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Abstract: The slurry pump is one of the most important pieces of equipment in mineral transportation and separation systems, and it has complex two-phase flow characteristics and wear mechanisms. By employing numerical and experimental methods, the solid–liquid two-phase flow characteristics and wear patterns were investigated in this study. A two-way coupling discrete phase model (DPM) method was used to predict the flow pattern and the wear location and shows good agreement with the experimental observations. The pump performance characteristics of numerical results under pure water conditions were consistent with the experimental results. The effects of particle parameters and operating conditions on the internal flow field and wear were compared and discussed. The results show that the wear degree increased with the increase in volume flow rate and solid volume fraction. With the increase in particle size, the wear range at the impeller inlet became significantly smaller, but the wear degree became obviously larger. This study provides a basis for reducing the wear and improving the hydraulic performance of slurry pumps.

Keywords: slurry pump; separation; wear; discrete phase model; two-way coupling

1. Introduction

Marine minerals have the potential to become a new and important industry. As key equipment in the separation and transportation system of solid–liquid two-phase materials, slurry pumps are widely used in the mining industry and other fields. Slurry pumps also play an important role in channel dredging, coastal construction, and other fields of marine engineering [\[1](#page-19-0)[–4\]](#page-19-1). The presence of solid particles directly leads to the prevalence of high energy consumption, severe wear of flow components, and poor operational reliability of solid–liquid slurry pumps [\[5](#page-19-2)[–8\]](#page-19-3). Due to the complex mechanism of fluid, particle, and material, the design of solid–liquid slurry pumps is often based on a pure water condition and results in the slurry pump having low efficiency and a short service life. At present, the wear mechanisms and design methods of slurry pumps have not been fully grasped. Therefore, the in-depth study of two-phase flow and wear mechanism in solid–liquid slurry pumps to improve their efficiency and service life is the goal of many researchers in the world [\[9,](#page-19-4)[10\]](#page-19-5).

A large number of scholars have used experimental methods to study the wear in slurry pumps. Fan et al. [\[11\]](#page-19-6) analyzed the worn impellers of slurry pumps by scanning electron microscopy, and they concluded that the wear mechanisms were different from each region. The amount of wear of the whole impeller and reduction of the blade thickness was obtained by 3000 h of wear tests. It was proposed that extending the service life of slurry pumps more likely required improving the wear resistance of the material.

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Bhupendra et al. [\[12\]](#page-19-7) found that the amount of wear was less when the pump was operating in its highest efficiency points, and this implied a longer pump operating life. Kadambi et al. [\[13\]](#page-19-8) measured the particle velocity in the impeller channel of a slurry pump by particle image velocimetry (PIV). The highest particle velocities were found to be generated on the suction side of the impeller and in the trailing edge region of the impeller. The particle fluctuation kinetic energy of particles at different speeds was compared to analyzing the wear mechanism inside the impeller of the slurry pump. Khalid et al. [\[14\]](#page-19-9) found the weight losses of the impeller were caused by the impact of solid particles on the surface area. The area near the center of the impeller encountered less wear compared to the area at the edge of the impeller. Madadnia et al. [\[15\]](#page-19-10) investigated the phenomenon of localized inhomogeneous wear in slurry pumps by using electron microscopy and numerical methods. The wear areas included the impeller inlet, inside and outside the impeller shroud, inside the pump body cover, and along the impeller outlet. With the rapid development of computer technology and theoretical mechanics, computational fluid dynamics (CFD) are becoming the most efficient methods to study solid–liquid flow mechanisms. Currently, the commonly used multiphase flow models include the Euler–Euler method and the Euler–Lagrange method. For the Euler–Euler method, Liao et al. [\[16\]](#page-19-11) investigated the effect of particle diameter on the internal flow field and performance of centrifugal pumps by a two-phase particle flow model. The volume fraction distribution of sediment particles and the velocity distribution of sediment particles were analyzed at different sediment particle diameters. Liu et al. [\[17\]](#page-19-12) used the mixture multiphase flow model to study the effect of transporting flow field and particle distribution in the impeller passage of a flue gas desulfurization system with limestone aqueous slurry. It was based on single particles of different concentrations and sizes. It was found that the particle distribution was more severely influenced by particle diameter than solid phase concentration. Wang et al. [\[18\]](#page-20-0) used the Euler–Euler multiphase flow model to study the effect of particle concentration on the conveying performance of slurry pumps. The degradation of pump performance at high volume concentrations was mainly due to increased frictional losses and collisions between particles and between particles and impeller surfaces. Wadnerkar et al. [\[19\]](#page-20-1) used the Eulerian–Eulerian model to simulate an agitated tank containing a solid–liquid mixture and investigated the turbulent energy and cloud height in the agitated tank. The presence of solid phase leads to a weakening of turbulence. The maximum deviation was observed in the impeller plane. Zhang et al. [\[20\]](#page-20-2) used the mixture multiphase model to analyze the effect of particle diameter, volume concentration, and particle density on the hydraulic performance of a low specific speed centrifugal pump. As the particle diameter and solid phase concentration increased, the pump head and efficiency decreased. However, the effect of particle density on the external characteristics of the pump was smaller. Although the Eulerian–Eulerian model is simpler and easier to converge, the calculation cannot reflect the differences of particle density, size, shape, and other physical properties of the particles. Therefore, this model cannot fully and accurately reveal the state of motion, location of the particles, and degree of wear. In summary, the Eulerian–Eulerian model can only obtain the wear status indirectly through the solid volume fraction near the wall and the solid phase slip speed, but cannot obtain the wear status of the pump through the erosion model.

