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Generation of Geometric Extra Phase and Accompanying Temporal Effects in Asymmetric Optically Compensated IPS-LCDs and FFS-LCDs

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Abstract: This paper reports the generation of an extra phase and the accompanying temporal effects in an asymmetric optically compensated in-pane-switching (IPS) liquid crystal (LC) system and a fringe-field-switching (FFS) liquid crystal display (LCD) exhibiting a twofold faster response speed in the switching-off process compared with that in single and symmetric IPS-LCDs and FFS LCDs for the first time. To explain the experimental results, we derived an approximate analytical formula for the optical output intensity that includes an extra phase advancement and conducted simulations to achieve normally black operation using a dynamic optical retarder.

Keywords: IPS; FFS-LCD; asymmetric optical -LCD system; optical compensation; response time

1. Introduction

Liquid crystal displays (LCDs) are widely used as information displays for televisions, computer monitors, and various instruments. Among the various LCD operation modes, in-plane-switching (IPS) [1] and fringe-field-switching (FFS) [2] are currently widely used because of the excellent wide viewing angles of IPS and FFS displays. However, their response speed needs further improvement. To solve this problem, several methods have been adopted (1) using liquid crystal (LC) materials with low viscosity [3,4], (2) devising electrode structures [5,6], (3) doping with chiral agents [7], (4) doping with nanoparticles [7], (5) using ferroelectric LC materials [8,9], (6) using an LCD with a narrow gap cell [10], and (7) overdriving to enhance the response speed in the switching-on process [11].

Our approach reported in the present paper represents an alternative to the abovementioned approaches. In a preliminary previous work, we reported the electro-optical performances of asymmetric optically compensated tunable birefringent (TB) mode LCD exhibiting two-fold fast response speed in the switching off process [7]. For this reason, this present work was conducted aiming at the extension of the previous work to this proposed asymmetric optically compensated IPS and FFS LCDs (Asymmetric OC-IPS and FFS-LCD). The present work is related to the geometry phase, originally described by Pancharatnum [12] and Berry [13] in their historically representative papers. The present paper also updates previous work by introducing a dynamic optical retarder [7] for realizing electro-optical characteristics with by introducing a dynamic retarder, we realized normally black state normally black operation, accompanying with a simulation utilizing a simulator, SHINTECH LCD MASTER. In conjunction with a simulation using a SHINTECH LCD MASTER simulator, we introduce a dynamic retarder, to realize normally black operation.



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2. Theory

2.1. Optical Configuration of our Asymmetric IPS-LC System

In Figure 1, angle φ_1 is the switching azimuthal angle of LC molecules and φ_2 is the setting azimuthal angle of a compensator (+A Plate) whose angle is set to $\varphi_2 = 3\pi/4 - \alpha$, which can be controlled using a dynamic retarder, when we use an IPS cell for the compensator.



Figure 1. Optical configuration of our asymmetric IPS-LCD (FFS-LCD) system.

The whole system is sandwiched between crossed polarizers, where the incident light field into the LCD system is polarized in the x-direction and the transmitted wave is polarized in the y-direction. Labels e_1 and o_1 in Figure 1 indicate that the axes of the extra-ordinary wave and the ordinary wave, respectively, and labels e_2 and o_2 indicate those of the compensator, respectively.

When $\varphi_2 = 3\pi/4$ and $\varphi_1 = \pi/4$, the system is symmetric and generates a black state; under these conditions, the extra phase is not generated and the high response speed is not achieved.

2.2. Analytical Derivation of Normalized Optical Transmission of a Single IPS-Cell

Using the system shown in Figure 1, we here carry out a 2 \times 2 Jones matrix calculation [14] for I_1 of a single cell.

The input electrical field is

$$\left[\begin{array}{c}1\\0\end{array}\right]E$$

We then have a projected field in the *x*-*y* coordinate:

$$\begin{bmatrix} E_1(x) \\ E_1(y) \end{bmatrix} = \begin{bmatrix} c_1 & -s_1 \\ s_1 & c_1 \end{bmatrix} \begin{bmatrix} -G_{11} & 0 \\ 0 & G_{12} \end{bmatrix} \begin{bmatrix} c_1 & s_1 \\ -s_1 & c_1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} E_0 \sin(\omega t)$$
(1)

where c_1 and s_1 are $\cos(\varphi_1)$ and $\sin(\varphi_1)$, respectively, here $G_{11} = \exp(-i \Gamma_{11})$ and $G_{12} = \exp(-i \Gamma_{12})$; in addition, $\delta_1 = \Gamma_{11} - \Gamma_{12}$, where each of which is the phase delay in the e_1 and o_1 , from the incident wave, respectively.

