

Review **Development of Thermoelectric Half-Heusler Alloys over the Past 25 Years**

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Abstract: Half-Heusler alloys are among the most promising thermoelectric materials. In the present review, thermoelectric properties (at 300 K and 800 K) of more than 1100 compositions from more than 220 publications between 1998 and 2023 were collected and evaluated. The dependence of the peak figure of merit, *ZT*max, of p- and n-type half-Heusler alloys on the publishing year and the peak temperature is displayed in several figures. Furthermore, plots of *ZT* vs. the electrical resistivity, the Seebeck coefficient and the thermal conductivity at 300 K and 800 K are shown and discussed. Especially thermal conductivity vs. power factor leads to a good overview of *ZT*. For both p- and n-type individually separated into systems, *ZT*s and peak *ZT*s in dependence on the composition are displayed and discussed. This overview can help to find the ideal half-Heusler alloy for practical use.

Keywords: half-Heusler alloys; physical properties; figure of merit (*ZT*)

1. Introduction

Thermoelectric (TE) modules and generators have the ability to directly convert waste heat into electric power and thus can be considered alternative, sustainable and "green" energy sources. There exist various classes of TE materials such as tellurides, selenides, clathrates, silicides, oxides, Zintl phases, antimonides, skutterudites, organic semiconductors, Heusler and half-Heusler alloys, etc., each of them qualifying in a defined temperature range. To judge the TE quality, the figure of merit, $ZT = S^2T/\rho(\lambda_e + \lambda_{ph})$, is used, where S is the Seebeck coefficient, ρ is the electrical resistivity, λ is the thermal conductivity, consisting of an electronic part, λ_e , and a phonon (lattice) part, λ_{ph} , and T is the temperature. To arrive at a high *ZT*, the power factor, $pf = S^2/\rho$, should be high, whereas λ should be low.

Half-Heusler (HH) alloys are promising TE materials intended for mid-to-high temperature power generation applications with already high *ZT* values. In addition, the starting materials are available, abundant, and, if hafnium and noble metals are avoided, cheap, a fact very important in respect of mass production.

Half-Heusler (HH) alloys are named after Friedrich Heusler (1866–1947), a German chemist and mining engineer. Half-Heusler alloys are members of the vast family of Heusler alloys with the general composition X_2 YZ, consisting of three interpenetrating face-centered cubic (fcc) sublattices (space group *Fm*3*m*). Reduction of symmetry (non-centrosymmetric space group $F-43m$) splits the X_2 sublattice of multiplicity eight into two sublattices of multiplicity four, of which one is empty, resulting in the formula XYZ of the so-called HH phases. The big advantage in the optimization of thermoelectric properties of HH compounds is the opportunity to dope each of the four sublattices individually. Besides that, HH alloys tend from metallic to half-metallic alloys, exhibit interesting magnetic properties and are small band gap semiconductors. Generally, one can divide HH alloys into the following main groups: MNiSn, MCoSb ($M = Ti$, Zr , Hf), MFeSb ($M = V$, Nb, Ta), the 19-electron system and other HH alloys not covered by these four systems.

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Reviews as papers, chapters in a book, or even a book have been published [\[1–](#page-17-0)[8\]](#page-17-1) focusing mainly on the crystal structure, magneto-optical properties, electron–phonon interactions, mobility, ways to enhance the power factor, new advances, applications and so on. Freer et al. [\[9\]](#page-17-2) recently published tables of TE materials: the chapter about HH alloys was handled by S. Han, C. Fu and T. Zhu. They collected TE parameters from about 90 publications of p- and n-type HH alloys, however, from each publication for the compound with the highest *ZT* only (for room temperature and the temperature of the peak *ZT*). These peak *ZT*s were also displayed as a figure.

In the current review, measured data of the Seebeck coefficient, the electrical resistivity and the thermal conductivity of all compositions in each published paper were collected, as well as the calculated power factor and *ZT* (at 300 K and 800 K) and particularly the peak *ZT*. As more than 220 papers were evaluated, this adds up to about 300 p-type and 810 n-type compositions. These data were plotted and evaluated to gain a deeper insight, e.g., *ZT* vs. Seebeck coefficient, resistivity, thermal conductivity, composition and pf vs. thermal conductivity.

2. Experimental Production Methods

There exist various methods to produce half-Heusler alloys, the basics of the most popular ones will be described as follows.

