


Article

Digital Shearing Speckle Pattern Interferometry Based on Rochon Prism and Its Application

Xu Wang , Zhan Gao *, Chenjia Gao, Jieming Zhao, Yuhao Niu, Zihui Liu, Yuan Zhang and Yuchen Liu

Key Laboratory of Luminescence and Optical Information of Ministry of Education, Beijing Jiaotong University, Beijing 100044, China; xuwang@bjtu.edu.cn (X.W.); 16118443@bjtu.edu.cn (C.G.); 18121679@bjtu.edu.cn (J.Z.); 18121665@bjtu.edu.cn (Y.N.); 17121652@bjtu.edu.cn (Z.L.); 17121665@bjtu.edu.cn (Y.Z.); 17121651@bjtu.edu.cn (Y.L.)

* Correspondence: zhangao@bjtu.edu.cn

Received: 20 May 2019; Accepted: 20 June 2019; Published: 22 June 2019



Abstract: This paper presents improved digital shearing speckle pattern interferometry based on Rochon Prism. The advantage of this method is that it avoids the shadow noise caused by transmission light, which improves the shearography accuracy. By reducing the angle of divergence and coordinating the measuring distance and the precision coordinate, the beam-splitting path of the interference optical path is improved. Experimental results prove that the improved method contributes to the practical and instrumental application of shearography.

Keywords: speckle; shearography; Rochon Prism; Wollaston Prism

1. Introduction

It has been nearly 40 years since Hung and Leendertz introduced shearography [1,2]. As an important optical measurement method, shearography has been widely used in the fields of shape measurement [3], vibration analysis [4,5], and strain detection [6–8]. Generally, a basic shearographic measurement system uses a shearing device, such as a Michelson interferometer, a Wollaston Prism (WP), or an optical wedge [9]. By passing through a shearing device, two identical but slightly misaligned images are generated and interfere with each other.

In a previous study, we introduced a shearographic measurement system based on the WP and the temporal wavelet transformation (TWT) [10]. Compared with the classical shearing devices, such as the Michelson interferometer, WP has an advantage in terms of common optical path arrangement, which improves the measurement stability and repeatability. WP also does not require the shifting structure of the Michelson interferometer. In further experiments, we proved the feasibility of combining the polarization splitting effect with the electro-optical crystal materials in temporal heterodyne shearing speckle pattern interferometry [11]. However, there are still some problems to be solved. For example, the polarization splitting effect of WP is not specially perfect under an illuminant of high intensity, and the lights emergent from the prism are not only an e-o light and an o-e light, but also a transmission light, which confuse the interference signals and reduce the signal-to-noise ratio (SNR) and the image contrast [12].

To solve these problems, we introduced the Rochon Prism (RP) to replace the WP. RP is also a traditional optical beam splitter, but different from the WP, the o-light is transmitted without refraction angle passing through the RP. As such, only two kinds of interference light occur in the interference field generated by the RP. In this study, we analyzed the advantages of the shearography based on RP, and conducted relevant experiments to compare the errors between WP and RP. The experimental

results show that the RP meets the shearographic requirements and improves the accuracy of the measurement system.

2. Principle of DSSPI

After subtraction, the fringe patterns are visible and can be described as [11]:

$$I_d = I(t) - I(0) = I_0 \left[\gamma \sin \left(\varphi_0 + \frac{\Delta\varphi}{2} \right) \sin \left(\frac{\Delta\varphi}{2} \right) \right], \tag{1}$$

where I_0 is the light intensity of background, γ is the amplitude of modulation, φ_0 is the random initial phase, and $\Delta\varphi$ is the phase variable quantity that depends on the small deformation of the measured object by the loading. Equation (1) means that a high-frequency carrier $\sin(\varphi_0 + \frac{\Delta\varphi}{2})$ is modulated by a low-frequency factor $\sin(\frac{\Delta\varphi}{2})$.

The deformation derivative can be expressed by [11]:

$$\frac{\partial w}{\partial x} = \frac{\lambda \Delta\varphi}{4\pi \Delta x'}, \tag{2}$$

where $\frac{\partial w}{\partial x}$ is the out-of-plane displacement derivative, λ is the wavelength of the illuminant, and Δx is the shearing distance. This means the out-of-plane displacement derivative can be obtained if we acquire the distribution of the phase difference.

In our previous research, we introduced the TWT to determine the wrapped phase [10]. The continuous wavelet coefficient can be expressed as:

$$W_f(a, b) = |a|^{-1/2} \int I(x, y, t) \psi^* \left(\frac{t-b}{a} \right) dt, \tag{3}$$

where a is the scale parameter, b is the shift parameter, and ψ^* is the mother wavelet. The phase and the amplitude are given by:

$$A(a, b) = \sqrt{\{Im[W_f(a, b)]\}^2 + \{Re[W_f(a, b)]\}^2}, \tag{4}$$

$$\varphi(a, b) = \arctan \frac{Im[W_f(a, b)]}{Re[W_f(a, b)]}, \tag{5}$$

where $Im(W_f(a, b))$ is the imaginary part of $W_f(a, b)$ and $Re(W_f(a, b))$ is the real part. After unwrapping the phase, $\frac{\partial w}{\partial x}$ can be solved using Equation (2).