We have the *x*-*y* components such that:

$$E_1(x) = [(c_1)^2 \exp(-i\delta_1) + (s_1)^2]E_0,$$
⁽²⁾

$$E_1(y) = [s_1 c_1(\exp(-i\delta_1) - 1)]E_0,$$
(3)

Then,

$$E_1(y) = E_0 \sin(\varphi_1) \cos(\varphi_1) \exp(-i\delta_1).$$
(4)

$$= 2E_0 \sin(2\varphi_1)\exp(-i\delta_1). \tag{5}$$

Thus, we have the optical output intensities and the normalized optical transmission for the single cell:

$$I_1 = I_0 \sin^2(2\varphi_1) \sin^2(\delta_1/2).$$
(6)

The normalized transmission is then

$$\Gamma_1 = \sin^2(2\varphi_1). \tag{7}$$

A δ_1 = 1.3 π for n-IPS and n, p-FFS and 1.04 π for p-IPS were determined based on the properties of the LC materials and the LCD cell fabrication details (Section 3.3).

2.3. Derivation of Normalized Optical Transmission Intensity of OC-IPS with 2 \times 2 Jones Matrix Calculations

We here derive the equation for the proposed asymmetric OC-IPS and FFS shown in Figure 1. This derivation is carried out by replacing the c_1 and s_1 terms in Equation (1) with c_2 and s_2 and $\Gamma_{21} - \Gamma_{22}$. In addition, by inserting $E_1(x)$ and $E_1(y)$, we obtain

$$E_1(x) = (s_1)^2 - (c_1)^2$$
 and $E_1(y) = -2s_1c_1$

We derive $E_2(y)$ as follows:

$$\begin{bmatrix} E_2(x) \\ E_2(y) \end{bmatrix} = \begin{bmatrix} c_2 & -s_2 \\ s_2 & c_2 \end{bmatrix} \begin{bmatrix} -G_{21} & 0 \\ 0 & G_{22} \end{bmatrix} \begin{bmatrix} c_2 & s_2 \\ -s_2 & c_2 \end{bmatrix} \begin{bmatrix} E_1(x) \\ E_1(y) \end{bmatrix}$$
(8)

where δ_2 is again the difference in phase delay between the e_2 and o_2 axes:

$$\delta_2 = \Gamma_{21} - \Gamma_{22}.$$

We then have:

$$E_2(y) = 2[s_2c_2(c_1^2 - s_1^2) - s_1c_1(c_2^2 - s_2^2)] \exp(-i\delta_2).$$
(9)

And $\delta_2 = \delta_1 = 1.3\pi$ for n-IPS and n, p-FFS and 1.04π for p-IPS as evaluated based on the physical properties and the LCD cell fabrication conditions described in Section 3.3. Then, with $\sin^2(\omega t) = 1/2$ and without $\exp(-i\delta_2)$ to normalize the optical intensity and transmission.

$$I_2 = I_0 \sin^2(2\varphi_2 - 2\varphi_1) \sin^2(\delta_1/2) \sin^2(\delta_2/2)$$
(10)

$$T_2 = \sin^2 \left(2\varphi_2 - 2\varphi_1 \right), \tag{11}$$

If we substitute $\varphi_2 = \pi/4 - \alpha$ into Equation (11), we obtain an equation with a phase-advancement shift, of 2α :

$$T_2 = I_0 \cos^2(2\varphi_1 + 2\alpha)$$
(12)

In our system, the symmetry breaking introduces a phase shift, which leads to a fast response in the switching-off process. This effect is shown later in Figure 3. When LC-cell and the A-plate are crossed, $\varphi_1 = \varphi_2 = \pi/4$ and $\delta_1 = \delta_2 = 1.3\pi$, the symmetric system then produces the black (dark) state and no ultrafast switching process occurs.