The historically favored route, and still the mostly used one, is the melting-annealing one. With some variations, the essential steps are to arc-melt the defined amounts of high-purity elemental pieces under argon, flip the reguli and remelt them several times for good homogenization. Oxygen contamination must be avoided; therefore, a certain vacuum level must be kept in combination with high-purity argon (the use of oxygen-getter material, such as Zr or Ti, which is melted prior to the actual samples, can further improve the sample quality). In some cases, the samples are additionally melted in a high-frequency induction furnace several times. Between the melting steps, it is possible to frequently break them into pieces and turn them upside down and outside-inside. As a last step, they are wrapped into protective Mo foils, vacuum sealed in a quartz ampulla and annealed at high temperatures, usually between 900–1000 ℃ for at least 48 h to equilibrate the samples and to further increase homogeneity.

This method is also suitable for big quantities (ingots with a weight of >1 kg) when stoichiometric amounts of pure elements are synthesized together in a vacuum induction furnace, and the melt is cast into a mold.

A similar production method is levitation melting (e.g., [\[10](#page-17-3)[–12\]](#page-18-0)), which is becoming more and more popular in recent years. The main advantage of levitation melting is that (i) the sample has no contact with the environment as well as (ii) the eddy currents, which improve the homogeneity of the melt.

All these above-described methods are followed by breaking the solidified samples and pulverizing them manually or, in recent years more often practiced, ball-mill (BM) or high-energy ball-mill (HBM) and consolidating them either via hot-pressing (HP) or spark plasma sintering (SPS). For half-Heusler alloys containing materials with a high vapor pressure at elevated temperatures, such as antimony, it is necessary to compensate for antimony vaporization losses before HP.

Mechanical alloying (MA) is a fast method (e.g., $[13-16]$ $[13-16]$) to prepare homogeneous powder mixtures prior to consolidation in HP or spark plasma sintering devices.

In the so-called solid-state reaction process (e.g., [\[17\]](#page-18-3)), the powder mixture of highpurity constituents is heated (almost 1000 ◦C) under a flowing argon atmosphere for up to a week, usually followed by SPS.

Self-propagation high-temperature synthesis (SHS) or combustion synthesis (CS) is a scalable, rapid preparation method. It is based on a propagating self-sustained exothermic reaction in consolidated elemental powders of the desired composition, and it yields a homogeneous product similar to the ones obtained by the well-known zone melting process. In addition, the SHS process is so fast that evaporation of low-melting-point elements is

dramatically suppressed, leading to a precisely controlled composition (for details, see, e.g., [\[18–](#page-18-4)[20\]](#page-18-5)).

In some cases, SHS is combined with the selective laser melting (SLM) method. For SLM, a laser beam is used to melt a single-layer powder bed that rapidly solidifies as soon as the laser moves away. Layer upon layer, a three-dimensional object can be formed.

Microwave preparation (e.g., [\[21](#page-18-6)[–24\]](#page-18-7)) is another fast way to prepare half-Heusler alloys. Stoichiometric amounts of metal powders are cold-pressed into bar-shaped pellets, directly sealed into evacuated quartz tubes and placed into a crucible filled with some granular carbon, which acts as the microwave susceptor material. The sample is surrounded by alumina insulation foam to minimize heat loss. Finally, the whole set-up is placed into a commercial microwave reactor with a rotating plate. The reactions are allowed to run at 100% power (e.g., 700 W) for 1 min.

In some cases, the bulk HH alloy is processed by high-pressure torsion to further improve the TE performance by severe plastic deformation resulting in a drastic reduction towards ultra-low thermal conductivity, which occurs due to grain refinement and deformation-induced defects (for details, see references [\[25](#page-18-8)[–27\]](#page-18-9)).

Whenever the phonon part of the thermal conductivity was not published, it was calculated by the authors: $\lambda_{ph} = \lambda - \lambda_e$ with λ as the total thermal conductivity and λ_e as the electron part; $\lambda_e = LT/\rho$ (Wiedemann–Franz law) with L as the Lorenz number (calculated as suggested by Kim et al. $[28]$) and ρ is the electrical resistivity. In all cases where the power factor and/or *ZT* were not published, it was calculated by the authors from $pf = S^2/\rho$ and $ZT = S^2T/(\rho\lambda)$, respectively.

3. Results and Discussion

3.1. General Overview

In this chapter, all data from references [\[10–](#page-17-3)[27,](#page-18-9)[29](#page-18-11)[–218\]](#page-25-0) were used for figures and discussions. The references [\[29](#page-18-11)[–218\]](#page-25-0) are numbered by the publication year, within one year, in alphabetical order.

