

## Article

# A Helmholtz Free Energy Equation of State of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> Fluid Mixtures (ZMS EOS) and Its Applications

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**Abstract:** An equation of state (EOS) of CH<sub>4</sub>-N<sub>2</sub> fluid mixtures in terms of Helmholtz free energy has been developed by using four mixing parameters, which can reproduce the pressure-volume-temperature-composition (*PVTx*) and vapor-liquid equilibrium (VLE) properties of CH<sub>4</sub>-N<sub>2</sub> fluid mixtures. The average absolute deviation of all the *PVTx* data available up to 673.15 K and 1380 bar from this EOS is 0.38%. Combining this EOS of CH<sub>4</sub>-N<sub>2</sub> fluid mixtures and the EOS of CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> fluid mixtures in our previous work, an EOS of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures has been developed, which is named ZMS EOS. The ZMS EOS can calculate all thermodynamic properties of ternary CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures and the average absolute deviation of the *PVTx* data from the ZMS EOS is 0.40% for the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> system. The ZMS EOS can be applied to calculate excess enthalpies of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures, predict the solubility of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures in brine and water, and quantitatively estimate the impact of the impurities (CH<sub>4</sub> and N<sub>2</sub>) on the CO<sub>2</sub> storage capacity in the CO<sub>2</sub> capture and storage (CCS) processes. The ZMS EOS can also be applied to calculate the isochores of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> system in the studies of fluid inclusions. All Fortran computer codes and Origin drawing projects in this paper can be obtained freely from the corresponding author.

**Keywords:** CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> system; equation of state; CCS; excess enthalpies; fluid inclusion



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## 1. Introduction

The CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub> are important natural fluids, and non-aqueous CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixtures are often reported in the studies of fluid inclusions [1–3]. In recent decades, the amount of greenhouse gas CO<sub>2</sub> in the atmosphere has gradually increased because of the development of industry. Numerous studies have shown that CO<sub>2</sub> capture and storage (CCS) is an effective method to reduce the amount of CO<sub>2</sub> in the atmosphere [4–6]. However, CO<sub>2</sub> in industrial exhaust gases is usually impure, mixing with other components, such as CH<sub>4</sub> and N<sub>2</sub> [7,8]. The impurities can affect the design of the CCS processes [9–11]. Therefore, predicting thermodynamic properties of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures, especially the pressure-volume-temperature-composition (*PVTx*) and vapor-liquid equilibrium (VLE) properties at different temperatures and pressures, is of great significance for the related CCS and fluid inclusion studies [12–15].

Although experimental thermodynamic data are the most reliable in the corresponding applications, they are costly and time-consuming to obtain. On the other hand, the predictive EOS is a better choice for us. In the past century, the cubic EOSs (e.g., Peng-Robinson (PR) EOS [16], Soave-Redlich-Kwong (SRK) EOS [17]) originated from the van der Waals EOS [18] are highly preferred for the CCS researchers because of their simplicity. Based on the standard PR EOS, some more accurate EOSs have been presented for the CCS mixtures, such as a model that integrates the classical PR EOS with advanced mixing rules, called the “Peng-Robinson + residual part of excess Helmholtz energy model” [19], and the Enhanced-Predictive-PR78 (E-PPR78) EOS [20]. However, the cubic EOSs cannot well reproduce the *PVTx* properties of fluids under high pressure-temperature conditions, as can be seen in

Section 2. Based on Heyen EOS [21], Heyen et al. [22] simulated the phase equilibria of the CH<sub>4</sub>-CO<sub>2</sub> system below 50 °C and 100 bar; Darimont and Heyen [23] predicted the phase equilibria of the CO<sub>2</sub>-N<sub>2</sub> system below 20 °C and 90 bar. Because cubic EOSs cannot well reproduce the *PVTx* properties, Thiery et al. [24,25] and Thiery and Dubessy [26] modeled the phase equilibria of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> system including liquid, vapor and solid by the cubic EOSs and the *PVTx* properties by the Lee–Kesler correlation [27]. However, these models have the disadvantage of using two EOSs that may cause incoherency of various fluid parameters [28].

In recent years, the Statistical Associating Fluid Theory (SAFT) EOS [29] from Wertheim’s thermodynamic perturbation theory (TPT) has become popular for CCS fluid mixtures. Because it is based on statistical mechanics and has an explicit physical meaning, it has more reliable extrapolation and more accurate predictive capability than traditional empirical models. Many successful modifications of the SAFT-based EOSs have been presented, such as the EOS of perturbed-chain SAFT (PC-SAFT) [30,31], variable range SAFT (SAFT-VR) [32], and Lennard-Jones SAFT (LJ-SAFT) [33]. However, the SAFT EOSs cannot well reproduce experimental VLE data in the near-critical region of the CO<sub>2</sub>-N<sub>2</sub> fluid mixture, as can be seen in the previous work [34]. Up to now, the EOS in terms of Helmholtz free energy is also commonly used to calculate the thermodynamic properties of CCS fluid mixtures [35–39], because the EOS in the form of Helmholtz free energy has high precision, wide temperature-pressure range, and can calculate all thermodynamic properties of the fluids by the derivation from the Helmholtz free energy EOS, such as enthalpy, heat capacity, entropy, etc. The derivation from the Helmholtz free energy EOS is not complex, as can be seen later.

In our previous work [34,40], the EOS in terms of dimensionless Helmholtz free energy for the CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> fluid mixtures have been established by using four mixing parameters, which have been applied to the studies of the CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> fluid inclusions. In this work, firstly, the EOS in terms of dimensionless Helmholtz free energy of the CH<sub>4</sub>-N<sub>2</sub> fluid mixture has been developed by the same approach. The ZMS EOS is obtained by combining the binary interaction parameters of CH<sub>4</sub>-CO<sub>2</sub>, CO<sub>2</sub>-N<sub>2</sub> and CH<sub>4</sub>-N<sub>2</sub> systems. Experimental data available for the binary CH<sub>4</sub>-N<sub>2</sub> and ternary CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures are used to verify the accuracy of the ZMS EOS. Then the ZMS EOS is applied to calculate the excess enthalpies, the solubility of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> gas mixtures in brine and water, the impact of impurities (CH<sub>4</sub> and N<sub>2</sub>) on the CO<sub>2</sub> storage capacity, and the isochores of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid inclusions.

## 2. The ZMS EOS

The EOS of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixtures is in terms of dimensionless Helmholtz free energy  $\alpha$ , which is represented by

$$\alpha = \alpha^0 + \alpha^r \quad (1)$$

where  $\alpha^0$  and  $\alpha^r$  are the ideal-gas part and the residual part of dimensionless Helmholtz free energy.  $\alpha^0$  and  $\alpha^r$  are defined by

$$\alpha^0 = \sum_{i=1}^n x_i [\alpha_i^0 + \ln x_i] \quad (2)$$

$$\alpha^r = \sum_{i=1}^n x_i \alpha_i^r(\delta, \tau) + \alpha^E(\delta, \tau, x) \quad (3)$$

where  $n$  is the number of components in the mixture,  $x_i$  is the mole fraction of component  $i$  in the mixture, the superscripts “0”, “r” and “E” denote the ideal-gas part, the residual part, and the excess of dimensionless Helmholtz free energy, respectively. The subscript  $i$  denotes the component  $i$ .  $\alpha_i^r$  can be calculated by the EOSs of pure CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub> fluids [41–43], which are in terms of dimensionless Helmholtz free energy and are recommended as the standard EOSs of pure CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub> fluids by the National Institute of Standards and Technology (NIST).

$\delta$  and  $\tau$  of Equation (3) are the reduced mixture density ( $\delta = \rho/\rho_c$ ) and the inverse reduced mixture temperature ( $\tau = T_c/T$ ), where  $\rho$  and  $T$  are the density and temperature of the mixture, and  $\rho_c$  and  $T_c$  are the pseudo-critical density and temperature of the mixture.  $\rho_c$  and  $T_c$  are defined by

$$\rho_c = \left[ \sum_{i=1}^n \frac{x_i}{\rho_{ci}} + \sum_{i=1}^{n-1} \sum_{j=i+1}^n x_i x_j \zeta_{ij} \right]^{-1} \tag{4}$$

$$T_c = \sum_{i=1}^n x_i T_{ci} + \sum_{i=1}^{n-1} \sum_{j=i+1}^n x_i^{\beta_{ij}} x_j \zeta_{ij} \tag{5}$$

where  $x_j$  is the mole fraction of component  $j$  in the mixture,  $\rho_{ci}$  and  $T_{ci}$  are the critical density and temperature of component  $i$ , respectively. The critical temperatures and densities of the three components considered in this study are listed in Table 1.  $\zeta_{ij}$ ,  $\beta_{ij}$  and  $\zeta_{ij}$  are binary parameters associated with components  $i$  and  $j$ .