3. Comparing RP with WP

The RP and WP schemes are shown in Figure 1. The difference between RP and WP is that their optical axes of the crystal in the first prism (P1) are mutually perpendicular. In Figure 1a, the unpolarized incident beam is divided into o-light and e-light, where the o-light travels along the original optical axis of the measurement system (x-axis) without deviation and the e-light deviates from the optical axis with an angle of α . However, in Figure 1b, the incident beam is divided into o-e light and e-o light, neither of which remain along the x-axis.

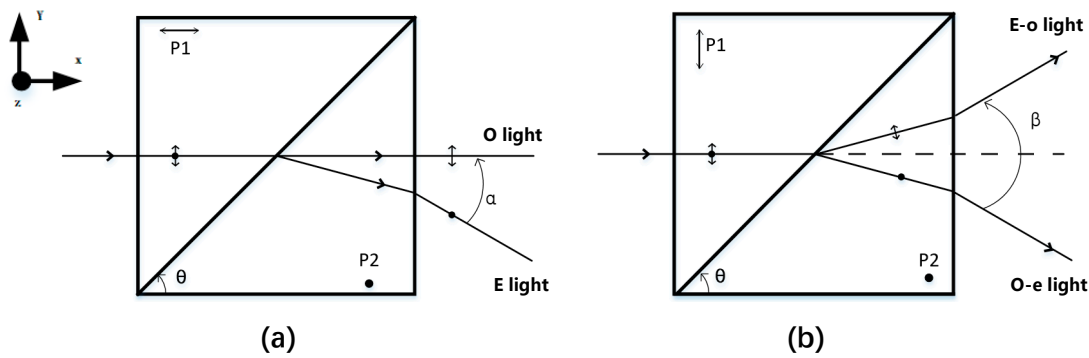


Figure 1. Schemes of (a) a Rochon Prism and (b) a Wollaston Prism, where the line of dashes is the optical axis of the system.

This difference produces RP’s two advantages. First, maintaining the o light along the optical axis improves the SNR. As shown in Figure 2b, the sheared images generated by WP produce three overlapping images: the e-o light, the o-e light, and the transmission light, which is along the x-axis in Figure 1b. Due to unexpected transmission light, there are three overlapping fringe patterns rather than one when the WP is applied in the shearography system. Contrary to WP, RP solves this problem automatically because of its optical path structure shown in Figure 1a. Figure 2a shows that there are only two sheared images (the o-light and the e-light) and no transmission light, which protects the fringe pattern from background noise.

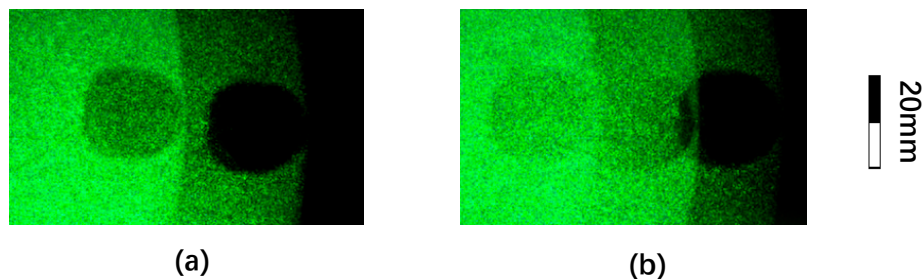


Figure 2. Images after shearing devices under 150-mW power laser illumination: (a) sheared by Rochon Prism (RP) and (b) sheared by Wollaston Prism (WP).

Second, RP provides a more suitable angle of divergence. As shown in Figure 1a, the optical path difference (OPD) is composed of the following two parts: the OPD in P2 where o-light and e-light are separated, and the optical path difference in the air caused by the angle of emergence between the o-light and e-light.

Next, we deduced the OPD in an RP. In the P1 of an RP, both the beams have the same refractive index and are not affected by the birefringence effect, which is different from the WP. When the beam enters the second prism (P2), the o-light has the same refractive index and is undeviated. The e-light has a different refractive index and bends away from the x-axis in Figure 1a, causing the o-light and e-light to separate after their emergence. If the angle between the o-light and e-light is θ' , the OPD in P2 is:

$$\Delta[l]_1 = y' \left(\frac{n_e}{\cos \theta'} - n_o \right), \tag{6}$$

where y' is the distance that the o-light travels in P2, which can be described as:

$$\begin{aligned} y' &= L - \left(\frac{L}{2 \tan \theta} - x_1 \right) \tan \theta \\ &= \frac{L}{2} + x_1 \tan \theta \end{aligned} \tag{7}$$

where x_i is the distance to the geometric axis of the prism (positive above the central axis and negative below) and θ is the prisms' wedge angle. The relationship between θ' and θ is:

$$\frac{\cos \theta}{\sin(\theta - \theta')} = \frac{n_o}{n_e'} \quad (8)$$

After emergence, the angle of divergence between o- and e-light is α . As o- and e-light have the same refractive index, the OPD in the air can be described as:

$$\Delta[l]_2 = y \left(\frac{1}{\cos \alpha} - 1 \right), \quad (9)$$

where y is the distance between the prism and the CCD.