Looking at the highest *ZT*s published during the last 25 years, one can see in Figure [1](#page-3-0) that, as already mentioned in the introduction, considerably more n-type half-Heusler alloys were investigated than p-type (this review comprises 294 p-type compounds and 816 n-type compounds). It is also clearly visible that the trend of high peak *ZT*s increases almost linearly from 2004 to 2018. There are exceptions for the n-type, such as (i) the work of Sakurada et al. [\[43\]](#page-19-0), who published already in 2005 outstanding high values (*ZT* = 1.25 to 1.5) for $Ti_{0.5}Zr_{0.25}Hf_{0.25}N$ iSn doped with small amounts of antimony, or (ii) from Shen et al. [\[17,](#page-18-3)[33\]](#page-18-12) and Kawasaki et al. [\[37\]](#page-19-1), who reached almost *ZT*s = 1 at 800 K in 2001 and 2004 for $ZrNiSn_0.99Sb_{0.01}$. It is also worth mentioning that Yu et al. [\[150\]](#page-23-0) reached $ZT = 1.6$ at 1200 K for his p-type HH series ($Nb_{1-x}Ta_{x}$)_{0.8}Ti_{0.2}FeSb. Furthermore, it seems that after a boom during 2015–2020, the number of half-Heusler publications has dropped significantly.

Figure [2](#page-3-1) displays the highest *ZT*s as a function of temperature. It is necessary to point out that in some cases *ZT*max is not really the peak *ZT* of the respective sample but the *ZT* at the highest measured temperature. This is especially the case for publications earlier than 2005. Besides that, in most cases *ZT*max is published for temperatures between 700 K and 1000 K; the highest temperatures for ZT_{max} for p- as well as n-type are around 1200 K.

Figures [3](#page-3-2) and [4](#page-4-0) depict the dependence of the power factor on the thermal conductivity at 300 K and 800 K, respectively. Figure [3](#page-3-2) shows that at 300 K, almost all power factor values for the n-type are within the range of 4 mW/mK^2 for thermal conductivities up to 250 mW/cmK, whereas the p-type displays much higher power factors, reaching values higher than 10 mW/mK 2 for thermal conductivities up to about 150 mW/cmK, e.g., for the $Nb_{1-x}Ti_xFeSb$ -series of He et al. [\[131\]](#page-22-0) or for $Nb_{0.95}Hf_{0.05}FeSb$ and $Nb_{0.95}Zr_{0.05}FeSb$ of Ren and colleagues [\[160\]](#page-23-1).

Figure 1. Peak *ZT*, *ZT*max, of p- and n-type half-Heusler alloys vs. year of publication. **Figure 1.** Peak *ZT*, *ZT*max, of p- and n-type half-Heusler alloys vs. year of publication.

Figure 2. Peak *ZT*, *ZT*max, of p- and n-type half-Heusler alloys vs. temperature, T. **Figure 2.** Peak *ZT*, *ZT*max, of p- and n-type half-Heusler alloys vs. temperature, T.

Figure 3. P- and n-type half-Heusler alloys: power factor, pf, vs. thermal conductivity, λ, at 300 K. **Figure 3.** P- and n-type half-Heusler alloys: power factor, pf, vs. thermal conductivity, λ, at 300 K.

Figure 4. P- and n-type half-Heusler alloys: power factor, pf, vs. thermal conductivity, λ, at 800 K. **Figure 4.** P- and n-type half-Heusler alloys: power factor, pf, vs. thermal conductivity, λ, at 800 K.

Figure 3. P- and n-type half-Heusler alloys: power factor, pf, vs. thermal conductivity, λ, at 300 K.

The situation has changed comparing this figure now with the one for 800 K (Figure [4\)](#page-4-0), 4), and as for p- and n-type, all data are mainly evenly distributed within a power factor– and as for p- and n-type, all data are mainly evenly distributed within a power factor– t_{bound} conductivity window of almost $0, 5.5$ mW/mK² by 20–80 mW thermal conductivity window of almost $0-5.5$ mW/mK² by 20–80 mW/cmK.

Lines for the resulting *ZT* values are drawn in both Figures [3](#page-3-2) and [4,](#page-4-0) revealing that at 300 K, no compound exceeds $ZT = 0.5$. Whilst many of the p-type compounds range between $ZT = 0.1$ and 0.25 or have values as low as almost 0 , one can find the values for the n-type all over, the very low values due to high thermal conductivities. At 800 K, both p- and n-type exceed the *ZT* = 1.5 line. The majority of data can be found for thermal conductivities between 20 and 80 mW/cmK, representing *ZT*s from almost 0 to *ZT* = 1.5. This figure also illustrates that p-type HH materials generally exhibit lower thermal conductivities than the Figures 5 and 6 depict the relation between electrical resistivity and *ZT* at 300 K and n-type.