3. Results and Discussion

3.1. Optical Configurations and Switching of Single and OC-IPS, and OC-FFS Cells with $\Delta \varepsilon > 0$ and $\Delta \varepsilon < 0$

Figure 2a,b show the optical configurations and switching of single and OC-IPS and OC-FFS-LCD cells used in the present research, where p and n denote $\Delta \varepsilon > 0$ and $\Delta \varepsilon < 0$, respectively.



(a) Switching of p-IPS and p-FFS



(b) Switching of n-IPS and n-FFS

Figure 2. Optical configuration of switching in single and optically compensated IPS- and FFS-LCDs with negative dielectric anisotropy.

We next explain the switching process in our OC-LCDs, along with the contents of Figure 2.

Figure 3 shows simulation results for the normalized transmission and compares T_1 (I_1) with T_2 (I_2) against the switching angles calculated using Equations (7) and (12). The top and central parts in this figure show the behaviors in the switching-off process, and the bottom and right parts show the behaviors in the switching-on process. Thus, the φ_1 for T_2 (I_2) starts from a degree of $-\alpha$, and the φ_1 for the T_1 (I_1) starts from -45° .

The results in Figure 3 are interpreted as follows (1) The $T_2(I_2)$ has a phase advancement of α over the $I_1(T_1)$, (2) According to the top area, the $T_2(I_2)$ starts the decay process with a finite inclination, whereas the $T_1(I_1)$ starts the decay process with no inclination. This decay in the absence of inclination is attributed to the fast response in the decay process, which was experimentally demonstrated in the present research. This phenomenon is a purely optical effect, where the decay process occurs as the common molecular relaxation for both the symmetric and asymmetric systems. (3) The $I_2(T_2)$ has a finite value at the right bottom because of energy conservation, which means that the operation is not normally black. The realization of a normally black operation will be described in Section 3.4.



Figure 3. Simulation results of a comparison of switching characteristics of I_1 (T_1) with I_2 (T_2) starting from the switching off process.

3.2. V-T Curves of Single p-IPS and p-OC-IPS Devices

Figure 4 shows the experimentally obtained transmittance-voltage (V-T)-curves of IPS LCDs, as recorded at 25 °C using a DMS-703 (Autronic Melchers, GMb) instrument. Interestingly, in the symmetric compensation system, the introduction of optical compensation introduces a wide operating voltage range, as shown in Figure 4. Regarding the V-T curve, refer to [14], where the Freedericksz transition is described.



Figure 4. Experimentally obtained *V*-*T* curves for p-IPS ($\Delta \varepsilon > 0$) and p-OC-IPS devices with $\alpha = 0$ (symmetric).

Figure 5 shows the *V*-*T* curves of the p-OC-IPS device without optical compensation ($\alpha = 0$) and the p-OC-IPS device with optical compensation, ($\alpha = -6.8^{\circ}$). The proposed asymmetric OC-IP produces *T* = 0 at 4 V; which means a non-normally black operation. In general, an LCD is operated normally black or normally white; thus, a dynamic optical compensator, which has been updated from the device described in our previous work [7], was used (see Section 3.5).

3.3. Experimental Results of Response Time

The physical properties of the LC materials and the specifications of the LCD cells used in the present work are described as follows: The cell gap of the sample cells was $3.5 \ \mu m$ and the wavelength of the optical light source was 589 nm. The anisotropy of refractive index,

 Δn (at 589 nm, 25 °C) and $\Delta \varepsilon$ (25 °C) (After Datasheet of DIC)

0.088	6.2
0.110	-4.1
0.110	11.9
0.110	-4.1
	0.088 0.110 0.110 0.110

Using these quantities,

The experimental results in Tables 1 and 2, indicate that the response time of the proposed asymmetric OC-IPS and OC-FFS devices are reduced by 42–56% compared with those of single cells and symmetric OC-IPS and OC-FFS devices. Approximately the same results were obtained for the OC-FFS device, and we will report the corresponding data elsewhere. These results indicate that the response time in the switching-on process will be greatly reduced by adopting the overdriving technique [12].



Figure 5. Experimentally obtained *V*-*T* curves of a p-OC-IPS device without optical compensation, ($\alpha = 0$) and a p-OC-PS device with optical compensation, ($\alpha = -6.8^{\circ}$).