As most of the parameters are determined in a specific system, the general form of the phase difference can be obtained after combining the OPD of two parts, that is:

$$\varphi_{\Delta} = \frac{\Delta[l]_1 + \Delta[l]_2}{\lambda} = c + \frac{\tan \theta}{\lambda} \cdot x_1, \quad (10)$$

where c is constant and λ is the wavelength. In a particular measuring system, the phase difference at a certain point only depends on the distance from the axis of the prism.

If the shearing distance is Δx , which means that each point interferes with a corresponding point that is Δx away, the inherent phase difference of the interference system is:

$$\varphi_{\Delta} = \frac{\tan \theta}{\lambda} \cdot \Delta x. \quad (11)$$

An immobile phase difference means that RP would not introduce phase difference noise into the system, which is the same as WP. However, although the inherent phase difference of RP does not change due to the deformation of the object, the shearing distance of the prism is still reflected in the fringe distribution through Equation (10). The larger the shearing distance, the narrower the measured zone. The fringe pattern is sparser and the SNR is reduced with a large shearing distance. A smaller shearing distance, as shown in Equation (2), leads to a tighter pattern and lower contrast, which reduces the intensity and accuracy of the result [9].

To verify the conclusion, we conducted a simulation, as shown in Figure 3. Different shearing distances led to different distributions in the phase variety. In Figure 3a, the surface after loaded is supposed to be Gaussian function. From Figure 3b–d, the shearing distance was set from 3 to 12, which progressively lowered the peak of the difference distribution.

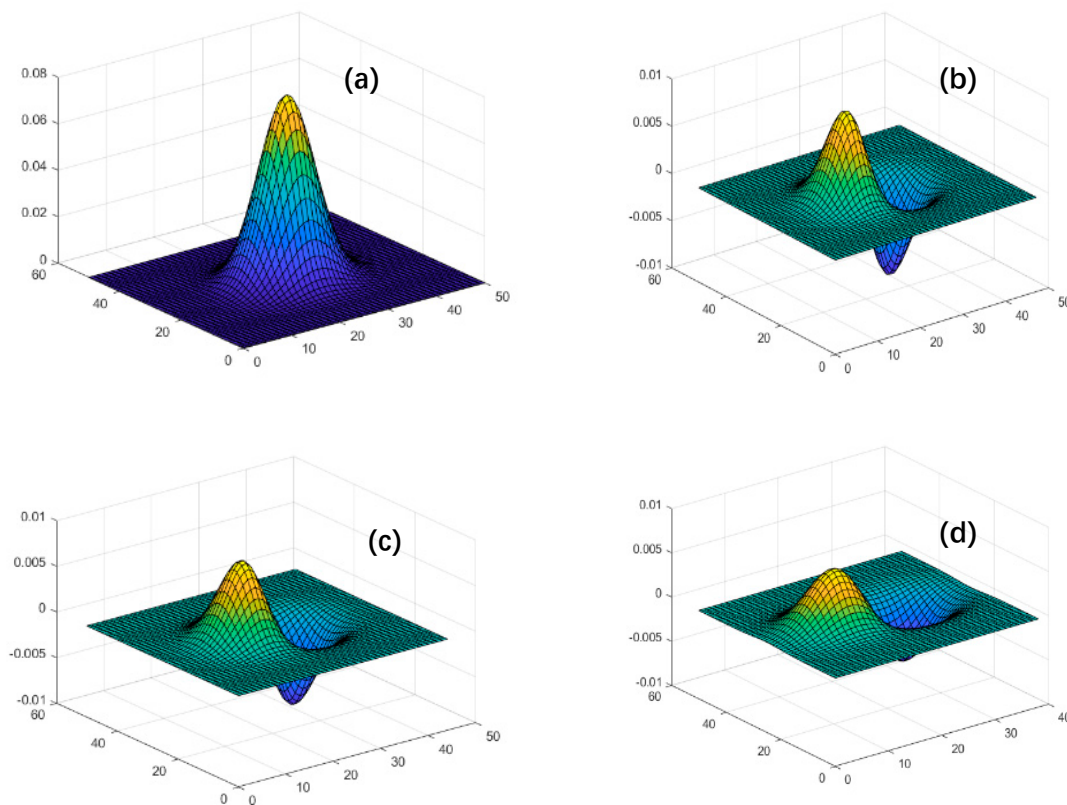


Figure 3. The simulation of difference distribution in various shearing distance: (a) Gaussian function, (b) $\Delta x = 3$, (c) $\Delta x = 6$, and (d) $\Delta x = 12$. Without loss of generality in conclusion, the units are omitted.