Figures [5](#page-4-1) and [6](#page-5-0) depict the relation between electrical resistivity and *ZT* at 300 K and 800 K, respectively. It shows that, especially in the 300 K temperature region, individual specimens have very high resistivities, resulting, as a consequence, in very low ZTs. Generdisplays lower resignation *ZTs* in the range of 10−3 or even find *ZTs* in the range of 10−3 or even ally, the electrical resistivity decreases with increasing temperature. Figure [6](#page-5-0) displays lower lower. The insert, a cutout, in each *ZT*-resistivity figure, gives a better overview, indicating resistivity values. Still, here one can find *ZT*s in the range of 10−³ or even lower. The insert, a cutout, in each *ZT*-resistivity figure, gives a better overview, indicating that, especially for 800 K, most data can be found between 500 and 2500 µΩcm. In this region are also the region are also the highest *ZT*s. highest *ZT*s.

Figure 5. ZT at 300 K vs. electrical resistivity, ρ , for p- and n-type half-Heusler alloys. Insert: cutout for $ρ < 15000 μΩ$ cm.

Figure 6. ZT at 800 K vs. electrical resistivity, ρ , for p- and n-type half-Heusler alloys. Insert: cutout for ρ < 12000 µΩcm. for ρ < 12000 µΩcm.

Figur[es](#page-5-1) 7 a[nd](#page-6-0) 8 present the dependence of ZT on the Seebeck coefficient. Especially at at 300 K, the area of Seebeck values is much wider for the n- than for the p-type. Generally, 300 K, the area of Seebeck values is much wider for the n- than for the p-type. Generally, it seems that an absolute Seebeck value of 200 \pm 20 μ V/cm leads to high *ZTs,* a great finding for the production of new half-Heuslers. for the production of new half-Heuslers.

Figure 7. *ZT* at 300 K vs. Seebeck coefficient, S, for p- and n-type half-Heusler alloys. **Figure 7.** *ZT* at 300 K vs. Seebeck coefficient, S, for p- and n-type half-Heusler alloys.

Whereas for the high-temperature regime (Figure [9\)](#page-6-1), the data *ZT* vs. λ are more or less evenly mixed at 300 K (insert in Figure [9\)](#page-6-1), most p-type half-Heuslers have thermal conductivities between 20 and 70 mW/cmK. Of course, in both regions, *ZT*s are high for low thermal conductivities. Comparing the total and the lattice thermal conductivities in Figures [9](#page-6-1) and [10,](#page-6-2) the distribution of the data looks very similar, indicating that the electron part of the thermal conductivity is rather low, and, as a consequence, the phonon part is in the range of the total thermal conductivity.

Figure 8. *ZT* at 800 K vs. Seebeck coefficient, S, for p- and n-type half-Heusler alloys. **Figure 8.** *ZT* at 800 K vs. Seebeck coefficient, S, for p- and n-type half-Heusler alloys.

Figure 7. *ZT* at 300 K vs. Seebeck coefficient, S, for p- and n-type half-Heusler alloys.

Figure 9. ZT at 800 K vs. thermal conductivity, λ , for p- and n-type half-Heusler alloys. Insert: ZT at 300 K vs. thermal conductivity, λ. 300 K vs. thermal conductivity, λ.

Figure 10. *PT* at 800 K vs. lattice thermal conductivity, γ _{ph}, for p- and n-type half-Heusler allows. Insert: *ZT* at 300 K vs. lattice thermal conductivity, λph. Figure 10. ZT at 800 K vs. lattice thermal conductivity, λ_{ph} , for p- and n-type half-Heusler alloys. Insert: *ZT* at 300 K vs. lattice thermal conductivity, λ_{ph} .

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Table 1. N-type half-Heusler alloys.

Table 1. *Cont.*

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Table 2. P-type half-Heusler alloys.

Table 2. *Cont.*

3.2. N-Type Half-Heusler Alloys

All compounds with the respective references can be found in Table [1.](#page-10-0) In some cases (indicated by *), measurements were only performed below room temperature.

3.2.1. (Ti,Zr,Hf)NiSn-System

Comparing the *ZT* values of TiNiSn, ZrNiSn and HfNiSn, as shown in Figures [11](#page-12-0) and [12,](#page-12-1) one can immediately see that ZrNiSn is the most investigated alloy of these three, whereas HfNiSn is not so popular. It is evident from Figures [11](#page-12-0) and [12](#page-12-1) that the highest *ZT*s at 300 K as well as at 800 K in this system are for (Zr,Hf)NiSn and (Ti,Zr,Hf)NiSn and even more so when Sb-doped at the Sn site.

Figure 11. *ZT* at 300 K for the HH n-type (Ti,Zr,Hf)NiSn-system. **Figure 11.** *ZT* at 300 K for the HH n-type (Ti,Zr,Hf)NiSn-system.

Figure 12. *ZT* at 800 K and *ZT*max for the HH n-type (Ti,Zr,Hf)NiSn-system. **Figure 12.** *ZT* at 800 K and *ZT*max for the HH n-type (Ti,Zr,Hf)NiSn-system. **Figure 12.** *ZT* at 800 K and *ZT*max for the HH n-type (Ti,Zr,Hf)NiSn-system.