Erosion models based on CFD numerical simulation have become more and more widely used recently. One of the earliest erosion models is the Finnie erosion model [\[21\]](#page-20-3), based on the hypothesis of micro-cutting theory, which believes that when the particle impact angle was low, the tube wall material deformed by cutting. However, the micro-cutting theory was more suitable for a low angle of attack conditions. Grant and Tabakoff [\[22\]](#page-20-4) used the Monte Carlo method to simulate the erosion process of solid particles and liquid droplets in a turbomachine. The simulation only considered the drag force that determined the particle trajectory. By calculating the erosion caused by each collision, the model can predict the amount and location of material erosion. McLaury [\[23\]](#page-20-5) fully considered the erosion caused by direct collisions and random collisions. A CFD-based general erosion prediction process was established and its applicability was verified in two-dimensional

shapes, which is in good agreement with the experimental results. McLaury and Shirazi [\[24\]](#page-20-6) established an erosion model for predicting the maximum permeability of the pipe wall based on CFD for pipes of different shapes. The model was based on a large amount of empirical data from Tulsa University, LSU, Harwell, and DNV for sand movement in multiphase flows. Chen et al. [\[25\]](#page-20-7) added Grant and Tabakoff's particle random rebound model to simulate the erosion severity of elbows and tees. The above wear model studies mainly focused on pipes, elbows, and tees, with less focus on the wear of pumps.

The discrete phase model (DPM) belongs to the Euler–Lagrange method, which is more suitable for particle track study and wear prediction. Pagalthivarthi et al. [\[26\]](#page-20-8) used a two-dimensional model of the pump body with two-way coupling to obtain the dilute slurry flow field by combining the DPM with the standard $k-\varepsilon$ model. The qualitative effects of these parameters on the erosion rate were investigated by varying the pump flow rate, rotating speed, particle diameter, and various geometric conditions. Li et al. [\[27\]](#page-20-9) numerically simulated the solid–liquid two-phase turbulent flow by DPM method. The results showed that the particle diameter has a significant effect on the wear characteristics, and the wear could be reduced by reducing the exit blade angle. Noon et al. [\[28\]](#page-20-10) performed a three-dimensional numerical analysis of wear in a slurry pump and found that the locations subject to the greatest wear were the tongue and passage of the volute. Shen et al. [\[29\]](#page-20-11) used the DPM model to study the screw centrifugal pump that transported the solid–liquid two-phase flow of large particles. The wear distribution under the conditions of solid phase parameters with different density, particle sizes, and solid volume fractions was analyzed. The wear distribution of low-density particles was uniform, and the wear of high-density particles was concentrated. Shen et al. [\[30\]](#page-20-12) analyzed the wear of a double-suction centrifugal slurry pump. By combining the analysis with the scanning electron microscope (SEM) images, the wear of the impeller blades at different locations was obtained. By comparing the wear distribution of numerical calculation with the experimental results, it was concluded that the DPM can be used to predict the wear of the pump. Zhu et al. [\[31\]](#page-20-13) studied the effects of different turbulence models and rebound models on the wear of the three-stage electric submersible pump (ESP) through DPM. The difference between turbulence models was mainly due to the particle incidence angle and impact velocity. The effect of different rebound models can be ignored. Zhu et al. [\[32\]](#page-20-14) also used numerical simulations and wear experiments to compare the accuracy of erosion models in CFD. The wear experiments were to calculate the average erosion rate by measuring the weight loss and verify it with the erosion model. It was concluded that the Oka erosion model had the highest accuracy in the ESP, but the computational resources were the most expensive. Jiang et al. [\[33\]](#page-20-15) conducted numerical simulation and experimental research on the wear characteristics of centrifugal pumps with different solid volume fractions. Through the comparison of wear pictures and numerical simulation results, the accuracy of the numerical simulation was verified, and the energy gradient and particle motion were combined to explain the wear characteristics. However, only one-way coupling between solid–liquid phases were considered in these studies, which ignored the interaction of solid particles and the effects of particle motion on the flow field of the fluid phase.

