly contributed by the 3d electrons of Cu, and the conduction band is mainly contributed by the 5s electrons of Sn. When the stress changes from -6 GPa to 20 GPa, CZTS always

presents a direct band gap semiconductor at the high-symmetry point and maintains a tetragonal crystal system. At -6 GPa, the maximum static dielectric constant is 18.88, and the dielectric peak shifts to the high-energy direction with the increase in stress. The absorption coefficient of CZTS in the visible light band is greater than 10^4 cm^{-1} , and the reflectivity is lower than 30% in the energy range of 0.5–4.8 eV, showing good light absorption characteristics and low reflectivity. With the increase in compressive stress, the absorption peak of CZTS shifts towards the direction of high energy, and the peak value gradually increases. When the energy is greater than 16 eV, the refractive index is essentially constant, and the extinction coefficient is approximately 0. The maximum value of energy loss function is 57.22 near 15.89 eV, and the peak position of the energy loss function shifts to the direction of high energy with the increase in compressive stress, while the peak value decreases slowly with the increase in compressive stress. In summary, the band gap, optical absorption, dielectric constant, and other properties of CZTS can be adjusted by applying stress; this work provides a strong reference for the experimental preparation of high-quality CZTS films and can help in the development of efficient CZTS solar cells.

Author Contributions: Contributions, X.Y.; experimental design, X.Y.; simulation calculation, X.Y.; writing—review and editing, X.Q.; model building, X.Q.; data analysis, W.Y.; overall planning, W.Y.; review and revision of the thesis, C.Z.; drawing, C.Z.; literature review, D.Z.; software, D.Z.; analysis, B.G.; design of the question, B.G. All authors have read and agreed to the published version of the manuscript.

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References

1. Yue, Q.; Liu, W.; Zhu, X. N-Type Molecular Photovoltaic Materials: Design Strategies and Device Applications. *J. Am. Chem. Soc.* **2020**, *142*, 11613–11628. [[CrossRef](#)] [[PubMed](#)]
2. Sun, J.M.; Zhao, E.; Liang, J.; Li, H.; Zhao, S.; Wang, G.; Gu, X.; Tang, B.Z. Diradical-Featured Organic Small-Molecule Photothermal Material with High-Spin State in Dimers for Ultra-Broadband Solar Energy Harvesting. *Adv. Mater.* **2022**, *34*, 2108048. [[CrossRef](#)] [[PubMed](#)]
3. Harijan, D.; Gupta, S.; Ben, S.K.; Srivastava, A.; Singh, J.; Chandra, V. High photocatalytic efficiency of α -Fe₂O₃-ZnO composite using solar energy for methylene blue degradation. *Phys. B Condens. Matter* **2022**, *627*, 413567. [[CrossRef](#)]
4. Hu, Y.-H.; Li, M.-J.; Zhou, Y.-P.; Xi, H.; Hung, T.-C. Multi-physics investigation of a GaAs solar cell based PV-TE hybrid system with a nanostructured front surface. *Sol. Energy* **2021**, *224*, 102–111. [[CrossRef](#)]
5. Celline, A.C.; Subagja, A.Y.; Suryaningsih, S.; Aprilia, A.; Safriani, L. Synthesis of TiO₂-rGO Nanocomposite and its Application as Photoanode of Dye-Sensitized Solar Cell (DSSC). *Mater. Sci. Forum* **2021**, *1028*, 151–156. [[CrossRef](#)]
6. Guo, W.H.; Zhu, Y.H.; Zhang, M.; Du, J.; Cen, Y.; Liu, S.; He, Y.; Zhong, H.; Wang, X.; Shi, J. The Dion-Jacobson perovskite CsSbCl₄: A promising Pb-free solar-cell absorber with optimal bandgap 1.4 eV, strong optical absorption 10^5 cm^{-1} , and large power-conversion efficiency above 20%. *J. Mater. Chem. A* **2021**, *9*, 16436–16446. [[CrossRef](#)]
7. Das, B.; Hossain, S.M.; Nandi, A.; Samanta, D.; Pramanick, A.K.; Chapa, S.O.M.; Ray, M. Spectral conversion by silicon nanocrystal dispersed gel glass: Efficiency enhancement of silicon solar cell. *J. Phys. D Appl. Phys.* **2021**, *55*, 025106. [[CrossRef](#)]
8. Yan, F.R.; Yang, P.Z.; Li, J.B.; Guo, Q.; Zhang, Q.; Zhang, J.; Duan, Y.; Duan, J.; Tang, Q. Healing soft interface for stable and high-efficiency all-inorganic CsPbI₂Br₂ perovskite solar cells enabled by S-benzylisothiourea hydrochloride. *Chem. Eng. J.* **2021**, *430*, 132781. [[CrossRef](#)]