 $ZT = 0.9$ at 773 K for ZrNiSn and Liu et al. $[119] ZT = 0.8$ at 1000 K for HfNiSn. In comparison to these values, those for compounds with Ti and Zr, Ti and Hf, Zr and Hf or with Ti and Zr
and Hf and have high proclamation of the Life MiCa ZT, 12, t 272 K (Chanker and Hf reach much higher values, such as for $Zr_{0.7}Hf_{0.3}N$ iSn $ZT = 1.2$ at 873 K (Chauhan with Ti and Times in the Column of the Studies o
The studies of Light and *Zames and Zames and 2000 K* (Decleted Light) K (κ ₂₁ μ ₁₅), or CVCR21 = 1.3 at 660 K (κ ₀2 C can μ ₁₁₀). Gürth et al. [129] published *ZT* = 1 at 823 K for TiNiSn and for ZrNiSn, Misra et al. Gürth et al. [\[129\]](#page-22-7) published *ZT* = 1 at 823 K for TiNiSn and for ZrNiSn, Misra et al. [\[106\]](#page-21-1) et al. [\[154\]](#page-23-9)), for Ti_{0.5}Zr_{0.25}Hf_{0.25}NiSn *ZT* = 1.3 at 693 K (Sakurada and Shutoh [\[43\]](#page-19-0); Shutoh and Sakurada [\[47\]](#page-19-12)) or even $ZT = 1.5$ at 800 K (Rogl et al. [\[146\]](#page-22-9)). [43]; Shutoh and Sakurada [47]) or even *ZT* = 1.5 at 800 K (Rogl et al. [146]).

Doping in the (Ti, Zr, Hf) NiSn-system with X ($X = Sc$, Y , V , Nb , $V + Nb$, Ta , Mn , Co , Cu, $Co + Cu$, Pd, Pt, Pb) (see also Table [1\)](#page-10-0) was not really successful as generally, all these alloys did not have higher *ZT*s; however, from Figure [11](#page-12-0) as well as from Figure [12](#page-12-1) it is obvious that with an additional substitution of Sn in many cases *ZT* is enhanced. Usually, Sn is substituted by a very small amount of Sb or Ge, but also, in rarer cases, In, Si, Pb, Bi or Te was used. The highest *ZT* values were reached for $Ti_{0.5}Zr_{0.25}Hf_{0.25}NiSn_{0.998}Sb_{0.002}$ at 700 K (*ZT* = 1.5, Shutoh and Sakurada [\[47\]](#page-19-12)) and $Ti_{0.5}Zr_{0.25}Hf_{0.25}NiSn_{0.99}Sb_{0.01}$ at 825 K (*ZT* = 1.45, Rogl et al. [\[146\]](#page-22-9)). Besides that, there are several compounds with *ZT* ∼ 1.3, such as $(Hf_{0.6}Zr_{0.4})_{0.99}V_{0.01}NiSn_{0.995}Sb_{0.005}$ at 900 K (Chen et al. [\[140\]](#page-22-17)) or Ti_{0.49}Zr_{0.49}V_{0.02}NiSn_{0.98} $Sb_{0.02}$ at 823 K (Rogl et al. [\[146\]](#page-22-9)) (see also Figure [12\)](#page-12-1).

Various groups used additions (not shown in any figure but presented in Table [1\)](#page-10-0), mainly nanoparticles, to reduce the thermal conductivity and to enhance *ZT*. Schwall and Balke [\[80\]](#page-20-12) incorporated Nb into $Zr_{0.5}Hf_{0.5}N$ iSn but presented data only below room temperature. Akram et al. [\[108\]](#page-21-2) admixed La with ZrNiSn: the enhancement in *ZT* was within the error bar of *ZT* (from 0.52 to 0.54), which also applies to Visconti et al. [\[149\]](#page-23-8), who used HfO₂ for $Ti_{0.5}Zr_{0.5}NiSn_{0.994}Sb_{0.006}$ (ZT rises from 0.9 to 1). A better result was achieved by Chauhan et al. [\[153\]](#page-23-3) by mixing ZnO with ZrNiSn (highest *ZT* = 1). Nanoparticles of ZrO_2 were added to ZrNiSn (Huang et al. [\[35\]](#page-18-20)) as well as to $Zr_{0.5}Hf_{0.5}Ni_{0.8}Pd_{0.2}Sn_{0.99}Sb_{0.01}$ and $Zr_{0.25}Hf_{0.6}Ti_{0.15}NiSn_{0.995}Sb_{0.005}$ (Chen et al. [\[49\]](#page-19-8) and [\[127\]](#page-22-13), respectively). For $ZrNiSn$, the "enhancement" of *ZT* was within the error bar; however, it was much better for the other two compounds with $ZT_{\text{max}} = 0.74$ and even $ZT_{\text{max}} = 1.3$ for the Ti-doped half-Heusler alloy. Additionally, a rather high $ZT_{\text{max}} = 1.35$ at 873 K was achieved by Kang et al. [\[197\]](#page-25-10) for $Zr_{0.4}Hf_{0.6}NiSn_{0.99}Sb_{0.01}$ plus 2 weight percent tungsten nanoparticles. Makongo et al. [\[77,](#page-20-10)[78\]](#page-20-11) and Douglas et al. [\[83\]](#page-20-1) added to $Zr_{0.25}Hf_{0.75}N$ iSn and TiNiSn the corresponding full Heusler phase, i.e., $Zr_{0.25}Hf_{0.75}Ni_2Sn$ and TiNi₂Sn, respectively; however, the outcome was rather disappointing with low *ZT*s < 1, in some cases even lower than *ZT* without the full-Heusler phase.