Two-way coupling not only considers the effect of the flow field on the particle motion but also the effect of the particles on the flow field. In particular, continuous phase flow under the influence of particle motion was analyzed in this study to explain the effect of particles on the continuous phase. This study took a typical centrifugal slurry pump as the research object. The accuracy of the numerical simulation results was verified by experimental methods. The effects of particle parameters and operating conditions on the internal flow field and wear were compared and discussed.

2. Parameters and Methods

2.1. Geometry and Parameters

The 500 MZ centrifugal slurry pump was selected as the object of study. As shown in Table [1,](#page-3-0) its design volume flow rate was $Q = 26 \text{ m}^3/\text{h}$, rotating speed was $n = 1500 \text{ r/min}$, and design head was $H = 11$ m. Therefore, its specific speed was:

$$
n_s = \frac{3.65n\sqrt{Q}}{H^{3/4}} = 77\tag{1}
$$

Table 1. Geometric specifications of the impeller.

Three-dimensional (3D) modeling is based on the real physical model using Unigraphics NX10.0 software (Siemens PLM Software, Plano, TX, USA). The fluid domain model was simplified into four parts: inlet section, impeller, volute, and outlet section, in which
developed. The specific three-dimensional calculation dothe inlet and outlet section were extended to five times the impeller diameter in order to ensure the flow fully developed. The specific three-dimensional calculation domain is ensure the flow fully developed. The specific three-dimensional calculation domain is shown in Figure [1.](#page-3-1) model was simplified into four parts: inlet section, impeller, volute, and outlet section, in

Figure 1. 3D model of the slurry pump.

Figure 1. 3D model of the slurry pump. *2.2. Grids*

used for grids generation. The quality of the grids has great influence on the convergence
and reliability of the computational results. A relatively reasonable number of grids could meet the required accuracy of the simulation and save computational resources. The
hards duel study was a side with a connect was explored and a Fine of a magnetic different grid numbers were selected for grid [in](#page-4-0)dependence analysis. As shown in Table 2, the head and efficiency gradually stabilize as the number of grids gradually increase. After
comparison, it gould be seen that the number of grids atabilizes at 2 million. The total grid number of 3.1 million was adopted for the calculation of single-phase pure water and In this study, ANSYS-ICEM 19.0 software (ANSYS Inc., Canonsburg, PA, USA) was and reliability of the computational results. A relatively reasonable number of grids could hexahedral structural grids with near-wall refinement was employed. Five schemes with comparison, it could be seen that the number of grids stabilizes at 3 million. The total

solid–liquid two-phase flow. The structural grids of the impeller and volute are shown in Figure [2.](#page-4-1)

Figure 2. Grid generation of the computing domain: (**a**) impeller; (**b**) volute. **Figure 2.** Grid generation of the computing domain: (**a**) impeller; (**b**) volute.

2.3. Numerical Methods

2.3.1. Turbulence Model

The standard *k*-*ε* model is the most widely used model for industrial applications. Its model parameters were calibrated by experimental data and have good stability and reasonable accuracy for most applications [\[34\]](#page-20-16). This model implements the Reynolds 1,565,848 11.95 72.12 dissipation rate *ε* into the transport equation. The values for *k* and *ε* are solved by the turbulent kinetic energy dissipation and turbulent energy equations [\[35](#page-20-17)[,36\]](#page-20-18): time-averaged two-equation model by introducing turbulent kinetic energy *k* and turbulent