**Table 1.** Critical parameters of pure fluids.

Component i	$T_{ci}$ (K)	$\rho_{ci}$ (mol/dm <sup>3</sup> )	References
CH <sub>4</sub>	90.6941	10.139	Setzmann and Wagner [42]
CO <sub>2</sub>	216.592	10.625	Span and Wagner [41]
N <sub>2</sub>	63.151	11.184	Span et al. [43]

$\alpha^E$  of Equation (3) takes the general form developed by Lemmon and Jacobsen (LJ-1999 EOS) [35].

$$\alpha^E(\delta, \tau, x) = \sum_{i=1}^{n-1} \sum_{j=i+1}^n x_i x_j F_{ij} \sum_{k=1}^{10} N_k \delta^{d_k} \tau^{t_k} \tag{6}$$

where  $F_{ij}$  is a binary parameter associated with components  $i$  and  $j$ ,  $N_k$ ,  $d_k$  and  $t_k$  are the general parameters independent of fluids, which can be obtained from Lemmon and Jacobsen [35]. There are four binary parameters ( $\zeta_{ij}$ ,  $\beta_{ij}$ ,  $\zeta_{ij}$  and  $F_{ij}$ ) in the above equations. Binary interaction parameters of CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> have been determined in previous work [34,40]. In this work, the binary interaction parameters of CH<sub>4</sub>-N<sub>2</sub> systems are obtained by the nonlinear regression of selected experimental data. Since experimental  $PVTx$  data available are much more than experimental VLE data, only part of new and accurate experimental  $PVTx$  data are selected to fit the parameters to make the  $PVTx$  and VLE data keep similar weights in the fitting. In this work, the Levenberg–Marquardt algorithm of the nonlinear least squares method is used for fitting, which is an efficient and widely used mathematical optimization technique. The fitting condition is to minimize the weighted sum of squares of the calculation errors of the equation on the selected experimental  $PVTx$  and VLE data. Compared to Newton’s method (e.g., The Newton–Raphson method used in REFPROP [44]), it combines the advantages of two numerical minimization algorithms: the gradient descent method and the Gauss–Newton method, and it has good convergence and robustness. Regressed parameters of the CH<sub>4</sub>-N<sub>2</sub> system are listed in Table 2, which also includes the parameters of CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> systems from previous studies.

**Table 2.** Binary interaction parameters of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> system.

References	Binary Mixtures	$F_{ij}$	$\zeta_{ij}$ (dm <sup>3</sup> /mol)	$\zeta_{ij}$ (K)	$\beta_{ij}$
Mao et al. [40]	CH <sub>4</sub> -CO <sub>2</sub>	$0.12844025 \times 10$	$0.35751245 \times 10^{-2}$	$-0.43720344 \times 10^2$	$0.10358865 \times 10$
Zhang et al. [34]	CO <sub>2</sub> -N <sub>2</sub>	$0.16671494 \times 10$	$0.58411078 \times 10^{-2}$	$-0.22952094 \times 10^2$	$0.16878787 \times 10$
This work	CH <sub>4</sub> -N <sub>2</sub>	0.63739997	$0.3812517 \times 10^{-2}$	$-0.17790001 \times 10^2$	$0.1009001 \times 10$

Note: Subscripts  $i$  and  $j$  refer to the first component and the second component, respectively.

### 2.1. The Binary CH<sub>4</sub>-N<sub>2</sub> Mixture

The calculated deviations of the *PVTx* and VLE properties from each experimental data set [45–66] by this EOS for the CH<sub>4</sub>-N<sub>2</sub> fluid mixture are listed in Table 3. Detail calculation method for the *PVTx* and VLE properties can be found in Mao et al. [40]. The deviations of the *PVTx* data are the density deviations, and the temperature and pressure of the experimental *PVTx* data are up to 673.15 K and 1380 bar, respectively. The deviations of VLE property are the deviations of the vapor-liquid phase equilibrium composition, and the experimental VLE composition covers the full range. The average absolute deviation of all *PVTx* data from this EOS is 0.38% for the CH<sub>4</sub>-N<sub>2</sub> fluid mixture.

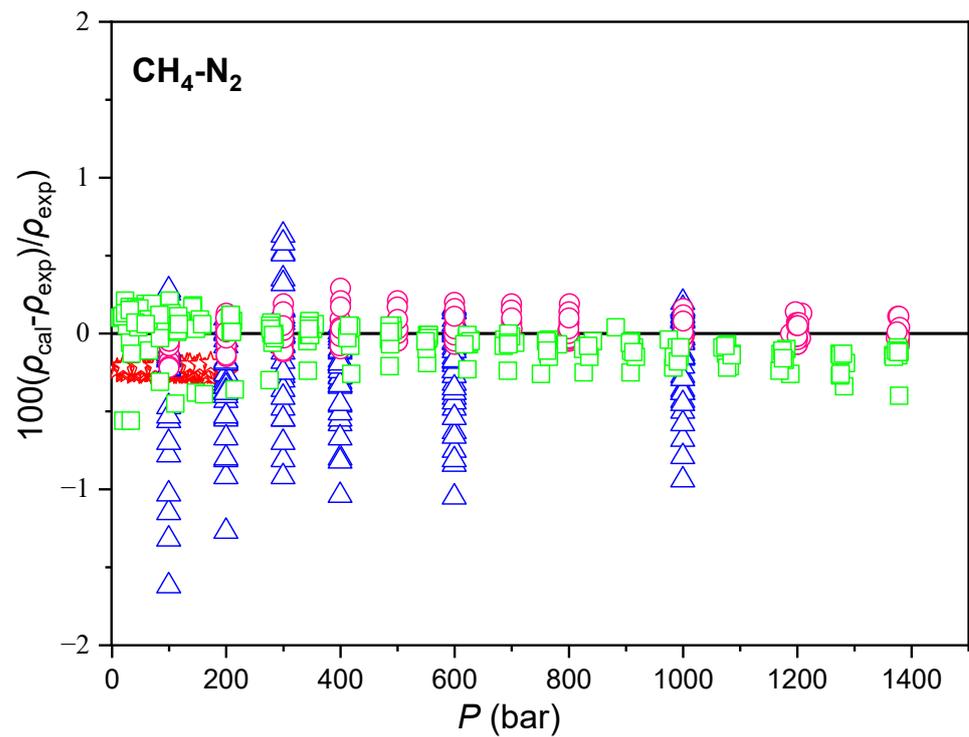
**Table 3.** Calculated *PVTx* and VLE deviations for the CH<sub>4</sub>-N<sub>2</sub> mixture.

References	Number of Data Points		T-P-x Range for xCH <sub>4</sub> -(1-x) N <sub>2</sub>			AAD%
	Styles	Total (Used)	T (K)	P (bar)	x	
Liu and Miller [45]	<i>PVTx</i>	7 (0)	91–115	3–12	0.5	0.57
Rodosevich and Miller [46]	<i>PVTx</i>	8 (0)	91–115	43–454	0.8–0.95	3.83
Pan et al. [47]	<i>PVTx</i>	7 (0)	91–115	1–11	0.5–0.86	0.11
Hiza et al. [48]	<i>PVTx</i>	21 (0)	95–140	1–21	0.5–0.95	0.10
Da Ponte et al. [49]	<i>PVTx</i>	182 (50)	110–120	15–1380	0.3–0.68	0.12
Straty and Diller [65]	<i>PVTx</i>	578 (0)	82–320	6–356	0.32–0.7	0.24
Haynes and McCarty [56]	<i>PVTx</i>	85 (85)	140–320	10–164	0.3–0.71	0.09
Seitz et al. [64]	<i>PVTx</i>	190 (90)	323–573	99–999	0.1–0.9	0.28
Seitz and Blencoe [63]	<i>PVTx</i>	43 (0)	673.15	199–999	0.1–0.9	5.36
Ababio et al. [50]	<i>PVTx</i>	83 (83)	308–333	9–120	0.5–0.78	0.12
Chamorro et al. [52]	<i>PVTx</i>	242 (56)	230–400	9–192	0.8–0.9	0.23
Janisch et al. [57]	<i>PVTx</i>	17 (0)	129–180	15–50	0.4–0.9	2.10
Li et al. [60]	<i>PVTx</i>	27 (27)	17–270	1–16	0.9	0.06
Gomez-Osorio et al. [54]	<i>PVTx</i>	142 (42)	304–470	100–1379	0.25–0.75	0.07
Brandt and Stroud [51]	VLE	23 (0)	128–179	34	0.05–0.98	1.48
Cheung and Wang [53]	VLE	20 (0)	92–124	0.2–6	0.85–1.0	5.02
Pan et al. [47]	VLE	60 (60)	95–120	0.2–25	0.05–1	5.05
Miller et al. [62]	VLE	11 (0)	112	2–13	0.2–0.97	3.30
Kidnay et al. [59]	VLE	83 (83)	112–180	2–49	0.1–0.99	1.31
McClure et al. [61]	VLE	8 (8)	91	1–3	0.1–0.8	5.60
Jin et al. [58]	VLE	10 (10)	123	4–26	0.1–0.95	2.96
Parrish and Hiza [66]	VLE	48 (43)	95–120	2–20	0.1–0.9	3.14
Janisch et al. [57]	VLE	16 (6)	130–180	0.5–5	0.4–0.96	1.14
Han et al. [55]	VLE	77 (60)	110–123	4–13	0.7–1.0	2.49

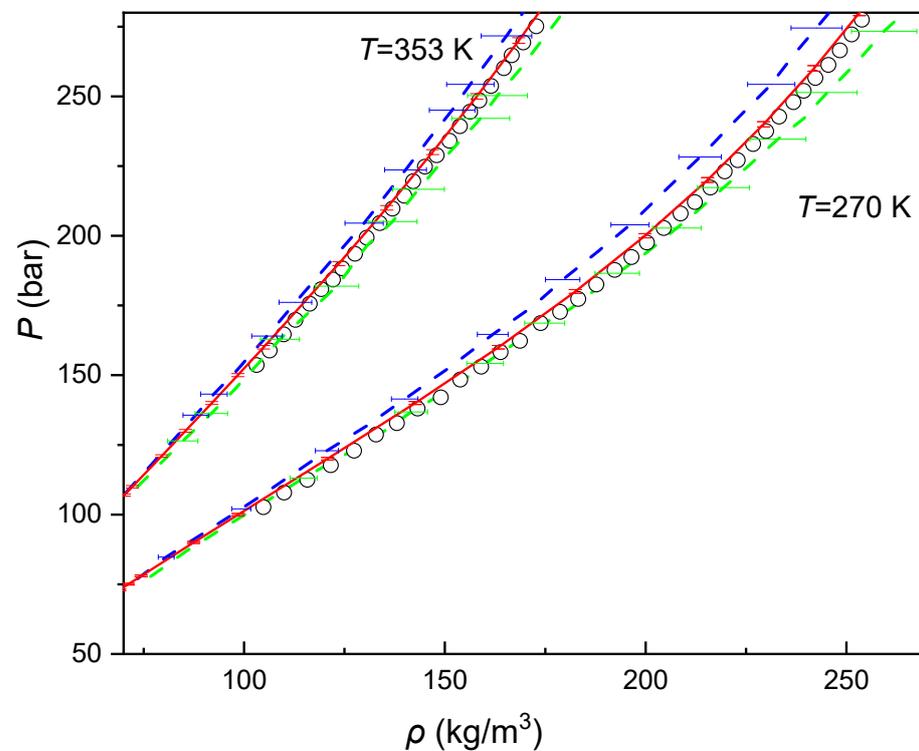
Note: Deviations of the *PVTx* data are the density deviations, and the deviations of VLE data are the equilibrium composition deviations. AAD is the abbreviation of average absolute deviation.