Rogl et al. [\[25,](#page-18-8)[26\]](#page-18-25) processed arc-melted, annealed, ball-milled and hot-pressed $Ti_{0.5}Zr_{0.5}$ NiSn and $Ti_{0.5}Zr_{0.5}NiSn_{0.98}Sb_{0.02}$ samples by high-pressure torsion to improve their thermoelectric performance via a drastic reduction of the thermal conductivity due to grain refinement and a high concentration of defects inferred by severe plastic deformation. Whereas for $Ti_{0.5}Zr_{0.5}N$ ish, the thermally stable alloy showed an enhancement of ZT_{max} of 20%, for $Ti_{0.5}Zr_{0.5}NiSn_{0.98}Sb_{0.02} Zr_{\text{max}}$ was, within the error bar, about the same.

As a summary, one can say that the (Ti,Zr,Hf)NiSn-system comprises TE materials of high interest.

3.2.2. (Ti,Zr,Hf)CoSb-System, (V,Nb,Ta)FeSb-System, 19-Electron System and Other N-Type Half-Heusler Alloys

In Figures [13](#page-14-0) and [14,](#page-14-1) the *ZT*s are shown for the (Ti,Zr,Hf)CoSb-system, the (V,Nb,Ta) FeSb-system, the 19-electron system and of all those half-Heuslers, which do not fit into any of the so-far mentioned categories. For all these systems, only a few half-Heuslers with Sb/Sn substitutions were published; generally, *ZT*s are much lower in comparison with the (Ti,Zr,Hf)NiSn-system for 300 K as well as for 800 K.

At 800 K (Figure [14\)](#page-14-1), almost all *ZTs* are lower than 0.7 and also *ZT*_{max} rarely exceeds 1. He et al. [\[141\]](#page-22-18) reached $ZT > 1$ for $Zr_{0.5}Hf_{0.5}Co_{0.9}Ni_{0.1}Sb$ at 1073 K, Liu et al. [\[158\]](#page-23-12) acquired *ZT*s of almost 1 and around 1 for the series $(Zr_{1-x}Hf_x)_{0.88}Nb_{0.12}CoSb$. Xia et al. [\[116](#page-21-15)[,165\]](#page-23-15) investigated Nb1−xCoSb and Nb0.8Co1−xNixSb with *ZT* = 0.8–0.9. Zhu et al. [\[187\]](#page-24-8) obtained *ZT* = 0.85 at 800 K with a peak *ZT* = 1.04 at 972 K for *Zr*Co_{0.9}Ni_{0.1}Bi_{0.85}Sb_{0.15}.

Figure 13. *ZT* at 300 K for the HH n-type (Ti,Zr,Hf)CoSb-system, (V,NbTa)FeSb-system, 19-electron **Figure 13.** *ZT* at 300 K for the HH n-type (Ti,Zr,Hf)CoSb-system, (V,NbTa)FeSb-system, 19-electron system and other half-Heuslers. system and other half-Heuslers. $_{\rm ^{1}1}$ gure 13. ZI at 300 K for the HF

Figure 14. ZT at 800 K for the HH n-type (Ti,Zr,Hf)CoSb-system, (V,NbTa)FeSb-system, 19-electron system and other half-Heuslers. system and other half-Heuslers. system and other half-Heuslers.

At 800 K (Figure 14), almost all *ZT*s are lower than 0.7 and also *ZT*max rarely exceeds At 800 K (Figure 14), almost all *ZT*s are lower than 0.7 and also *ZT*max rarely exceeds *3.3. P-Type Half-Heusler Alloys*

All compounds with the respective references can be found in Table [2.](#page-11-0) In some cases, measurements were only performed below room temperature. Figures [15](#page-15-0) and [16](#page-15-1) represent the ZTs for p-type half-Heusler alloys at 300 K and 800 K as well as the peak ZTs, respectively.