4,041,363 11.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.78 71.7
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$$
\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_i}(\rho k u_i) = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + G_k + G_b - \rho \epsilon - Y_M + S_k \tag{2}
$$

$$
\frac{\partial}{\partial t}(\rho \varepsilon) + \frac{\partial}{\partial x_i}(\rho \varepsilon u_i) = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_{\varepsilon}} \right) \frac{\partial \varepsilon}{\partial x_j} \right] + C_{1\varepsilon} \frac{\varepsilon}{k} (G_k + C_{3\varepsilon} G_b) - C_{2\varepsilon} \rho \frac{\varepsilon^2}{k} + S_{\varepsilon} \quad (3)
$$

where G_k represents the generation of turbulence kinetic energy due to the mean velocity gradients, G_b is the generation of turbulence kinetic energy due to buoyancy, Y_M represents the contribution of the fluctuating dilatation in compressible turbulence to the overall the values of \mathcal{L} and \mathcal{L} are solved by the largest production of the overall $\frac{d\cos\phi}{dx}$ and $\cos\phi$ and $\frac{d\phi}{dx}$ and $\frac{d\phi}{dx}$ are user-defined source terms. dissipation rate, $C_{1\epsilon}$, $C_{2\epsilon}$, and $C_{3\epsilon}$ are constants, σ_k and σ_ϵ are the turbulent Prandtl numbers

> The turbulent viscosity μ_t can be computed by combining *k* and *ε* as follows: $\frac{1}{2}$ + $\frac{1$

$$
\mu_t = \rho C_\mu \frac{k^2}{\varepsilon} \tag{4}
$$

2.3.2. Discrete Phase Model

In the discrete phase model [\[37](#page-20-19)[,38\]](#page-20-20), the calculation of the particle track is independent. It is scheduled to be completed within the interval specified in the fluid phase calculation. The track of the discrete phase particles is solved by integrating the differential equations for the forces acting on the particles in Lagrangian coordinates. Solid particles moving in solid–liquid two-phase flow field are mainly subjected to the forces of gravity, winding resistance, additional mass force, additional forces induced by the pressure gradient of the flow field, Basset force, Saffman lift, and Magnus lift. The controlling equations of the forces on the particles in the x-direction are listed here as follows:

$$
m_s \frac{du_s}{dt} = \frac{\pi d_s^3 \rho_s}{6} \left[\frac{g_x (\rho_s - \rho_f)}{\rho_s} + F_D + F_V + F_P + F_x \right]
$$
(5)

$$
F_D = \frac{18\mu}{\rho_s d_s^2} \frac{C_D Re_s}{24} \left(u_f - u_s \right) \tag{6}
$$

$$
F_V = \frac{1}{2} \frac{\rho_f}{\rho_s} \frac{d}{dt} \left(u_f - u_s \right) \tag{7}
$$

$$
F_P = \frac{\rho_f}{\rho_s} u_s \frac{\partial u}{\partial x} \tag{8}
$$

where g_x is the x-direction gravitational acceleration; F_D is the fluid bypassing resistance per unit mass of solid particles; *F^V* is the additional mass force per unit mass of solid particles; *F^P* is the additional force induced by the pressure gradient of the flow field; *C^D* is the drag coefficient; and F_x is the sum of other external forces, mainly including Basset force, Saffman lift force, Magnus lift force, and other force.

Saffman lift and Magnus lift caused by the rotation of the particles due to the transverse velocity gradient (shear layer flow) are only valid for subscopic particles (1 to 10 μ m in diameter). Because the particle size of solid particles in the two-phase fluid conveyed by centrifugal slurry pumps is much larger than that of subscopic particles, the effect of the particles' own rotation can be neglected when numerical simulations are performed. The Basset force per unit mass of solid particles in the flow field can be calculated according to the following equation:

$$
F_B = \frac{9}{d_s \rho_s} \sqrt{\frac{\rho_f \mu}{\pi}} \int_0^1 \frac{\frac{du_f}{dt} - \frac{du_s}{dt}}{\sqrt{t - t'}} dt' \tag{9}
$$

In the rotating coordinate system (e.g., rotating impeller zone), the additional forces in the equation also include the centrifugal and Coriolis force exerted on the solid particles due to the impeller rotation. When the impeller rotation axis is the *z*-axis, the additional forces per unit mass of solid particles in the *x* and *y* directions in the Cartesian coordinate system are as follows:

$$
F_{s,x} = \left(1 - \frac{\rho_f}{\rho_s}\right)\omega^2 x + 2\omega \left(u_{y,s} - \frac{\rho_f}{\rho_s}u_{y,f}\right)
$$
(10)

$$
F_{s,y} = \left(1 - \frac{\rho_f}{\rho_s}\right)\omega^2 y + 2\omega \left(u_{x,s} - \frac{\rho_f}{\rho_s}u_{x,f}\right) \tag{11}
$$

where *ux*,*^s* , *uy*,*^s* are the velocities of solid phase particles in *x*, *y* direction; *ux*, *^f* , *uy*, *^f* are the velocities of liquid phase in *x*, *y* direction; and ω is the angular velocity of solid particles rotation. The calculation of force in the *y* and *z* direction is similar to the *x* direction.