Figure 1 shows the deviations between this EOS and the experimental density data of the CH<sub>4</sub>-N<sub>2</sub> mixture [49,52,54,64] at different pressures. It can be seen in Figure 1 that most of the density deviations are within  $\pm 1\%$ , with or close to experimental accuracy. The comparisons between experimental density and the calculated densities by this EOS are shown in Figure 2, where the calculated densities from the PR and SRK EOSs [67] are also included. The errors of three EOSs are also plotted in Figure 2. Figure 2 shows the PR and SRK EOSs cannot well reproduce the *PVTx* properties, especially under high-pressure conditions. In contrast, the calculated densities from this EOS are in good agreement with experiment data [68].

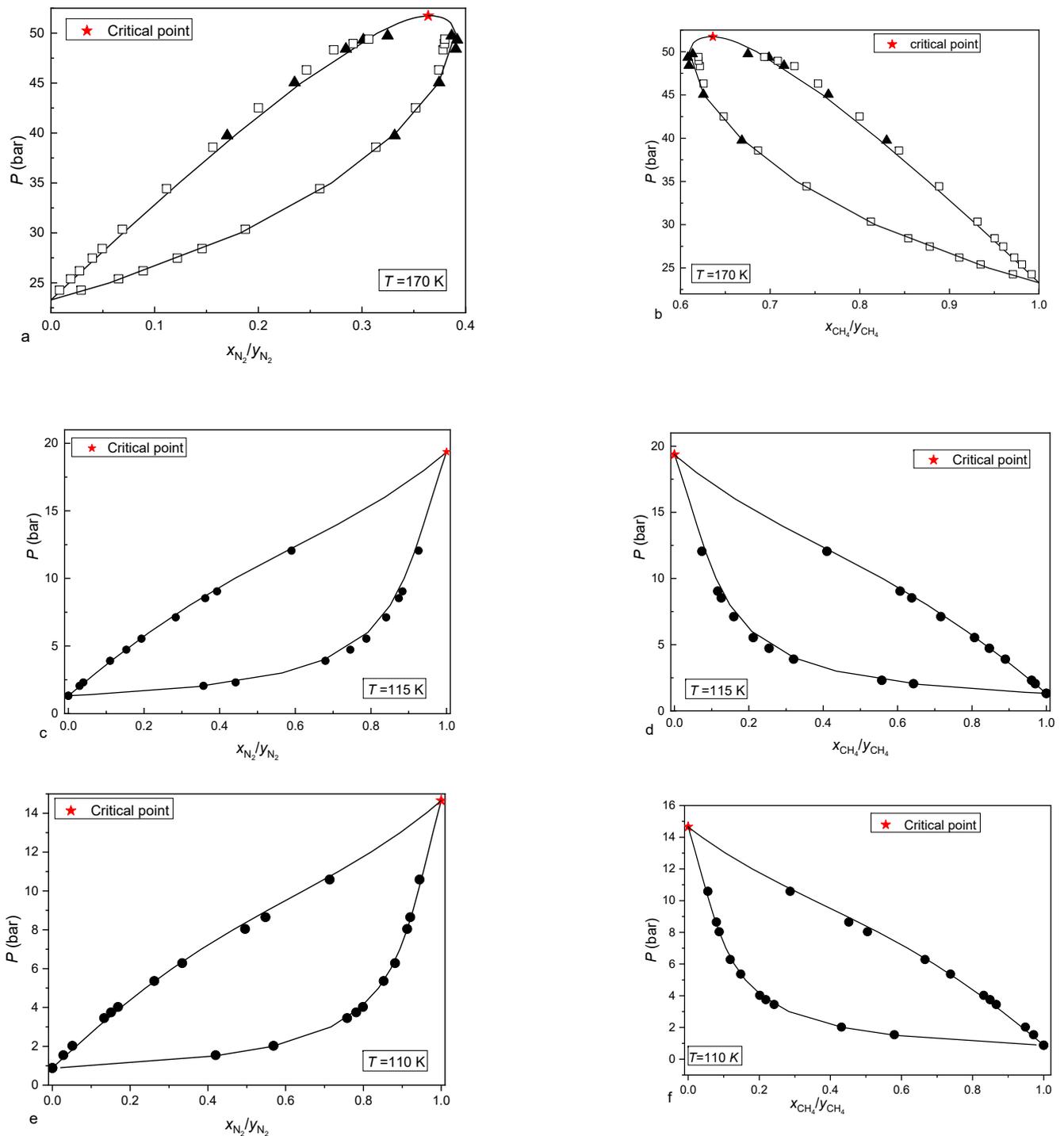
Figure 3 compares the vapor-liquid phase equilibrium curves calculated from this EOS with experimental data [55,57,59,66] at three temperatures of 170 K, 115 K and 110 K. It can be seen from Figure 3 that the calculated vapor-liquid phase equilibrium curves from this EOS are consistent with experimental data, indicating that this EOS can accurately calculate the vapor-liquid phase equilibrium of CH<sub>4</sub>-N<sub>2</sub> mixture.



**Figure 1.** Density deviations for the CH<sub>4</sub>-N<sub>2</sub> system at different pressures. Experimental data (triangles [64], stars [52], rounds [54], squares [49]):  $P$  is the pressure,  $\rho_{\text{cal}}$  is the calculated density, and  $\rho_{\text{exp}}$  is the experimental density.



**Figure 2.** Experimental and calculated densities for the 0.8 mole CH<sub>4</sub> + 0.2 mole N<sub>2</sub> mixture at different temperatures. The experimental density data (rounds) [68], the PR (green dash lines) and SRK (blue dash lines) EOSs [67], this work (red lines):  $P$  is the pressure,  $\rho$  is the density, and  $T$  is temperature.

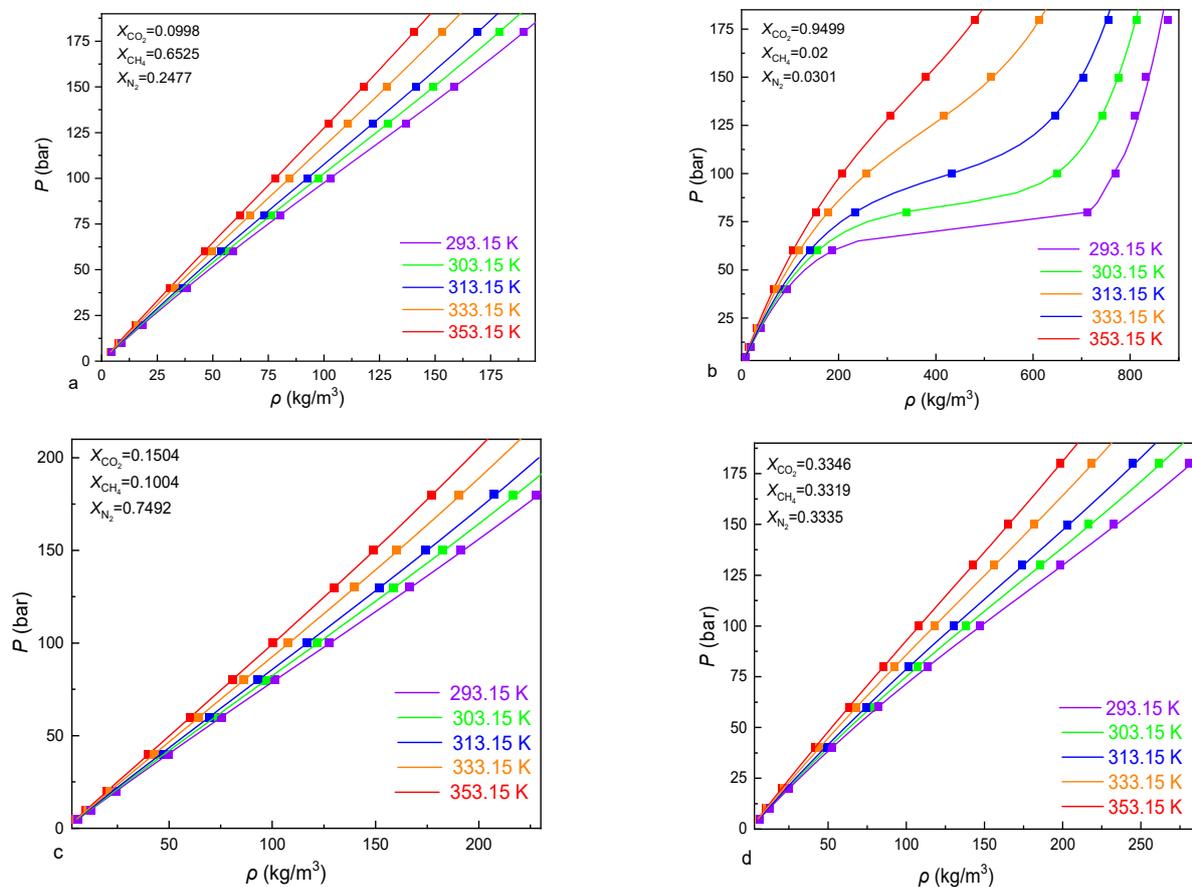


**Figure 3.** Vapor-liquid phase equilibria of the  $\text{CH}_4\text{-N}_2$  system at different temperatures, (a)  $P\text{-}x_{\text{N}_2}/y_{\text{N}_2}$  figure at 170 K; (b)  $P\text{-}x_{\text{CH}_4}/y_{\text{CH}_4}$  figure for 170 K; (c)  $P\text{-}x_{\text{N}_2}/y_{\text{N}_2}$  figure for 115 K; (d)  $P\text{-}x_{\text{CH}_4}/y_{\text{CH}_4}$  figure for 115 K; (e)  $P\text{-}x_{\text{N}_2}/y_{\text{N}_2}$  figure for 110 K; (f)  $P\text{-}x_{\text{CH}_4}/y_{\text{CH}_4}$  figure for 110 K. Experimental data (rounds [55], triangles [57], squares [59]), this work (solid lines):  $x_{\text{N}_2}$  and  $y_{\text{N}_2}$  are the mole fractions of  $\text{N}_2$  in the liquid and vapor phases, respectively  $x_{\text{CH}_4}$  and  $y_{\text{CH}_4}$  are the mole fractions of  $\text{CH}_4$  in the liquid and vapor phases, respectively.