In reference [13], the relationship mentioned above was analyzed and the fringe patterns under different shear were given, as shown in Figure 4. The shearing distance is set to 10 mm in Figure 4a, 20 mm in Figure 4b, and 30 mm in Figure 4c. When the shearing distance is too small, the stripes become dimmer and less visible because the phase difference between the interfered light sources is small and the overall phase changes is minimal, so the intensity at each pixel is not so different. As shown in Figure 2, an excessive shearing distance reduces the maximum value of the displacement gradient, which also reduces the contrast of the interference pattern. Therefore, we should consider both when adjusting the shearing distance.

In some cases, the sensitivity is desired to be adjustable [14,15]. In Michelson interferometers with one of the mirrors tilted [9], the shearing distance can be changed while the angle of the mirror is changed. As far as splitters such as WP and RP are concerned, their sensitivity, which depends on their shearing distance, is determined by their angles of divergence and their distance to the measured object. Under limited measuring conditions, such as a certain measuring distance and a fixed system resolution, we must choose a beam splitter with an appropriate emergent angle. Generally, the WP angle of divergence is larger than that of RP, which means that to maintain the proper shearing distance generated by WP, we have to reduce the measured distance and the measured area.

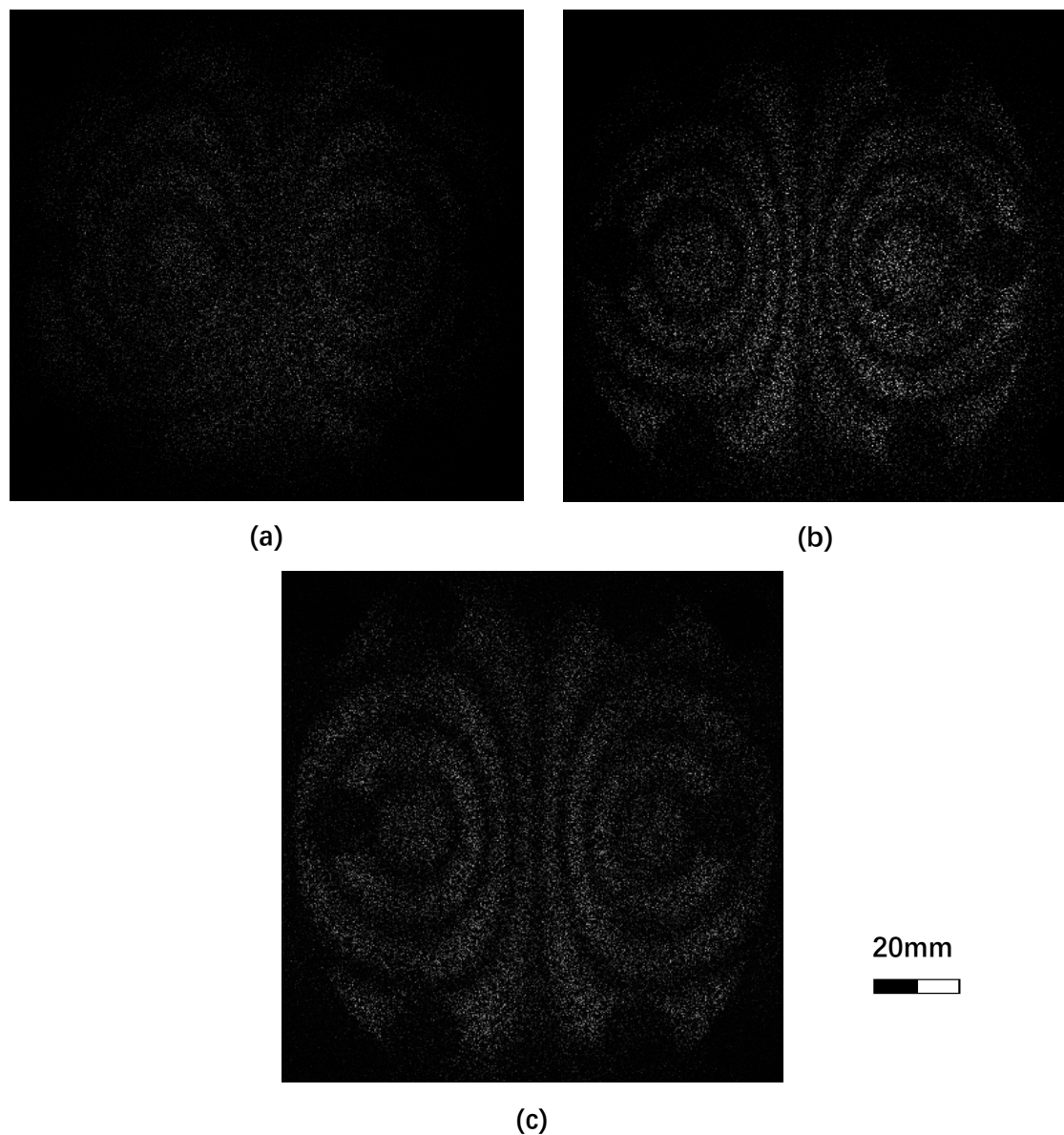


Figure 4. Fringe pattern under different shearing distance: (a) $\Delta x = 10$ mm, (b) $\Delta x = 20$ mm, and (c) $\Delta x = 30$ mm.

4. Experiments and Discussion

Here, we built a system to compare the performance of the two prisms under high power illumination. As shown in Figure 5, the system is based on digital shearing speckle pattern interferometry (DSSPI). The illuminant was a 532 nm green single longitudinal mode laser (MSL-FN-532-150mW). The shearing devices were based on a WP and a RP respectively, the angles of divergence were 1.52° and 1.50° , respectively

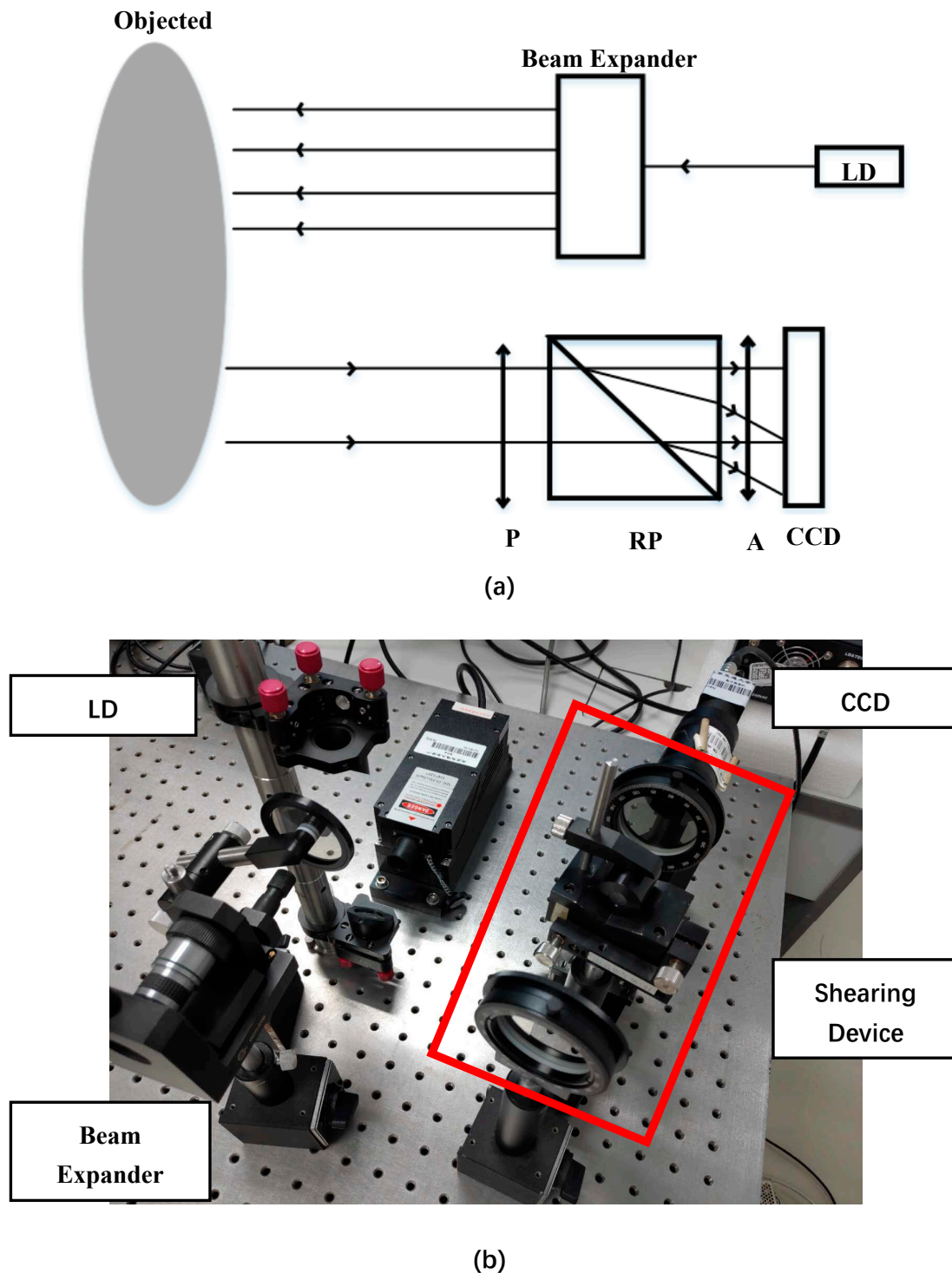


Figure 5. (a) The measurement scheme and (b) the setup of the measurement system.