9. Bi, P.; Zhang, S.; Chen, Z.; Xu, Y.; Cui, Y.; Zhang, T.; Ren, J.; Qin, J.; Hong, L.; Hao, X.; et al. Reduced non-radiative charge recombination enables organic photovoltaic cell approaching 19% efficiency. *Joule* **2021**, *5*, 2408–2419. [[CrossRef](#)]
10. Khataee, A.; Azevedo, J.; Dias, P.; Ivanou, D.; Dražević, E.; Bientien, A.; Mendes, A. Integrated design of hematite and dye-sensitized solar cell for unbiased solar charging of an organic-inorganic redox flow battery. *Nano Energy* **2019**, *62*, 832–843. [[CrossRef](#)]
11. Shen, L.; Li, H.; Meng, X.; Li, F. Transfer printing of fully formed microscale InGaP/GaAs/InGaNaSb cell on Ge cell in mechanically-stacked quadruple-junction architecture. *Sol. Energy* **2020**, *195*, 6–13. [[CrossRef](#)]
12. Lin, S.; Xie, S.; Lei, Y.; Gan, T.; Wu, L.; Zhang, J.; Yang, Y. Betavoltaic battery prepared by using polycrystalline CdTe as absorption layer. *Opt. Mater.* **2022**, *127*, 112265.
13. Yeojun, Y.; Sunghyun, M.; Sangin, K.; Jaejin, L. Flexible fabric-based GaAs thin-film solar cell for wearable energy harvesting applications. *Sol. Energy Mater. Sol. Cells* **2022**, *246*, 111930.
14. Fatemeh, G.Y.; Ali, F. Performance enhancement of CIGS solar cells using ITO as buffer layer. *Micro Nanostruct.* **2022**, *168*, 207289.
15. Minemoto, T.; Matsui, T.; Takakura, H.; Hamakawa, Y.; Negami, T.; Hashimoto, Y.; Uenoyama, T.; Kitagawa, M. Theoretical analysis of the effect of conduction band offset of window/CIS layers on performance of CIS solar cells using device simulation. *Sol. Energy Mater. Sol. Cells* **2001**, *67*, 83–88.
16. Pan, B.; Wei, M.; Liu, W.; Jiang, G.; Zhu, C. Fabrication of $\text{Cu}_2\text{ZnSnS}_4$ absorber layers with adjustable Zn/Sn and Cu/Zn+Sn ratios. *J. Mater. Sci. Mater. Electron.* **2014**, *25*, 3344–3352. [[CrossRef](#)]
17. Tablero, C. Effect of the oxygen isoelectronic substitution in $\text{Cu}_2\text{ZnSnS}_4$ and its photovoltaic application. *Thin Solid Films* **2012**, *520*, 5011–5013. [[CrossRef](#)]
18. Su, Z.H.; Liang, G.X.; Fan, P.; Luo, J.; Zheng, Z.; Xie, Z.; Wang, W.; Chen, S.; Hu, J.; Wei, Y.; et al. Device Postannealing Enabling over 12% Efficient Solution-Processed $\text{Cu}_2\text{ZnSnS}_4$ Solar Cells with Cd^{2+} Substitution. *Adv. Mater.* **2020**, *32*, 2000121. [[CrossRef](#)]
19. Andrea, C.; Ole, H. What is the band alignment of $\text{Cu}_2\text{ZnSn(S,Se)}_4$ solar cells. *Sol. Energy Mater. Sol. Cells* **2017**, *169*, 177–194.
20. Bao, W.; Ichimura, M. Prediction of the Band Offsets at the CdS/ $\text{Cu}_2\text{ZnSnS}_4$ Interface Based on the First-Principles Calculation. *Jpn. J. Appl. Phys.* **2012**, *51*, 10NC31. [[CrossRef](#)]
21. Su, Z.; Tan, J.; Li, X.; Zeng, X.; Batabyal, S.K.; Wong, L.H. Cation Substitution of Solution-Processed $\text{Cu}_2\text{ZnSnS}_4$ Thin Film Solar Cell with over 9% Efficiency. *Adv. Energy Mater.* **2015**, *5*, 1500682. [[CrossRef](#)]
22. Kim, S.; Kim, K.M.; Tampo, H.; Shibata, H.; Niki, S. Improvement of voltage deficit of Ge-incorporated kesterite solar cell with 12.3% conversion efficiency. *Appl. Phys. Express* **2016**, *9*, 102301. [[CrossRef](#)]
23. Chen, S.; Gong, X.G.; Wei, S.H. Band-structure anomalies of the chalcopyrite semiconductors CuGaX_2 versus AgGaX_2 (X=S and Se) and their alloys. *Phys. Rev. B* **2007**, *75*, 205209. [[CrossRef](#)]
24. Walsh, A.; Chen, S.; Wei, S.H.; Gong, X.G. Kesterite Thin-Film Solar Cells: Advances in Materials Modelling of $\text{Cu}_2\text{ZnSnS}_4$. *Adv. Energy Mater.* **2012**, *2*, 400–409. [[CrossRef](#)]
25. Tian, Y.; Zheng, Q.; Zhao, J. Tensile Strain-Controlled Photogenerated Carrier Dynamics at the van der Waals Heterostructure Interface. *J. Phys. Chem. Lett.* **2020**, *11*, 586–590. [[CrossRef](#)] [[PubMed](#)]
26. Fan, L.L.; Chen, S.; Luo, Z.L.; Liu, Q.H.; Wu, Y.F.; Song, L.; Ji, D.X.; Wang, P.; Chu, W.S.; Gao, C.; et al. Strain Dynamics of Ultrathin VO_2 Film Grown on TiO_2 (001) and the Associated Phase Transition Modulation. *Nano Lett.* **2014**, *14*, 4036–4043. [[CrossRef](#)]
27. LYU, L.; Yang, Y.Y.; CEN, W.F.; YAO, B.; OU, J.K. First-principles Study on Optical Properties of Cubic Ca_2Ge under Stress Effect. *Bull. Chin. Ceram. Soc.* **2019**, *38*, 3788–3795.
28. Yan, W.J.; Zhang, C.H.; Gui, F.; Zhang, Z. Electronic Structure and Optical Properties of Stressed $\beta\text{-FeSi}_2$. *Acta Opt. Sin.* **2013**, *33*, 243–249.
29. Lazarovits, B.; Kim, K.; Haule, K.; Kotliar, G. Effects of strain on the electronic structure of VO_2 . *Phys. Rev. B* **2010**, *81*, 115117. [[CrossRef](#)]
30. Manyk, T.; Rutkowski, J.; Kopytko, M.; Martyniuk, P. Theoretical Study of the Effect of Stresses on Effective Masses in the InAs/InAsSb Type-II Superlattice. *Eng. Proc.* **2022**, *21*, 16.
31. Schorr, S.; Hoebler, H.J.; Tovar, M. A neutron diffraction study of the stannite-kesterite solid solution series. *Eur. J. Mineral.* **2007**, *19*, 65–73.
32. Segall, M.D.; Lindan, P.; Probert, M.J.; Pickard, C.J.; Hasnip, P.J.; Clark, S.J.; Payne, M.C. First-principles simulation: Ideas, illustrations and the CASTEP code. *J. Phys. Condens. Matter* **2002**, *14*, 2717–2744. [[CrossRef](#)]
33. Perdew, J.P.; Burke, K.; Ernzerhof, M. Generalized Gradient Approximation Made Simple. *Phys. Rev. Lett.* **1996**, *77*, 3865–3868. [[CrossRef](#)] [[PubMed](#)]
34. Vanderbilt, D. Soft self-consistent pseudopotentials in a generalized eigenvalue formalism. *Phys. Rev. B* **1990**, *41*, 7892–7895. [[CrossRef](#)] [[PubMed](#)]
35. Liu, J. First-Principles Prediction of Structural, Elastic, Mechanical, and Electronic Properties of $\text{Cu}_2\text{ZnSnS}_4$ under Pressure. *ECS J. Solid State Sci. Technol.* **2022**, *11*, 073011. [[CrossRef](#)]
36. Kumar, M.; Zhao, H.; Persson, C. Cation vacancies in the alloy compounds of $\text{Cu}_2\text{ZnSn(S}_{1-x}\text{Se}_x)_4$ and $\text{CuIn(S}_{1-x}\text{Se}_x)_2$. *Thin Solid Films* **2013**, *535*, 318–321. [[CrossRef](#)]
37. Kheraj, V.; Patel, K.K.; Patel, S.J.; Shah, D.V. Synthesis and characterisation of Copper Zinc Tin Sulphide (CZTS) compound for absorber material in solar-cells. *J. Cryst. Growth* **2013**, *362*, 174–177.
38. Zhao, H.; Persson, C. Optical properties of Cu(In,Ga)Se_2 and $\text{Cu}_2\text{ZnSn(S,Se)}_4$. *Thin Solid Films* **2011**, *519*, 7508–7512. [[CrossRef](#)]