## 2.2. The Ternary CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> Mixture

Combining binary interaction parameters of the CH<sub>4</sub>-CO<sub>2</sub>, CO<sub>2</sub>-N<sub>2</sub> and CH<sub>4</sub>-N<sub>2</sub> systems, the EOS can also predict the thermodynamic properties of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixture. Here, the experimental *PVTx* and vapor-liquid equilibrium data are used to verify the accuracy of this EOS for the ternary mixture.

The calculated *PVTx* deviations for the ternary CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixture from each experimental data set [1,64,69–71] are given in Table 4, where the temperature and pressure of the *PVTx* data are up to 573.15 K and 1000 bar, and the composition almost covers the full range. The average absolute deviation of all *PVTx* data from this EOS is 0.40% for the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixture. Figure 4 shows the comparison between experimental data [71] and calculated densities for the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixture at different temperatures. It can be seen from Figure 4 that calculated densities from this EOS are in good agreement with experiment data, indicating that this EOS has a good predictive ability for the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixture.



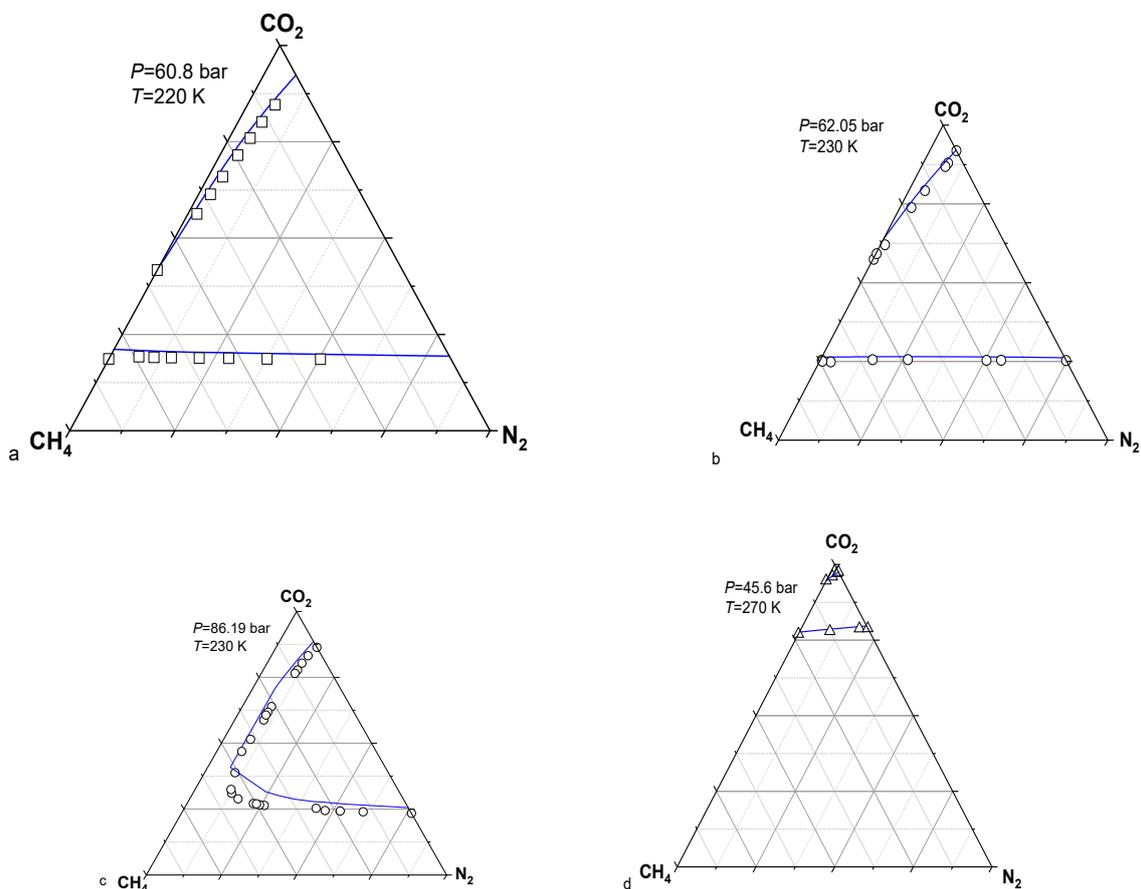
**Figure 4.** Comparisons between experimental and calculated densities for the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixture at different temperatures. (a)  $X_{\text{CO}_2} = 0.0998$ ,  $X_{\text{CH}_4} = 0.6625$ ,  $X_{\text{N}_2} = 0.2477$ ; (b)  $X_{\text{CO}_2} = 0.9949$ ,  $X_{\text{CH}_4} = 0.02$ ,  $X_{\text{N}_2} = 0.0301$ ; (c)  $X_{\text{CO}_2} = 0.1504$ ,  $X_{\text{CH}_4} = 0.1004$ ,  $X_{\text{N}_2} = 0.7492$ ; (d)  $X_{\text{CO}_2} = 0.3346$ ,  $X_{\text{CH}_4} = 0.3319$ ,  $X_{\text{N}_2} = 0.3335$ . Experimental data (squares [71]), this work (solid lines):  $X_{\text{CO}_2}$  is the mole fraction of CO<sub>2</sub> in the mixture,  $X_{\text{CH}_4}$  is the mole fraction of CH<sub>4</sub> in the mixture,  $X_{\text{N}_2}$  is the mole fraction of N<sub>2</sub> in the mixture.

**Table 4.** Calculated  $PVTx$  deviations for the ternary  $\text{CO}_2\text{-CH}_4\text{-N}_2$  mixture.

References	N	$T$ - $P$ - $x$ Range for $x\text{CO}_2$ - $y\text{CH}_4$ -( $1-x-y$ ) $\text{N}_2$				AAD%
		$T$ (K)	$P$ (bar)	$x$	$y$	
McElroy et al. [59]	242	303–333	6–126	0–0.9998	0–0.999	0.45
Seitz et al. [60]	42	474.15	1000	0.0–1.0	0.0–1.0	0.28
Seitz et al. [55]	271	323–573	199–999	0.1–0.8	0.1–0.8	0.28
Zhang et al. [61]	200	293.15–353.25	5–180	0.098–0.9949	0.02–0.6525	0.52
Le et al. [1]	84	305.15	5–600	0.499–0.899	0.0505–0.331	0.82

Note: Deviations of the  $PVTx$  data are the density deviations.  $N$  is the Number of data points, AAD is the abbreviation of average absolute deviation.

The vapor-liquid phase equilibrium properties of the ternary  $\text{CO}_2\text{-CH}_4\text{-N}_2$  system are shown in Figure 5, where the curves are calculated from this EOS, and the experimental data are from the literature of [72–74]. As can be seen from Figure 5, all the experimental data points in the non-critical region agree well with this EOS, but in the near-critical region (Figure 5c), the calculated values deviate largely from experimental data [72–74].



**Figure 5.** Vapor-liquid phase equilibria for the  $\text{CO}_2\text{-CH}_4\text{-N}_2$  system at different temperatures and pressures, (a)  $T = 220$  K,  $P = 60.8$  bar; (b)  $T = 230$  K,  $P = 62.05$  bar; (c)  $T = 230$  K,  $P = 86.19$  bar; (d)  $T = 270$  K,  $P = 45.6$  bar. Experimental data (squares [73], rounds [72], triangles [74]), this work (solid lines).