2.4. Erosion Model

The McLaury erosion model [\[39\]](#page-20-21) was employed to predict the erosion rate of the solid phase, which is one of the most common and widely used erosion models for dense flow. McLaury erosion rate is expressed as:

$$
E = AVn f(\gamma)
$$
 (12)

$$
A = F(B_h)^k \tag{13}
$$

where V is the particle impact velocity; F is the empirical constant; B_h is the Brinell hardness of the wall material; and *n* is the velocity index.

$$
f(\gamma) = b\gamma^2 + c\gamma \qquad \text{for } \gamma \le \gamma_{\text{lim}} \tag{14}
$$

$$
f(\gamma) = x \cos^2(\gamma) \sin(w\gamma) + y \sin^2(\gamma) + z \qquad \text{for } \gamma > \gamma_{\text{lim}} \tag{15}
$$

where *b*, *c*, *w*, *x*, *y*, *z* are the constants in the collision angle function. In this study, *A =* 2.38810−⁷ , *γlim* = 10, *n* = 1.73, *b* = −34.79, *c* = 12.3, *x =* 0.147, *y* = −0.745, and *w* = 5.205.

2.5. Boundary Conditions

For the single-phase pure water condition, the steady three-dimensional numerical simulation was carried out to predict the flow field and pump performance characteristics. The liquid medium in the computational domain was liquid water with a density of 998.2 kg/m³ , and the gravitational acceleration was set to 9.81 m/s² . The standard *k*-*ε* model was selected as the turbulence model. The boundary condition outlet was set as a mass flow outlet. The inlet boundary condition was set as a pressure inlet. The solid wall surface was set as a no-slip wall surface. Except for the impeller, which was set as rotating at 1500 r/min , all the other computational domains were set as stationary. The dynamic-static computational domain intersection was adopted as the frozen rotor model. The data were exchanged with the computational domains through the interface.

The solid–liquid two-phase flow calculation settings were partially adjusted for the pure water working condition settings. DPM tracks the mass and momentum of the particles along the orbit. These physical quantities can be acted upon in the subsequent calculation of the continuous phase. Thus, whereas the continuous phase affects the discrete phase, the effects of the discrete phase on the continuous phase were also considered. The control equations for the discrete phase and the continuous phase were solved alternately until both converged.

The inlet surface was specified as the particle incidence surface, and the particle incidence velocity was the same as the liquid phase inlet velocity. The solid particles were brown corundum with a density of 2900 kg/m³. The particle diameters were 1.0 mm, 0.8 mm, 0.6 mm, 0.3 mm, 0.15 mm, and 0.1 mm, respectively. The default particle shape was set as regular sphere. The outlet surface of the computational domain was designated as "Escape", and the rest of the wall surface was designated as "Reflect". The singlephase pure water flow field was used as the initial flow field to start the solid–liquid two-phase computation.

The particle collision model is an important part of the DPM model, which is used to describe the interaction between the particles and the wall. In this study, the collision model proposed by Pagalthivarthi and Gupta [\[40\]](#page-20-22) was utilized. The equations are as follows:

$$
\begin{cases}\n e_n = 0.993 - 0.0307\theta + 0.000475\theta^2 - 0.00000261\theta^3 \\
e_t = 0.998 - 0.029\theta + 0.000643\theta^2 - 0.00000356\theta^3\n\end{cases}
$$
\n(16)

where *eⁿ* is the normal restitution coefficient, *e^t* is the tangential restitution coefficient, and θ is the particle impact angle.

3. Experimental Verification and Simulation 3. Experimental Verification and Simulation

- *3.1. Experiments 3.1. Experiments*
- 3.1.1. Performance Experiments 3.1.1. Performance Experiments

The test bench $[41,42]$ $[41,42]$ for the pump performance experiment under pure water conditions is shown in Fig[ure](#page-7-0) 3. A turbine flow meter (accuracy ± 0.1 %) was applied to measure the volume flow rate, and a pressure sensor (accuracy $\pm 0.1\%$) was used to measure the pump inlet and outlet pressure. A torque sensor (accuracy of ± 0.5 %) and speed sensor (accuracy of $\pm 0.1\%$) were employed to measure torque and speed.

$$
H = \frac{p_2 - p_1}{\rho g} + \frac{v_2^2 - v_1^2}{2g} + (z_2 - z_1)
$$
\n(17)

$$
\eta = \frac{\rho g Q H}{P_s} \times 100\% \tag{18}
$$

where p_1 , v_1 , z_1 are the pump inlet parameters, p_2 , v_2 , z_2 are the pump outlet parameters, η is the pump efficiency, and P_s is the output power of the motor. According to the uncertainty calculation procedure of ISO 9906 $[43]$, the total measurement error of this test system is $D_s = \pm 0.37\%$.