In the experiment, a positive temperature coefficient (PTC) resistance was applied as the heat source to load the displacement. By associating the heat flux with the voltage applied to the PTC ceramics, we were able to quantify the load and ensure that the same amount of displacement was achieved in each experiment.

By changing the shearing device, we captured the speckle patterns under the same heating rate loaded, as shown in Figure 6. With a laser of power 150 mW, the images acquired with RP and WP are shown in Figure 6a,b, respectively.

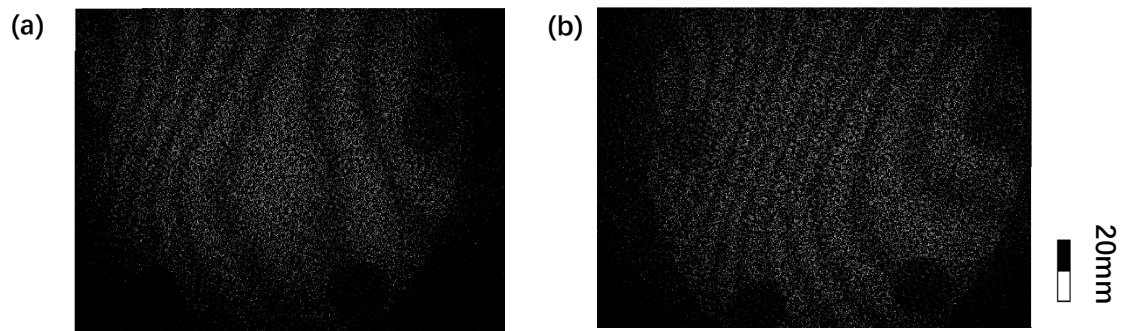


Figure 6. The speckle patterns and the fringe patterns generated by RP and WP: (a) the butterfly fringe patterns of RP; (b) the butterfly fringe patterns of WP.

After ensuring that each experiment was loaded with the same amount of load, we established a model based on the finite element method (FEM) to simulate the theoretical value of the measured displacement, as shown in Figure 7a. By dividing the surface of the measured object into many grid structures, we decomposed the overall changes of the measured object into the sum of local changes. The theoretical displacement gradient on the object surface should be distributed as shown in Figure 7b.