### 3. The Applications of the ZMS EOS

#### 3.1. Calculating Excess Enthalpies

The excess enthalpy ( $H^E$ ) is an important thermodynamic property of the mixture for the mixing and separation processes, which is defined by

$$H^E = \left\{ H_m - \sum_i x_i H_i \right\}_{P,T} \quad (7)$$

where  $H_m$  is the enthalpy of the mixture,  $H_i$  is the enthalpy of pure component  $i$ , and  $x_i$  is the mole fraction of component  $i$  in the mixture. According to the thermodynamic relationship between Helmholtz free energy and enthalpy,  $H_m$  and  $H_i$  are given by

$$H_m = RT \left( 1 + \tau \left( \alpha_\tau^0 + \alpha_\tau^r \right) + \delta \alpha_\delta^r \right) \quad (8)$$

$$H_i = RT \left( 1 + \tau_i \left( \alpha_{\tau_i}^0 + \alpha_{\tau_i}^r \right) + \delta_i \alpha_{\delta_i}^r \right) \quad (9)$$

where  $\alpha_\tau^0 = \left( \frac{\partial \alpha^0}{\partial \tau} \right)_\delta$ ,  $\alpha_\tau^r = \left( \frac{\partial \alpha^r}{\partial \tau} \right)_\delta$ ,  $\alpha_\delta^r = \left( \frac{\partial \alpha^r}{\partial \delta} \right)_\tau$ ,  $\alpha_{\tau_i}^0 = \left( \frac{\partial \alpha_i^0}{\partial \tau_i} \right)_{\delta_i}$ ,  $\alpha_{\tau_i}^r = \left( \frac{\partial \alpha_i^r}{\partial \tau_i} \right)_{\delta_i}$  and  $\alpha_{\delta_i}^r = \left( \frac{\partial \alpha_i^r}{\partial \delta_i} \right)_{\tau_i}$ .  $\delta_i$  and  $\tau_i$  are the reduced density and inverse reduced temperature of pure component  $i$ , which are defined by

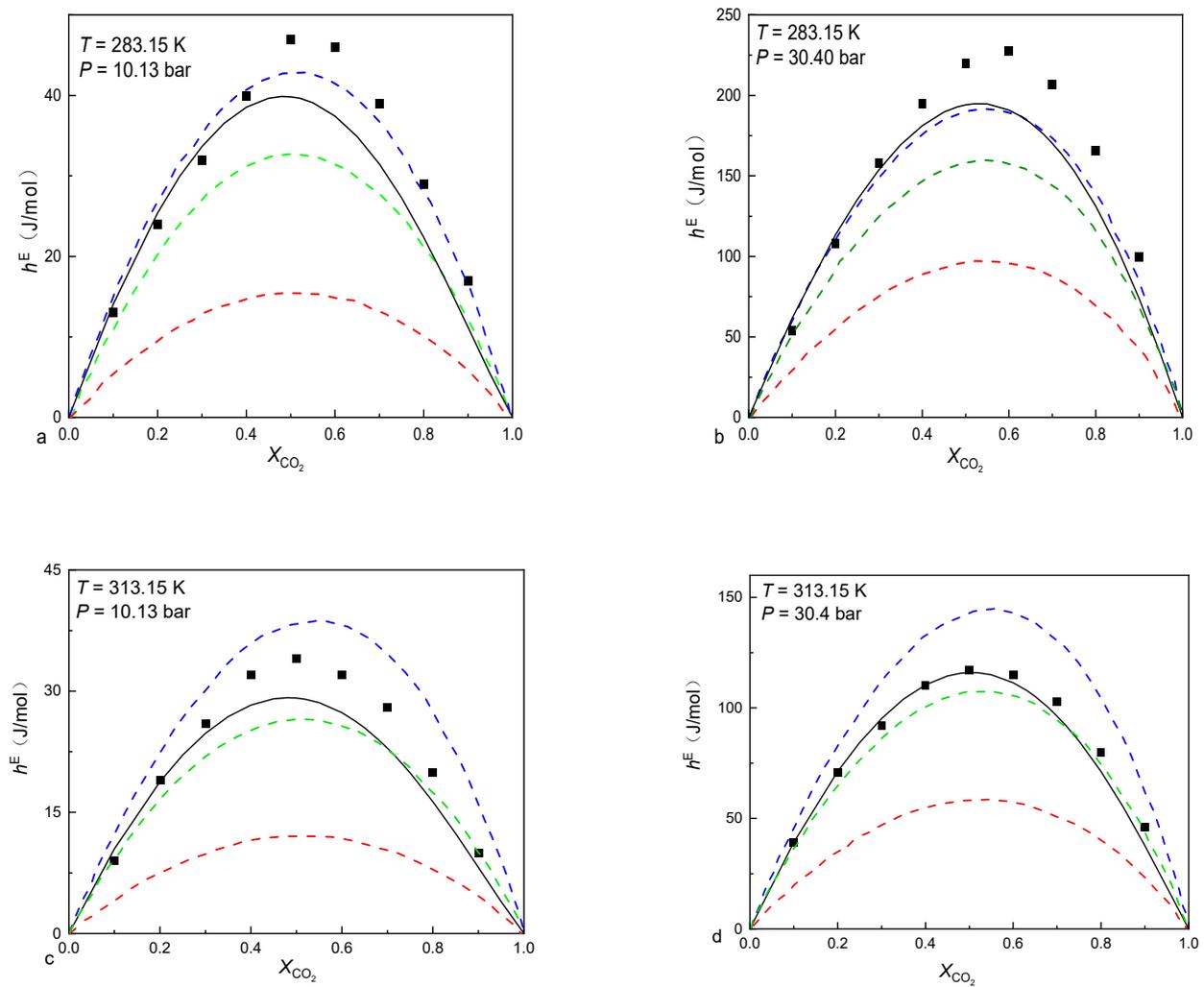
$$\delta_i = \rho_i / \rho_{ci} \quad (10)$$

$$\tau_i = T_{ci} / T \quad (11)$$

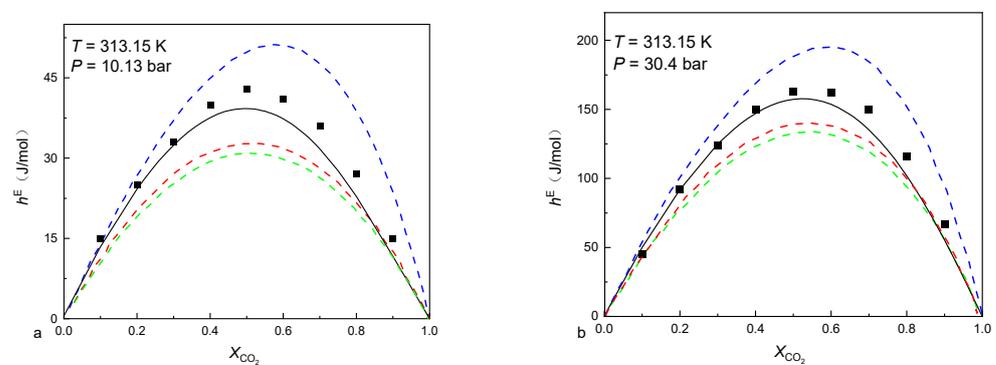
where  $\rho_i$  is the density of component  $i$ . Based on Equations (8)–(11),  $H_m$  can be calculated by this EOS of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures and  $H_i$  can be calculated by the above-mentioned equations of pure CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub> fluids.

The calculated excess enthalpy curves of the CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> mixtures are, respectively, shown in Figures 6 and 7, where the experimental data [75,76] are included for comparison. The following EOSs are also included for comparison: the standard Peng–Robinson (PR) EOS, either optimized (optimal  $k_{ij}$ ) or not ( $k_{ij} = 0$ ); the PR EOS with the residual part of the Wilson excess Helmholtz energy model (PR + EOS/ $a_{res}^{E,Wilson}$ ) [19]. It can be seen from Figures 6 and 7 that this EOS is more accurate than the standard PR EOS, either optimized (optimal  $k_{ij}$ ) or not ( $k_{ij} = 0$ ) at most cases. In general, the PR + EOS/ $a_{res}^{E,Wilson}$  and this EOS shows a similar predictive capability for the excess enthalpies of the CH<sub>4</sub>-CO<sub>2</sub> mixture. However, this EOS is slightly more accurate than the PR + EOS/ $a_{res}^{E,Wilson}$  for the CO<sub>2</sub>-N<sub>2</sub> mixture at low pressures.

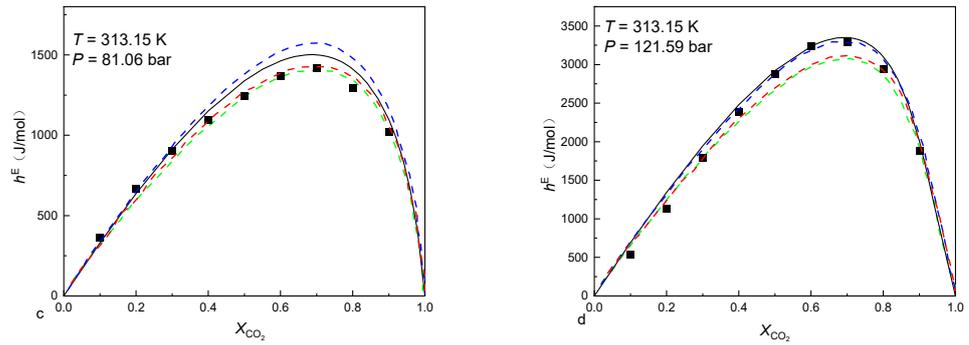
Figure 8 compares the excess enthalpy curves of the CH<sub>4</sub>-N<sub>2</sub> mixture calculated from the ZMS EOS with experimental data [77], and enthalpies calculated from the Enhanced-Predictive-PR78 (E-PPR78) EOS [20] are also added for comparison. Figure 8a shows the ZMSEOS are significantly more accurate than the (E-PPR78) EOS at high and middle pressures. Figure 8b shows the enthalpies of mixing calculated from the E-PPR78 EOS are slightly better than those of the ZMS EOS at low and middle pressures, but the enthalpies of mixing calculated from the ZMS EOS are better than those from the E-PPR78 EOS at  $P = 101.33$  bar.