*ZT*s, respectively.

Figure 15. *ZT* at 300 K for the p-type HH alloys. **Figure 15.** *ZT* at 300 K for the p-type HH alloys.

Figure 16. *ZT* at 800 K and *ZT*max for the p-type HH alloys. **Figure 16.** *ZT* at 800 K and *ZT*max for the p-type HH alloys.

When looking at Sectio[n 3.3](#page-14-2), it is obvious that many more n-type than p-type alloys When looking at Section 3.3, it is obvious that many more n-type than p-type alloys were investigated. As depicted in Figure 1[5, at](#page-15-0) 300 K, all *ZT*s are below 0.3; at 800 K (Figure 16), m[ost](#page-15-1) *ZT*s are lower than 1; however, several half-Heuslers in the (NbX")FeSb series with X" = Ti, Zr, Hf, Zr and Hf, as well as Ti + Ta, reach $ZT = 1.4$ at 800 K; the values *ZT*max in some cases reach or even exceed 1.4 and go up to 1.6. for *ZT*max in some cases reach or even exceed 1.4 and go up to 1.6.

It can be seen in Figure[s 15](#page-15-0) an[d 16](#page-15-1) that only data of two p-type compounds are It can be seen in Figures 15 and 16 that only data of two p-type compounds are available for the (Ti,Zr,Hf)NiSn-system, namely (Zr,Sc)NiSn [122] and (Zr,Y)NiSn [114] available for the (Ti,Zr,Hf)NiSn-system, namely (Zr,Sc)NiSn [\[122\]](#page-22-5) and (Zr,Y)NiSn [\[114\]](#page-21-5) with rather low *ZT*s. with rather low *ZT*s.

In the CoSb-based system, almost all compounds are substituted at the Sb-site by Sn, In the CoSb-based system, almost all compounds are substituted at the Sb-site by Sn, except for TiCoSb, which was also substituted by Ge [68]. Rausch et al. published in two except for TiCoSb, which was also substituted by Ge [\[68\]](#page-20-20). Rausch et al. published in two separate works [\[120](#page-21-18)[,121\]](#page-22-24) rather high *ZTs* of the (Ti,Zr,Hf)CoSb_{1−x}Sn_x, series: the highest $ZT_{\text{max}} = 1.4$ for $\text{Ti}_{0.12}\text{Zr}_{0.44}\text{Hf}_{0.44}\text{CoSb}_{0.8}\text{Sn}_{0.2}$ at 963 K [\[121\]](#page-22-24).

In the FeSb-based system, Yu et al. [\[150\]](#page-23-0) achieved outstanding high peak *ZT* values $(ZT = 1.2$ –1.6) for the series Ti_{0.2}(Nb_{1−x}Ta_x)_{0.8}FeSb; Zhu et al. gained *ZT*~1.4 at 1200 K for $Hf_{0.12}Nb_{0.88}FeSb$ [\[168\]](#page-23-21) and for Ti_{0.16}Ta_{0.84}FeSb [\[188\]](#page-24-22).

Whereas many p-type half-Heuslers have *ZT*s < 1, the ZrCoBi-based half-Heuslers showed excellent thermoelectric performance, even more so for substituting at the Bi-site with $ZT = 1.42$ for $ZrCoBi_{0.65}Sh_{0.15}Sn_{0.20}$ at 973 K as reported by Zhu et al. [\[169\]](#page-23-23).

Also, for p-type half-Heusler alloys, some research groups tried to enhance *ZT* via additions. Kimura et al. [\[71\]](#page-20-7) added Co and Ir to ZrNiSn and converted this way the n-type to a p-type half-Heusler alloy; however, *ZT*s were rather low, i.e., below 0.3. Mallick et al. [\[159\]](#page-23-18) investigated $(Hf_{0.7}Ti_{0.3}CoSb_{0.8}Sn_{0.2})_{1-x}(Cu_{1.98}Ni_{0.04}Te_{0.97}Se_{0.03})_x$: for x = 0.25 they could enhance *ZT* = 0.15 (without the chalcogenide) to *ZT* = 0.63 at 1023 K, which corresponds an enhancement of 320%. Hsu et al. [\[103\]](#page-21-21) were successful in producing nanostructured $HfO₂$ during the ball-milling and spark–plasma–sintering process of $Zr_{0.5}Hf_{0.5}CoSb_{0.8}Sn_{0.2}$ and reached as highest *ZT* = 0.75. Xie et al. [\[193\]](#page-24-20) enhanced *ZT* of Ti(Co,Fe)Sb with the addition of 1 atom percent nano inclusions of InSb by 450%; however, *ZT* = 0.33 is still very low.