Figure 3. Hydraulic performance test bench. **Figure 3.** Hydraulic performance test bench.

3.1.2. Wear Experiments 3.1.2. Wear Experiments

The slurry pump wear test bench [41] is shown in Figure [4](#page-8-0) and consisted of a slurry The slurry pump wear test bench [\[41\]](#page-20-23) is shown in Figure 4 and consisted of a slurry barrel, pump, pressure gauge, pipeline, valve, motor, and test system. The test system barrel, pump, pressure gauge, pipeline, valve, motor, and test system. The test system included electric parts such as a buck starting cabinet, system power distribution cabinet, included electric parts such as a buck starting cabinet, system power distribution cabinet, test conversion protection cabinet, and weak electric parts (such as intelligent display instrument, signal conversion device, and various signal sensors). The total duration of the test was 44.5 h. The operating speed of the slurry pump was 1500 r/min, and solid volume test was 44.5 h. The operating speed of the slurry pump was 1500 r/min, and solid volume fraction was 15%. The solid particles are brown corundum with the density of 2900 kg/m³ and diameter of 0.6 mm. and diameter of 0.6 mm.

The first phase of the experiment lasted for 24 h. When the pump was running for $\frac{1}{4}$ h, it was disassembled and the surface shape of the impeller and volute was inspected. 4 h, it was disassembled and the surface shape of the impeller and volute was inspected. The disassembly inspection revealed that obvious wear marks appeared on the blade inter,
the intersection of the blade suction surface, and the shroud. In the pump disassembly Inspection of the blade suction surface, and the shroud. In the pump disassembly inspection after 24 h operation, it was found that there was serious wear in the pump inspection after 24 h operation, it was found that there was serious wear in the pump impeller inlet and 1–2 mm oval deep pitting in the hub. There was more serious wear in mipeller linet and 1–2 mm oval deep priding in the hub. There was more serious wear in the intersection of the impeller blade suction surface and the shroud, the blade showed an the intersection of the impeller blade suction surface and the shroud, the blade showed an approximately 3 mm wear groove. The second phase of the experiment lasted for 20.5 h. approximately 3 mm wear groove. The second phase of the experiment lasted for 20.5 h. After the second phase of the wear experiment, the blade at the impeller inlet was badly The disassembly inspection revealed that obvious wear marks appeared on the blade inlet, worn. The impeller hub was almost destroyed.

worn. The impeller hub was almost destroyed.

3.2. Comparison and Verification The performance comparison of calculation results and experimental results is

3.2.1. Performance Comparison under Pure Water Conditions

in Figure 5. The simulated head and efficiency variation trends were in high agreement with the experimental values. The minimum deviation between the simulated and experimental
heads at the volume flow rate of $26 \text{ m}^3/\text{h}$ was only 0.02%. The deviation of efficiency near the design working condition was within 2.25%. The deviation of head calculation in this $\epsilon_{\rm M}$. full flow range was basically within 5%, so the calculation model could accurately predict
the pump performance. The performance comparison of calculation results and experimental results is shown the experimental values. The minimum deviation between the simulated and experimental the pump performance.

Figure 5. Pump performance curve comparison under pure water conditions. **Figure 5.** Pump performance curve comparison under pure water conditions.

3.2.2. Wear Comparison

The comparison of wear at the impeller inlet is shown in Figure [6.](#page-9-0) From the simulation results, it could be seen that there was a high degree of wear in the junction between the leading edge of the blade inlet and the hub. There were scattered wear marks in the shroud. In the experiments, the wear at the impeller inlet was mainly concentrated on the leading edge of the blades, especially arou[nd](#page-9-0) the hub side. The red mark on Figure 6 shows the leading edge of the impeller inlet had serious wear. The main reason might be that the particles enter the impeller and move toward the hub due to inertia [44], resulting in more severe wear at the junction of the impeller inlet and the hub than at the shroud. The wear at the impeller inlet was relatively consistent between the simulation and experiments.