Table 1. Results of res	ponse time for	the n-OC-IPS	device.
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	Single n-IPS	n-OC-IPS $(\alpha = 0^{\circ})$	$n-OC-IPS \\ (\alpha = -6.8^{\circ})$
$ au_{ m off} \ ({ m msec}) \ V_{ m op} \ ({ m V})$	25.6	25.1	14.9 (42% decrease)
	6.8	9.1	15.0

Table 2. Results of response time for the p-OC-IPS device.

Single p-IPS	$p-OC-IPS \\ (\alpha = 0^{\circ})$	p-OC-IPS $(\alpha = -6.8^{\circ})$
18.3	14.1	8.1 (56% decrease)
4.0	4.5	5.5

3.4. Temporal Derivatives of I₁ and I₂ and Their Comparison

We here derive temporal derivatives of I_1 and I_2

$$\frac{\partial I_1}{\partial t} = \frac{\partial I_1}{\partial \varphi_1} \frac{\partial \varphi_1}{\partial t} = 2I_0 \sin(4\varphi_1) \frac{\partial \varphi_1}{\partial t}$$
(13)

For an OC-IPS:

$$\frac{\partial I_2}{\partial t} = \frac{\partial I_2}{\partial \varphi_1} \frac{\partial \varphi_1}{\partial t} = -2I_0 \sin(4\varphi_2 - 4\varphi_1) \frac{\partial \varphi_1}{\partial t} = -2I_0 \sin(4\varphi_1) [\cos(4\alpha) - \sin(4\alpha) \cot(4\varphi_1)] \frac{\partial \varphi_1}{\partial t}$$
(14)

Thus, we obtain a formula, the *F*-function, by taking the ratio between Equations (13) and (14) as follows, where the two minus signs cancel each other:

$$F = \cos(4\alpha) - \sin(4\alpha)\cot(4\varphi_1)$$

The *F*-function $\left(\frac{\partial I_2}{\partial t} / \frac{\partial I_1}{\partial t}\right)$ is shown in Figure 6.



Figure 6. The *F*-functions with $\alpha = -0.12$ rad and $\alpha = -0.05$ rad.

Figure 6 shows that the temporal derivative of I_2 is much larger than that of I_1 during the switching-off process, consistent with the experimental results obtained in the present work. However, it also shows that this effect terminates at a switching angle of 15° and that the relation is thereafter inverted. However, if necessary, a dynamic compensator can be used to eliminate the phase-advancement after the decay process for the switching-on process, the removal of the applied voltage causes an overvoltage, resolving this problem [11].

3.5. Dynamic Compensator

A dynamic compensator changes the azimuthal angle of the *e*-axis, $\varphi_2 = 3\pi/4 - \alpha$, by applying a computer-controlled electric voltage, V_{appl} , according to the following equation:

$$\alpha = \alpha_0 \exp\left[1 - kV_{\rm th} / V_{\rm appl}\right], k > 1 \ (k = 2.5) \tag{15}$$

- (1) For the switching-on operation (starting from $V_{appl} \approx 0$), normally black operation will be induced with $\alpha = 0$ ($\varphi_2 = 3\pi/4$).
- (2) For the switching-off operation, where $V_{appl} = kV_{th}$ (k = 2.5), then $\alpha = \alpha_0$ (~7°), extra phase, α , is generated.

This practice is demonstrated in Figure 7.



Figure 7. Demonstration of normally black mode operation using a dynamic compensator.

4. Conclusions

- 1. The optical output intensities were formulated using a slightly modified Jones matrix calculation, and the results showed that an extra phase accompanying temporal behavior is generated in the switching angle for the proposed asymmetric optically compensated IPS -LCD and FFS-LCD systems.
- 2. Experiments showed that the optically compensated IPS-LCD has a broad dynamic range beyond the sine-squared oscillation in the ordinary IPS-LCD.
- 3. Simulation results revealed the existence of phase advancement in the optical output intensity of the proposed asymmetric OC-IPS relative to that of symmetric OC-IPS, depending on the selected optical retardation in the compensator.
- 4. Experimentally measured values of the response time of the proposed asymmetric OC-IPS-LCD and OC-FFS were reduced by 42–56% compared with those of symmetric OC-IPS, and OC-FFS, LCDs, and single IPS and FFS-LCDs.
- 5. Temporal derivatives of the optical intensity for the proposed asymmetric OC-IPS were shown to be larger than those of the symmetric OC-IPS.
- 6. By adopting a dynamic compensator, we demonstrated normally black operation for the proposed asymmetric OC-IPS.

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