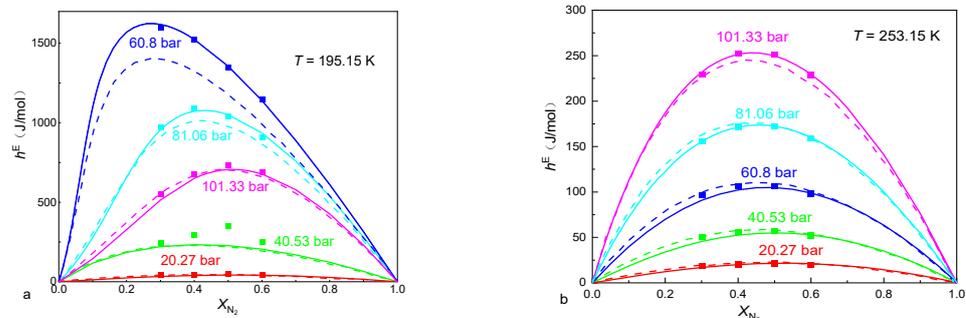
**Figure 6.** Comparisons between experimental and calculated mixing enthalpies for the CH<sub>4</sub>-CO<sub>2</sub> mixture, (a)  $T = 283.15\text{ K}$ ,  $P = 10.13\text{ bar}$ ; (b)  $T = 283.15\text{ K}$ ,  $P = 30.40\text{ bar}$ ; (c)  $T = 313.15\text{ K}$ ,  $P = 10.13\text{ bar}$ ; (d)  $T = 313.15\text{ K}$ ,  $P = 30.40\text{ bar}$ . Experimental data (squares [75]). PR ( $k_{ij} = 0$ ) (red dash curves); PR (optimal  $k_{ij}$ ) (green dash curves); PR + EOS/ $a_{res}^{E,Wilson}$  (blue dash curves); This work (black solid curves);  $h^E$  is the mixing enthalpy.



**Figure 7.** Cont.



**Figure 7.** Comparisons between experimental and calculated mixing enthalpies for the CO<sub>2</sub>-N<sub>2</sub> mixture, (a)  $T = 313.15\text{ K}$ ,  $P = 10.13\text{ bar}$ ; (b)  $T = 313.15\text{ K}$ ,  $P = 30.40\text{ bar}$ ; (c)  $T = 313.15\text{ K}$ ,  $P = 81.06\text{ bar}$ ; (d)  $T = 313.15\text{ K}$ ,  $P = 121.59\text{ bar}$ . Experimental data (black squares [76]). PR ( $k_{ij} = 0$ ) (red dash curves); PR (optimal  $k_{ij}$ ) (green dash curves); PR + EOS/ $a_{Tres}^{E,Wilson}$  (blue dash curves); This work (black solid curves).



**Figure 8.** Comparisons between experimental and calculated mixing enthalpies for the CH<sub>4</sub>-N<sub>2</sub> mixture at 195.15 K and 253.15 K. (a)  $T = 195.15\text{ K}$ ; (b)  $T = 253.15\text{ K}$ . Experimental data (squares [77]), this work (solid curves); the Enhanced-Predictive-PR78 (E-PPR78) EOS (dash curves).

### 3.2. Calculating the Solubility of CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> Mixtures in Aqueous Electrolyte Solution

The solubility of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> gas mixtures in the electrolyte solutions can provide the quantitative assessment for the storage of CO<sub>2</sub> in deep saline aquifers. According to the thermodynamic principle, when the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> gas mixtures reach dissolution equilibrium in the electrolyte solution, the chemical potential of component  $i$  in the liquid phase ( $\mu_i^L$ ) and its chemical potential in the vapor phase ( $\mu_i^V$ ) are equal. The chemical potential can be written in terms of fugacity in the vapor phase and activity in the liquid phase

$$\mu_i^L = \mu_i^{L(0)} + RT \ln m_i \gamma_i \tag{12}$$

$$\mu_i^V = \mu_i^{V(0)} + RT \ln \varphi_i y_i P \tag{13}$$

where  $i$  is the gas component in vapor mixtures, and  $P$  is the pressure in bar.  $y_i$  is the molar fraction of  $i$  in the vapor phase,  $\varphi_i$  is the fugacity coefficient of  $i$  in the vapor phase and  $m_i$  is the solubility of  $i$  in the liquid phase in mol/kg.  $\gamma_i$  is the activity coefficient of component  $i$  in the liquid phase.  $\mu_i^{L(0)}$  and  $\mu_i^{V(0)}$  are the standard chemical potential of  $i$  in the liquid and vapor phase, respectively.

At the dissolution phase equilibrium,  $\mu_i^V = \mu_i^L$ . From Equations (12) and (13), we can obtain

$$\ln m_i = \ln y_i P + \ln \varphi_i - \frac{\mu_i^{L(0)} - \mu_i^{V(0)}}{RT} - \ln \gamma_i \tag{14}$$

Since the water content in the vapor phase at equilibrium is generally small, it has a negligible effect on the fugacity coefficient. When calculating the fugacity coefficient, the water content in the vapor phase can be ignored, the vapor phase can be approximated as the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> system and the mole fraction of gas component *i* in the vapor phase is expressed as *y<sub>i</sub>*. Consequently, *φ<sub>i</sub>* can be calculated from the ZMS EOS developed in this work.

Pitzer activity coefficient model [78] is chosen to calculate *γ<sub>i</sub>*

$$\ln \gamma_i = \sum_c 2\lambda_{i-c} m_c + \sum_a 2\lambda_{i-a} m_a + \sum_c \sum_a \zeta_{i-c-a} m_c m_a \tag{15}$$

where *λ* and *ζ* are second-order and third-order interaction parameters, respectively; *c* and *a* refer to cation and anion, respectively. Substituting Equation (15) into Equation (14) yields

$$\begin{aligned} \ln m_i = & \ln y_i P + \ln \varphi_i - \frac{\mu_i^{L(0)} - \mu_i^{V(0)}}{RT} - \sum_c 2\lambda_{i-c} m_c \\ & - \sum_a 2\lambda_{i-a} m_a - \sum_c \sum_a \zeta_{i-c-a} m_c m_a \end{aligned} \tag{16}$$

As can be seen from Equation (17), *m<sub>i</sub>* is related to the difference between *μ<sub>i</sub><sup>L(0)</sup>* and *μ<sub>i</sub><sup>V(0)</sup>*, and is not related to the specific value of *μ<sub>i</sub><sup>L(0)</sup>* or *μ<sub>i</sub><sup>V(0)</sup>*. Therefore, to simplify the model, *μ<sub>i</sub><sup>V(0)</sup>* is assumed to be zero. Equation (16) can be simplified as

$$\begin{aligned} \ln m_i = & \ln y_i P + \ln \varphi_i - \frac{\mu_i^{L(0)}}{RT} - \sum_c 2\lambda_{i-c} m_c \\ & - \sum_a 2\lambda_{i-a} m_a - \sum_c \sum_a \zeta_{i-c-a} m_c m_a \end{aligned} \tag{17}$$

where  $\frac{\mu_i^{L(0)}}{RT}$ , *λ* and *ζ* are the function of temperature and pressure, and the parameters can be determined by the solubility model of pure gases (CO<sub>2</sub>, CH<sub>4</sub>, N<sub>2</sub>). Since the 1980s, many solubility models for CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub> gases in pure water and brine have been developed [79–91], most of which do not exceed 473 K and 1000 bar. Among them, the solubility models of CH<sub>4</sub>, CO<sub>2</sub> and N<sub>2</sub> established by Mao and co-workers [83,86,91] are applicable to a wider range of temperature, pressure, and salinity with higher accuracy. They have been chosen as the solubility models of pure CH<sub>4</sub>, CO<sub>2</sub> and N<sub>2</sub> gases in this work. Table 5 lists the applicable temperature, pressure, and salinity ranges of these solubility models for pure gases.

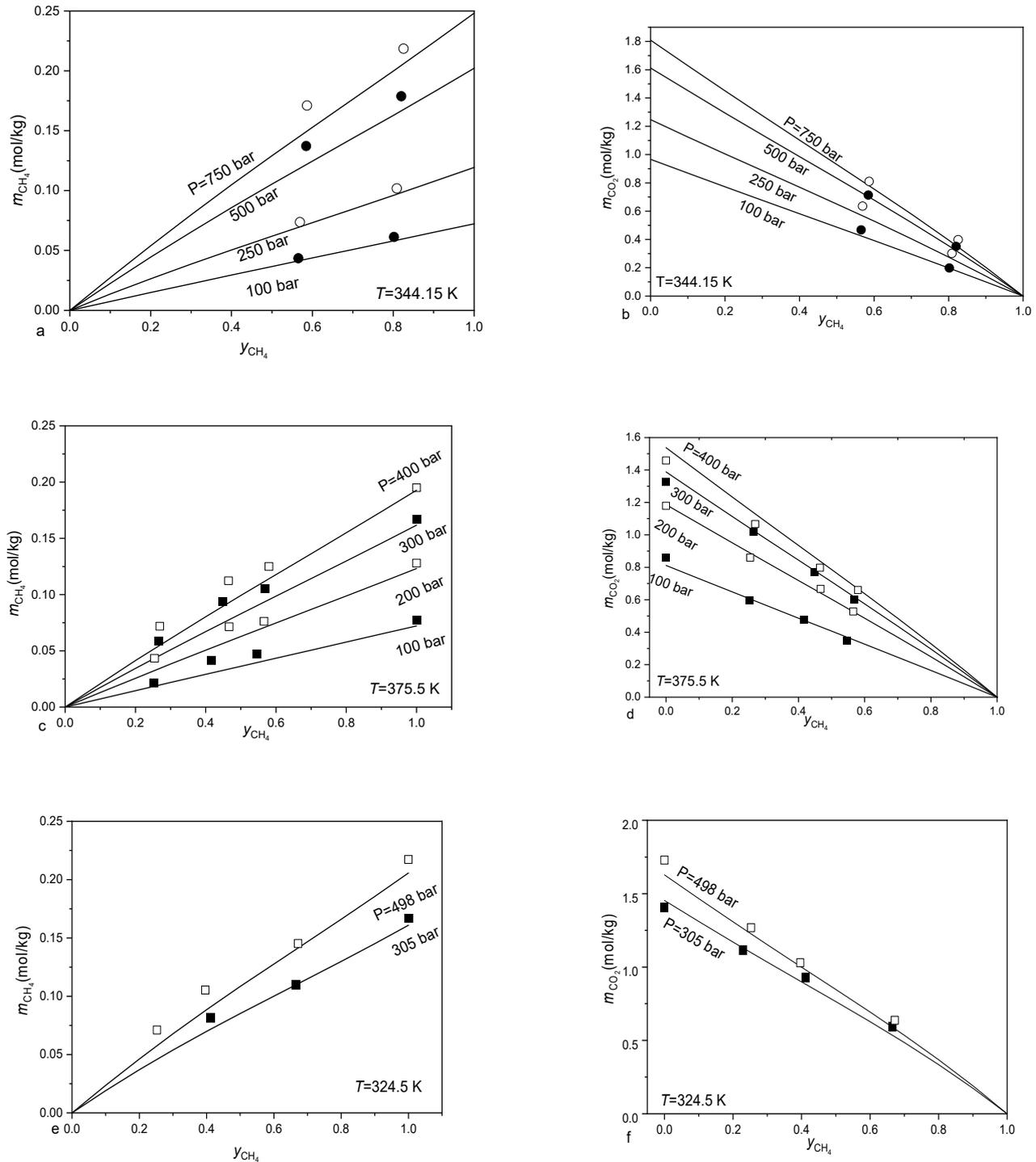
**Table 5.** Ranges of pure gas solubility models.