39. Hoo, Q.Y.; Xu, Y. Detection of dielectric screening effect by excitons in two-dimensional semiconductors and its application. *Acta Phys. Sin.* **2022**, *71*, 124–138.
40. Prokopidis, K.; Kallialakis, C. Physical interpretation of a modified Lorentz dielectric function for metals based on the Lorentz–Dirac force. *Appl. Phys. B* **2014**, *117*, 25–32. [[CrossRef](#)]
41. Nainaa, F.Z.; Bekkioui, N.; Abbassi, A.; Ez, Z.H. First principle study of structural, electronic optical and electric properties of $\text{Ag}_2\text{MnSnS}_4$. *Comput. Condens. Matter* **2019**, *22*, e00443. [[CrossRef](#)]
42. Scragg, J.J.; Dale, P.J.; Peter, L.M.; Zoppi, G.; Forbes, I. New routes to sustainable photovoltaics: Evaluation of $\text{Cu}_2\text{ZnSnS}_4$ as an alternative absorber material. *Phys. Status Solidi* **2010**, *245*, 1772–1778. [[CrossRef](#)]
43. Kahlaoui, S.; Belhorma, B.; Labrim, H.; Boujnah, M.; Regragui, M. Strain effects on the electronic, optical and electrical properties of $\text{Cu}_2\text{ZnSnS}_4$: DFT study. *Heliyon* **2020**, *6*, e03713. [[CrossRef](#)] [[PubMed](#)]