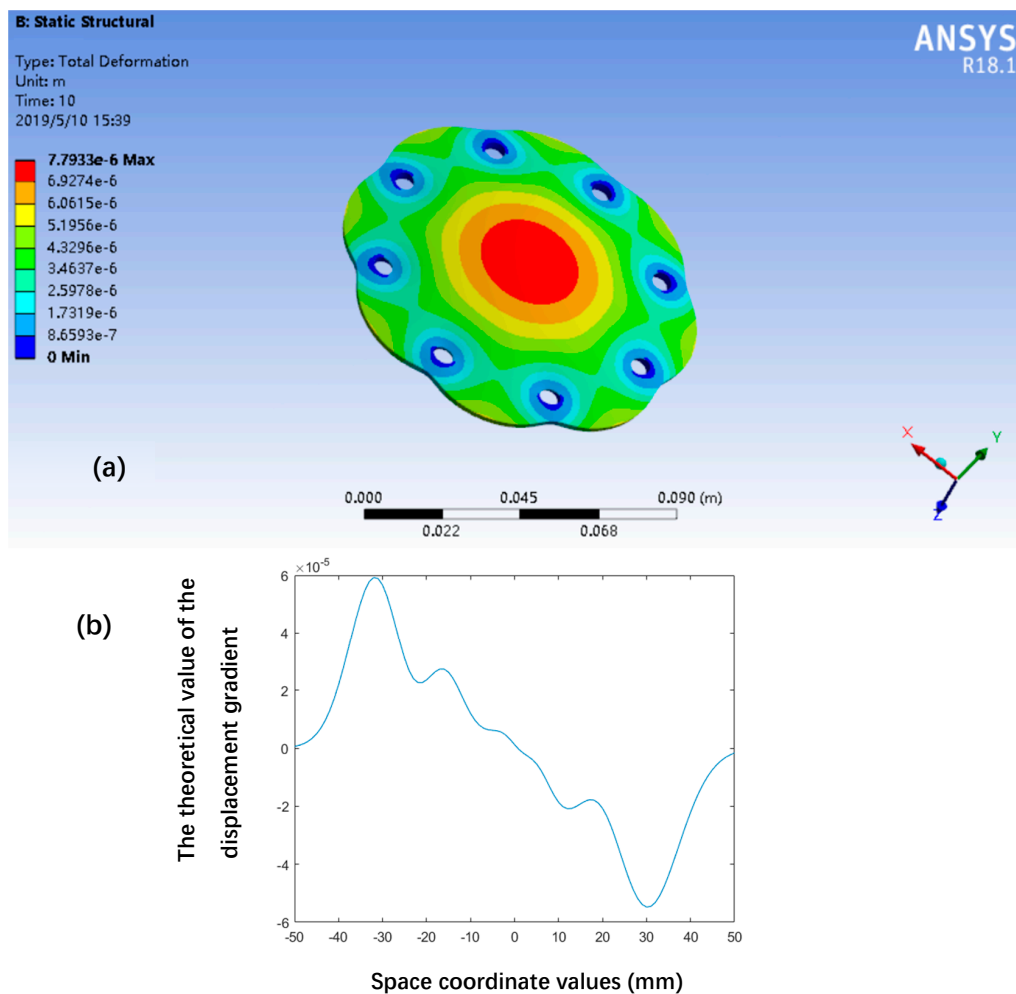


Figure 7. The model based on the finite element method (FEM) and the theoretical values: (a) displacement distribution, with the geometric center as the origin; and (b) the theoretical value of the displacement gradient along the x-axis.

Using WP and RP as shearing devices, we measured the objects under the same FEM. As shown in Figure 8, the measurement accuracies of the system with WP or RP as shear devices were different. The error distribution of WP had an obvious trend with the x-axis, and the overall error was larger than RP. In Figure 8b, the maximum relative measurement error of RP is less than 4%; however, the maximum relative measurement error of WP is about 6% and has a systematic distribution trend.

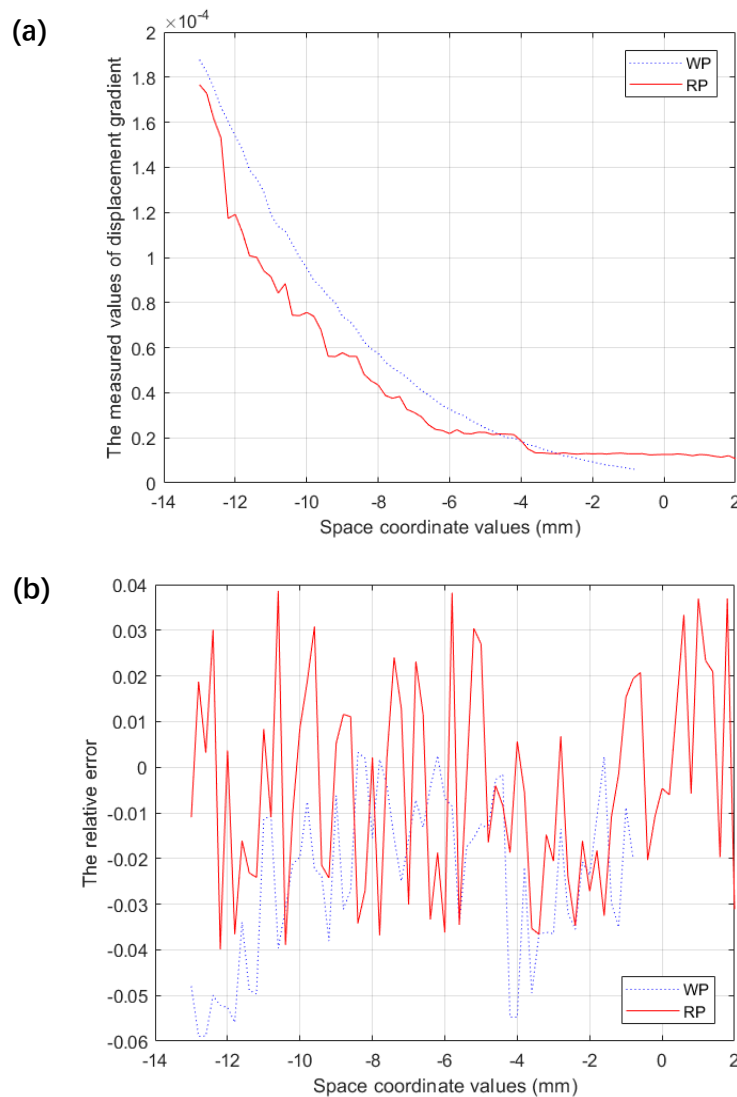


Figure 8. The (a) measured values and (b) relative errors of RP and WP along the x-axis.

The experiment proves that the RP can compensate for the system error caused by background noise. The accuracy of the system is limited by other factors, such as airflow and stray light. However, the improvements in the shearing device provides the possibility of using the shearographic measurement system outdoors, where the power of illumination is generally high.