Pure Gas	<i>T</i>	<i>P</i>	Salinity ( <i>m<sub>NaCl</sub></i> )	References
CH <sub>4</sub>	273–523 K	1–2000 bar	0–6 mol/kg	Duan and Mao [83]
CO <sub>2</sub>	273.15–723.15 K	1–1500 bar	0–4.5 mol/kg	Mao et al. [91]
N <sub>2</sub>	273–590 K	1–600 bar	0–6 mol/kg	Mao and Duan [86]

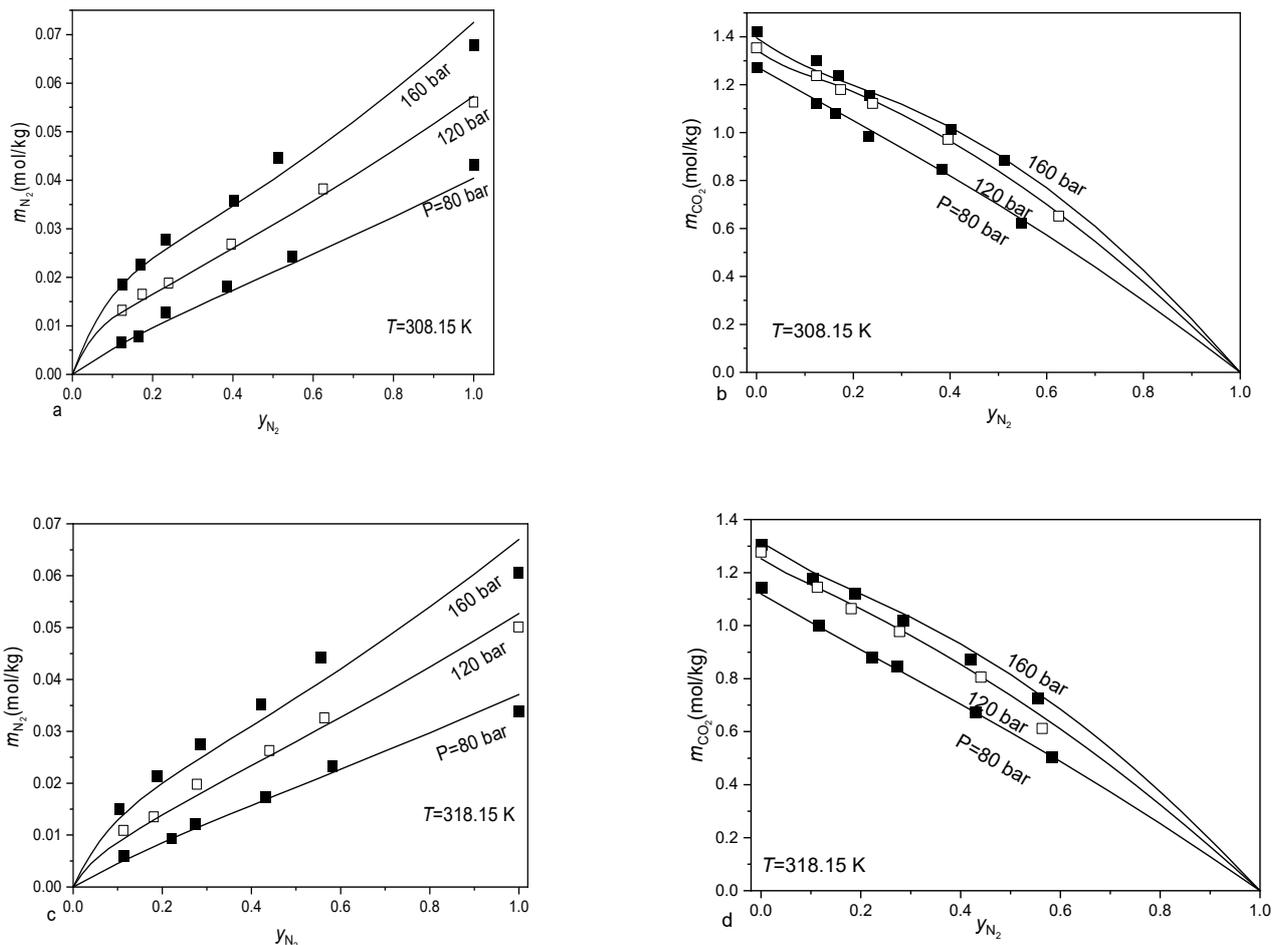
To verify the accuracy of the predictions, the experimental data for the solubility of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> gas mixtures are compared with the calculated results. Figure 9 compares the experimental data on the solubility of the CH<sub>4</sub>-CO<sub>2</sub> mixture in pure water [92,93] for the temperature range of 324.5–375.5 K and the pressure range of 100–750 bar, and it can be seen that the agreement is very good.

The solubility experimental data of the CO<sub>2</sub>-N<sub>2</sub> gas mixture in pure water [94] and the calculated results are compared in Figure 10, where the temperature range is 308.15–318.15 K and the pressure range is 80–160 bar. As can be seen from Figure 10 the predictive results are in good agreement with the experimental data. The average absolute deviations between the calculated N<sub>2</sub> and CO<sub>2</sub> solubility of this model and the experimental data are 6.04% and 1.52%, respectively. Figure 11 compares the solubility

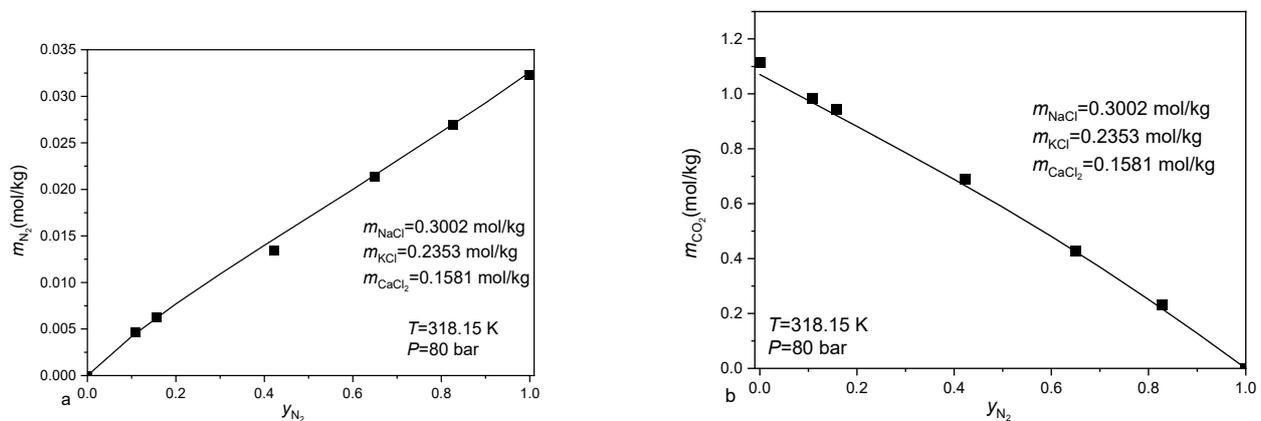
experimental data of the CO<sub>2</sub>-N<sub>2</sub> gas mixture in saline water. The average absolute deviations between the calculated N<sub>2</sub> and CO<sub>2</sub> solubility and the experimental data are 2.04% and 2.49%, respectively.



**Figure 9.** Comparisons between experimental and calculated solubilities for the CH<sub>4</sub>-CO<sub>2</sub> mixture in water. (a)  $m_{CH_4}$ - $y_{CH_4}$  figure at 344.15 K; (b)  $m_{CO_2}$  -  $y_{CH_4}$  figure at 344.15 K; (c)  $m_{CH_4}$  -  $y_{CH_4}$  figure at 373.5 K; (d)  $m_{CO_2}$  -  $y_{CH_4}$  figure at 373.5 K; (e)  $m_{CH_4}$  -  $y_{CH_4}$  figure at 324.5 K; (f)  $m_{CO_2}$  -  $y_{CH_4}$  figure at 324.5 K. Experimental data (rounds [92], squares [93]), this work (solid curves):  $m_{CO_2}$  is the solubility of CO<sub>2</sub>,  $m_{CH_4}$  is the solubility of CH<sub>4</sub>.



**Figure 10.** Comparisons between experimental and calculated solubilities for the CO<sub>2</sub>-N<sub>2</sub> mixture in water. (a)  $m_{N_2} - y_{N_2}$  figure at 308.15 K; (b)  $m_{CO_2} - y_{N_2}$  figure at 308.15 K; (c)  $m_{N_2} - y_{N_2}$  figure at 318.15 K; (d)  $m_{CO_2} - y_{N_2}$  figure at 318.15 K. Experimental data (squares [94]), this work (solid curves);  $m_{N_2}$  is the solubility of N<sub>2</sub>.