Not only for n-type half-Heusler alloys, as described in Section [3.2.1,](#page-11-1) but also for the p-type, NbFeSb and $Ti_{0.15}Nb_{0.85}FeSb$, the influence of severe plastic deformation via high-pressure torsion was investigated by Rogl and colleagues [\[25](#page-18-8)[,26\]](#page-18-25). For NbFeSb, the peak *ZT* could be enhanced by about 50% but is, with *ZT* = 0.00043, still very low. For $Ti_{0.15}Nb_{0.85}FeSb, ZT was enhanced from ZT = 0.68 to ZT = 0.74 for the thermally stable$ sample.

Generally, for n- as well as for p-type half-Heuslers, most *ZT*s at 300 K are below 0.3. Above 300 K, many peak *ZT*s are below 1 but can reach 1.6. Whereas the most successful system for the n-type is the (Ti,Zr,Hf)NiSn-system with Sb/Sn substitution, for the p-type, it is the (V,Nb,Ta)FeSb-system doped with Ti, Zr or Hf.

3.4. Miscellaneous

On 8-electron half-Heusler alloys: Several authors have screened 8-electron half-Heusler alloys in DFT throughput calculations, identifying a series of compounds of interest in thermoelectrics. Vikram et al., in one publication [\[219\]](#page-25-17), presented the TE properties of 21 half-Heusler alloys with a sum of 8 valence electrons (8-electron half-Heuslers); 6 of them showed *ZT*s between 0.8 and 1, such as LiSnAl, NaAlSn or LiAlGe. In another paper [\[220\]](#page-26-0), he calculated the thermoelectric properties of a large number of ternary compounds of the type: group IA or IB element—group IIA or IIB element—group V element. From these results, he highlighted even *ZT*s such as *ZT* = 1.37 for RbPBa and *ZT* = 1.56 for AgPMg. Hoat et al. [\[221\]](#page-26-1) also made theoretical investigations of the 8-electron half-Heusler alloy RbYSn, but they published only the power factors for the p- and n-types. So far, however, no experimental confirmation exists for the *ZT*s estimated from these DFT calculations.

On "double half-Heusler alloys", etc.: Some authors have coined the term "double half-Heusler alloys", "triple half-Heusler alloys", etc. These terms are not without problems. Firstly, double a half is full, but the authors certainly did not want to discuss full-Heusler alloys. Secondly, a formula such as $Ti₂FeCoSnSb$ from a structural chemical point of view indicates a fully ordered compound, where Fe,Co and Sn,Sb atoms fully occupy separate two-fold sites in a crystallographic subgroup of the space group of the half-Heusler phase. As such a subgroup cannot be found among the HH face-centered cubic space group type, the structure, therefore, needs to be transferred into a suitable body-centered tetragonal space group (cubic setting a,a,a needs to be transferred to a tetragonal cell: $a/\sqrt{2}$, $a/\sqrt{2}$, a) in order to find sites of two-fold multiplicity. Only if authors can prove the full atom order will the formula T₁₂FeCoSnSb be proper. Similar arguments hold for "triple half-Heusler alloys", etc.

4. Conclusions

The review presented here attempts to cover all measured thermoelectric properties of p- and n-type half-Heusler alloys of the last 25 years. Of course, even with the best literature search, it may be possible to miss some publications. It also became obvious that there exist about twice as many papers of n-type than of p-type half-Heusler alloys.

As the figure of merit, *ZT* is dependent on three key figures, the electrical resistivity, the Seebeck coefficient and the thermal conductivity, the evaluations in this review, focusing on these key figures, may help to understand the interrelations between them as well as their influence on *ZT*.

It was found out that the n-type (Ti,Zr,Hf)NiSn-sytem generally has the highest *ZT*s and that (Ti,Zr)NiSn, (Ti,Hf)NiSn, (Zr,Hf)NiSn and especially (Ti,Zr,Hf)NiSn have higher *ZT*s than TiNiSn, ZrNiSn and HfNiSn. Furthermore, this investigation showed that doping, in most cases, does not enhance *ZT* as much as substituting at the antimony site, preferably tin by antimony. Various additions of nanoparticles could, in some cases, push up *ZT* slightly.

For the p-type, the (V,Nb,Ta)FeSb-system in respect of high *ZT* is the best and even better with compounds doped with Ti or Hf, but also quite remarkable *ZT*s were found for the (Ti,Zr,Hf)CoSb-system substituted by tin, and for ZrCoBi substituted at the bismuth site by antimony and tin.

All these evaluated *ZT*s are, of course, primarily dependent on the preparation method warranting a proper microstructure with small (preferably nano) grain size and a high density.

The present compilation may provide the data for the individual selection of p- and n-type leg materials suitable for an efficient thermoelectric generator within a given temperature gradient.

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