5. Conclusions

Here, we analyzed the advantages of RP in shearography measurement. First, compared with WP, RP is more suitable for shearography interferometry. Second, RP has a smaller angle of divergence than WP, which ensures both the measurement distance and an appropriate shearing distance. The experimental results showed that replacing WP with RP can improve the accuracy of shearography system, which is expected to be an important step in the practical and instrumental application of the shearography.

Author Contributions: Conceptualization, X.W. and Z.G.; Methodology, X.W.; Software, X.W.; Validation, X.W. and Z.G.; Formal Analysis, J.Z. and Y.N.; Investigation, C.G.; Resources, X.W.; Data Curation, Z.L.; Writing-Original Draft Preparation, X.W.; Writing-Review & Editing, C.G., Y.L. and Z.G.; Visualization, Y.Z.; Supervision, Z.G.; Project Administration, Z.G.; Funding Acquisition, Z.G.

Acknowledgments: The authors acknowledge the financial supported from National Natural Science Foundation of China (Grant No. 51675038).

Conflicts of Interest: The authors declare no conflict of interest.

References

1. Leendertz, J.A.; Butters, J.N. An image-shearing speckle-pattern interferometer for measuring bending moments. *J. Phys. E Sci. Instrum.* **1973**, *6*, 1107–1110. [[CrossRef](#)]
2. Hung, Y.Y. Shearography: A new optical method for strain measurement and non-destructive testing. *Opt. Eng.* **1982**, *21*, 391–395. [[CrossRef](#)]
3. Hanayama, R.; Ishii, K. Surface shape measurement for small lens using phase shift shearing interferometer. *Int. J. Autom. Technol.* **2011**, *5*, 162–166. [[CrossRef](#)]
4. Perea, J.; Nehmetallah, G. Phase-resolved heterodyne shearographic vibrometer for observation of transient surface motion: Theory and model. *Opt. Express* **2017**, *25*, 6169–6181. [[CrossRef](#)] [[PubMed](#)]
5. Kirkove, M.; Guérit, S.; Jacques, L.; Loffet, C.; Languy, F.; Vandenberg, J.; Georges, M. Determination of vibration amplitudes from binary phase patterns obtained by phase-shifting time-averaged speckle shearing interferometry. *Appl. Opt.* **2018**, *57*, 8065–8077. [[CrossRef](#)]
6. Aebischer, H.A.; Waldner, S. Strain distributions made visible with image-shearing speckle pattern interferometry. *Opt. Lasers Eng.* **1997**, *26*, 407–420. [[CrossRef](#)]
7. Moreno, V.; Vázquez-Vázquez, C.; Gallas, M.; Crespo, J. Speckle shearing pattern interferometry to assess mechanical strain in the human mandible-jaw bone under physiological stress. In Proceedings of the International Conference on Applications of Optics and Photonics, Braga, Portugal, 3–7 May 2011.
8. Hiraoka, M.; Oshida, Y.; Iwahashi, Y. Speckle shearing interferometer for measuring strain distribution using laser diodes. In Proceedings of the International Conference on Experimental Mechanics 2014, Singapore, 15–17 November 2014.
9. Francis, D.; Tatam, R.P.; Groves, R.M. Shearography technology and applications: A review. *Meas. Sci. Technol.* **2010**, *21*, 102001. [[CrossRef](#)]
10. Feng, Z.; Gao, Z.; Zhang, X.; Wang, S.; Yang, D.; Yuan, H.; Qin, J. A polarized digital shearing speckle pattern interferometry system based on temporal wavelet transformation. *Rev. Sci. Instrum.* **2015**, *86*, 093102. [[CrossRef](#)] [[PubMed](#)]
11. Wang, X.; Gao, Z.; Qin, J.; Zhang, X.; Yang, S. Temporal heterodyne shearing speckle pattern interferometry. *Opt. Lasers Eng.* **2017**, *93*, 76–82. [[CrossRef](#)]
12. Kubacki, E. Perplexed by polarizers? *Photonics Spectra* **2006**, *40*, 74–78.
13. Feng, Z.; Gao, Z.; Zhang, X.; Wang, S. Research on key factors in the digital shearing speckle pattern interferometry. In Proceedings of the 2015 International Conference on Optical Instruments and Technology: Advanced Lasers and Applications, Beijing, China, 8–10 May 2015.
14. Khaleghi, M.; Dobrev, I.; Harrington, E.; Klausmeyer, P.; Cushman, M.; Furlong, C. Long-term effects of cyclic environmental conditions on paintings in museum exhibition by laser shearography. In *Advancement of Optical Methods in Experimental Mechanics*; Springer: Cham, Switzerland, 2014; Volume 3, pp. 283–288.
15. Chen, X.; Khaleghi, M.; Dobrev, I.; Tie, W.; Furlong, C. Structural health monitoring by laser shearography: Experimental and numerical investigations. In *Experimental and Applied Mechanics*; Springer: Cham, Switzerland, 2015.



© 2019 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>).