Figure 6. Wear comparison at impeller inlet: (a) simulation; (b) experiment.

The wear comparison at the blade pressure surface is shown in Figures 7 and 8. [Fr](#page-9-1)om the simulation result, it could be seen that the junction of the front half of the blade and the
half of the blade and high seconds. From the half of the blade late the front line of the the ground position shifted from the hub to the middle part. The experimental result shows that the blade near the hub side had serious wear, with a worn groove in the middle of the blade. The wear at the impeller was mainly concentrated on the leading edge of the blade closer to the hub side. By comparing Figures 7 and 8, it c[ou](#page-9-1)ld be concluded that the [w](#page-10-0)ear at the blade pressure surface and impeller outlet was relatively consistent between the simulation and experiments. hub showed higher wear. From the leading edge of the blade to the trailing edge, the wear hub showed higher wear. From the leading edge of the blade to the trailing edge, the wear the simulation result, it could be seen that the junction of the front half of the blade and

Figure 7. Wear comparison at blade pressure surface: (a) simulation; (b) experiment.

Figure 8. Wear comparison at middle section of blade pressure surface: (**a**) simulation; (**b**) experiment.

section to the eighth section of the volute, only the volute wall close to the shroud showed wear. In the experimental results, the wear area gradually became larger from the first section to the eighth section. In particular, the middle of the volute wall surface showed the most serious wear. The simulation results were consistent with the experimental results in terms of wear location. The wear comparison at the volute i[s](#page-10-1) shown in Figure 9. The simulation result shows that the whole volute wall surface near the shroud was heavily worn. From the second

cross-section; (c) experiment, sections I to II; (d) experiment, sections IV to V; (e) experiment, sections **Figure 9.** Wear comparison at middle section of blade pressure surface: (**a**) simulation; (**b**) volute cross-section; (**c**) exper-Figure 9. Wear comparison at middle section of blade pressure surface: (a) simulation; (b) volute VI to VIII.

4. Results and Discussion *4.1. Flow Field Analysis*

4.1. Flow Field Analysis

4.1. Flow Field Analysis
4.1.1. Pressure Distribution

The pressure distribution of the pump cross-section with different solid volume fractions under 1.0 Q_{BEP} is shown in Figure 10. The term Q_{BEP} represents the flow rate of the design operating condition. It was found that the pressure distribution in the cross-section of the volute under different solid volume fractions was relatively similar to the distribution in pure water conditions. The pressure at the impeller passage increased from the inlet to the outlet. Compared with the impeller passage, the pressure inside the volute was more influenced by the solid-liquid two-phase flow. A certain degree of pressure drop in the volute was caused by the presence of solid particles. As the solid volume fraction increased, the particles in the volute hit the wall more frequently, resulting in a large amount of energy loss. Therefore, the total pressure of the liquid that entrained with particles gradually decreased. Especially when the solid volume fraction was 15%, there was a more significant drop of the outlet pressure in the volute outlet diffusion section.

Figure 10. Pressure distribution of pump cross-section with different solid volume fractions under **Figure 10.** Pressure distribution of pump cross-section with different solid volume fractions under 1.0Q_{BEP}: (a) $C_V = 0\%$; (b) $C_V = 5\%$; (c) $C_V = 10\%$; (d) $C_V = 15\%$.

4.1.2. Velocity Streamlines 4.1.2. Velocity Streamlines

The inlet section streamline distributions of the slurry pump at 0%, 5%, 10%, and 15% in the inlet section streamline distributions of the slurry pump at 0%, 5%, 10%, and 15% solid volume fraction with particles of 0.6 mm diameter under 1.0*Q_{BEP}* are shown in Figure [11.](#page-12-0) Under the pure water conditions, the liquid entering the impeller would be affected be a fact the pure water conditions, the liquid entering the impeller would be a fact that the state of the state of the sta affected by the rotation of the impeller. The original axial flow then became radial flow due to the impeller rotating effects. Slight prewhirl occurred at the impeller inlet section and its direction was consistent with the direction of impeller rotation. When T_{rel} is alleged added to the liquid, the particles improved the viscosity of the medium. The backflow at the front chamber was intensified and the approximation at Thus of the ability of its direction was consistent with the direction of impeller rotation. When particles were the front chamber was intensified, and the prewhirl was strengthened. Thus, the ability of the impeller to drive the prewhirl near the inlet was enhanced compared with pure water

conditions. As solid volume fraction increased, the prewhirl became more intense. Under the 15% solid volume fraction condition, there was a stronger prewhirl in the impeller inlet. ditions. As solid volume fraction increased, the prewhirl became more intense. Under the following a solid volume fraction increased, the prewill increased because interest.