**Figure 11.** Comparisons between experimental and calculated solubilities for the CO<sub>2</sub>-N<sub>2</sub> mixture in saline water. (a)  $m_{N_2} - y_{N_2}$  figure; (b)  $m_{CO_2} - y_{N_2}$  figure. Experimental data (black squares [94]), this work (solid curves).

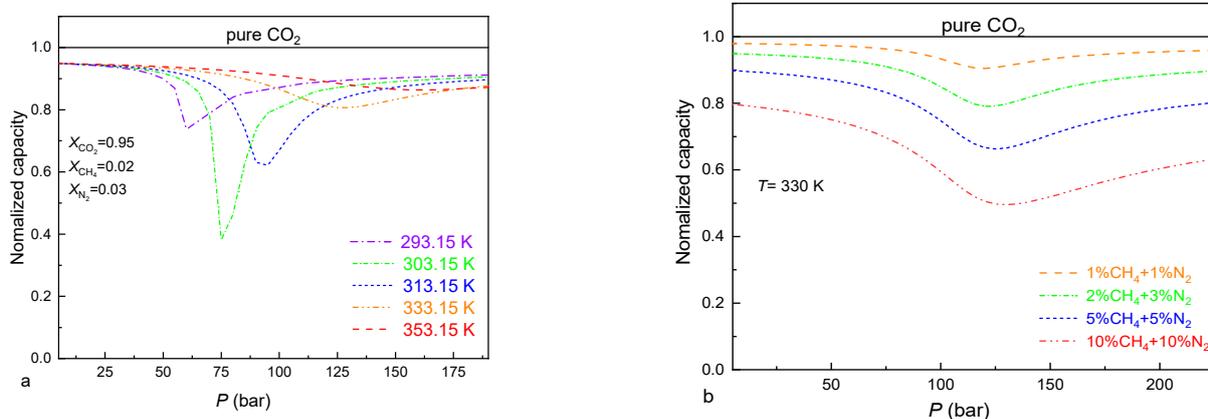
### 3.3. The Impact of Impurities ( $\text{CH}_4$ and $\text{N}_2$ ) on the $\text{CO}_2$ Storage Capacity

$\text{CH}_4$  and  $\text{N}_2$  are non-condensable impurities, which reduce the  $\text{CO}_2$  storage capacity in geological formations. Based on the EOS of the  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  fluid mixtures, the impact of  $\text{CH}_4$  and  $\text{N}_2$  on  $\text{CO}_2$  storage capacity can be calculated quantitatively by the normalized storage capacity proposed by Wang et al. [9]. For a certain reservoir with a certain volume, the normalized storage capacity of the  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  gas mixtures is shown as

$$\frac{A}{A_{\text{CO}_2}} = \frac{\rho}{\rho_{\text{CO}_2} \left(1 + \frac{M_{\text{CH}_4}}{M_{\text{CO}_2}} + \frac{M_{\text{N}_2}}{M_{\text{CO}_2}}\right)} \quad (18)$$

where  $A$  and  $A_{\text{CO}_2}$  are the mass of  $\text{CO}_2$  in the mixtures and the mass of pure  $\text{CO}_2$  under the same volume, respectively.  $\rho_{\text{CO}_2}$  and  $\rho$  are the density of the pure  $\text{CO}_2$  and  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  gas mixtures, respectively;  $M_{\text{CH}_4}$ ,  $M_{\text{N}_2}$  and  $M_{\text{CO}_2}$  are the mass of  $\text{CH}_4$ ,  $\text{N}_2$ , and  $\text{CO}_2$  in the mixture, respectively.  $A/A_{\text{CO}_2}$  is the ratio of the mass of  $\text{CO}_2$  per unit volume in the mixture to that in the pure state. The value of  $A/A_{\text{CO}_2}$  can be viewed as the normalized storage capacity for  $\text{CO}_2$  (i.e., the storage capacity for structural trapping of  $\text{CO}_2$ ). For pure  $\text{CO}_2$ , the normalized storage capacity ( $A/A_{\text{CO}_2}$ ) is 1. For the  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  gas mixtures of a given composition,  $M_{\text{CH}_4}$ ,  $M_{\text{N}_2}$  and  $M_{\text{CO}_2}$  are also known.  $\rho_{\text{CO}_2}$  can be calculated by the above-mentioned equation of pure  $\text{CO}_2$  fluid and  $\rho$  can be calculated by the ZMS EOS.

The normalized storage capacity of the  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  fluid mixture calculated by the ZMS EOS is plotted in Figure 12. Figure 12a shows the normalized storage capacity of a given composition at different temperatures. With the increase in pressure, the normalized storage capacity decreases at first and then increases, and there is a minimum point. When the temperature changes, the pressure corresponding to the lowest point of the normalized storage also changes. Figure 12b shows the normalized storage capacity of different compositions at the same temperature, which indicates that the content of impurities is much larger and the normalized storage capacity is much smaller.



**Figure 12.** Normalized  $\text{CO}_2$  storage capacity calculated by this EOS of  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  mixture. (a) the Normalized capacity- $P$  figure at different temperatures for the same composition; (b) the Normalized capacity- $P$  figure at different compositions for the same temperatures.

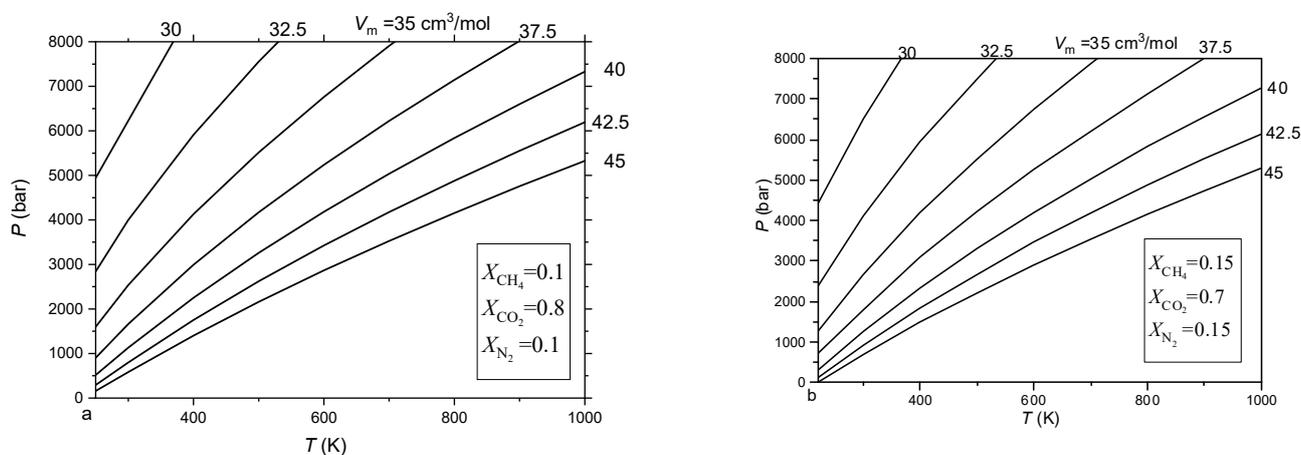
### 3.4. Isochores of the $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$ fluid Inclusions

In the studies of fluid inclusions, the isochores (pressure-temperature relation at constant density and composition) are frequently used to estimate the trapping temperatures and pressures.

Based on the EOS of the  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  fluid mixtures, isochores of the  $\text{CO}_2$ - $\text{CH}_4$ - $\text{N}_2$  inclusions can be calculated by the following equation:

$$P(\delta, \tau, x) = \rho RT [1 + \delta \alpha'_\delta] \quad (19)$$

The calculated isochores of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> inclusions at two different compositions are plotted in Figure 13, from which it can be seen that the isochores of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> inclusions are a bit curved.



**Figure 13.** Calculated isochores of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> mixture by this EOS, (a)  $X_{\text{CO}_2} = 0.8$ ,  $X_{\text{CH}_4} = 0.1$ ,  $X_{\text{N}_2} = 0.1$ ; (b)  $X_{\text{CO}_2} = 0.7$ ,  $X_{\text{CH}_4} = 0.15$ ,  $X_{\text{N}_2} = 0.15$ :  $V_m$  is the molar volume.

This section may be divided into subheadings. It should provide a concise and precise description of the experimental results, their interpretation, as well as the experimental conclusions that can be drawn.

#### 4. Conclusions

A fundamental EOS for the Helmholtz free energy of the CH<sub>4</sub>-N<sub>2</sub> mixture has been developed by using four binary interaction parameters. Comparisons with experimental  $PVTx$  and VLE data available showed that the EOS can satisfactorily reproduce the experimental volumetric and vapor-liquid phase equilibria data of binary CH<sub>4</sub>-N<sub>2</sub> mixtures up to 673.15 K and 1380 bar, with or close to experimental accuracy. Combining this EOS of the CH<sub>4</sub>-N<sub>2</sub> fluid mixtures and the EOS of the CH<sub>4</sub>-CO<sub>2</sub> and CO<sub>2</sub>-N<sub>2</sub> fluid mixtures developed in our previous work, an EOS of ternary CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures has been presented, which is named ZMS EOS. The ZMS EOS can be applied to calculate excess enthalpies, predict the solubility of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> gas mixtures in water and brine, estimate the impurities CH<sub>4</sub> and N<sub>2</sub> on the CO<sub>2</sub> storage capacity, and calculate the isochores of the CO<sub>2</sub>-CH<sub>4</sub>-N<sub>2</sub> fluid mixtures.

**Author Contributions:** J.Z.: Conceptualization, Software, Investigation, Writing—original. S.M.: Supervision, Methodology, Writing—review, and editing. Z.S.: Investigation, Software, Validation. All authors have read and agreed to the published version of the manuscript.

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