Figure 11. $\sqrt{8}$. $\sqrt{9}$ / $\sqrt{9}$ / $\sqrt{9}$ / $\sqrt{9}$ / $\sqrt{9}$ / $\sqrt{9}$ / $\sqrt{9}$ under 1.0 Q_{BEP} : (a) $C_V = 0\%$; (b) $C_V = 5\%$; (c) $C_V = 10\%$; (d) $C_V = 15\%$. **Figure 11.** Velocity streamline distributions in the inlet section at different solid volume fractions

Fraction are shown in Figure [12.](#page-13-0) Similar to the pure water conditions, the velocity of the flow field increased from the impeller inlet to the outlet. In the impeller passage near the suction surface, a large vortex appeared. Compared with the pure water conditions, there suction surface, a large vortex appeared. Compared with the pure water conditions, there was a decreasing trend of the vortex strength in the impeller passage under the solid–liquid condition. The reason might be that the total mass of the particles is larger, so the liquid rotational motion needs to overcome the larger inertia of particle motion with more energy consuming. The vortex development is somewhat inhibited by particle motion. Under the design operating conditions, the velocity streamline distributions at the cross-section of the volute was approximately identical under different solid volume fractions. There were more streamlines back into the spiral section passage at the location of the volute tongue. The velocity streamline distributions in the cross-section under different solid volume Due to the large inertia of particle motion and more frequent collision between the particles

and the wall, a certain flow separation occurred at the downstream of the volute diffusion section. The flow separation was more serious under the larger solid volume fraction. larger solid volume fraction.

stream of the volute diffusion section. The volute diffusion section serious under the flow separation was more
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Figure 12. Velocity streamline distributions of impeller-volute with different solid volume fractions **Figure 12.** Velocity streamline distributions of impeller-volute with different solid volume fractions under 1.0 Q_{BEP} : (a) C_V = 0%; (b) C_V = 5%; (c) C_V = 10%; (d) C_V = 15%.

the blade. With the liquid movement in the impeller passage, the vortex developed to the trailing of the suction surface $[44]$. The blade flow separation led to the occurrence of the Γ shown in Figure [13.](#page-14-0) The inter-impeller vortex had a large impact on the particle trajectory. The particles were subjected to great resistance at the leading edge of the pressure surface. Thus, the particles hit the wall at the leading edge of the pressure surface. The wear range The particles were affected by the vortex at the suction surface and hit the wall more frequently. Therefore, the wear of the suction surface concentrated on the middle and rear of the suction surface, and the degree of wear was larger. Compared with other impeller passages which it vortex, the impeller passages which the strong vortex showed less wear. It is because the blockage of the impeller passages by the vortex causes fewer particles to flow through the impeller passage, resulting in less frequent particle collision with the wall. The inter-impeller vortex started from the pressure surface near the leading edge of inter-impeller vortex. The vortex clockwise direction rotation in the impeller passage is of the pressure surface was concentrated on the leading edge, and the degree was slight. passages without vortex, the impeller passages with the strong vortex showed less wear. It

4.2. Wear Analysis

4.2.1. Effects of Volume Flow Rates

Figure [14](#page-14-1) shows the pump head and efficiency under different volume flow rates. The pump performance curve under pure water and 5% solid volume fraction conditions were compared. With particles, the head and efficiency drop faster than under pure water conditions. The reason might be that with the increase in the volume flow rate, the collisions between the particles and between the particles and the wall cause additional energy losses. The impeller consumes more energy under the solid–liquid two-phase flow condition. Figure [15](#page-15-0) shows the wear distribution of each part under different volume flow rates, with a particle diameter of 0.6 mm and solid volume fraction of 5%. As the volume flow rate increased, the wear range and degree of the impeller inlet became significantly larger. Remarkable wear appeared at the blade leading edge and the shroud. The wear on the blades near the hub was larger than at other parts. The reason might be as the volume flow rate increased, the flow velocity increased, and more solid particles move with the liquid from the shroud to the hub [\[44\]](#page-20-26).

Figure 13. Velocity streamlines and wear in the impeller.

Figure 14. Pump head efficiency under different volume flow rates ($d_p = 0.6$ mm, $C_V = 5\%$).

4.2.2. Effects of Particle Diameters

The comparisons of pump head and efficiency with different particle diameters $(d_p = 1.0$ mm, $d_p = 0.8$ mm, $d_p = 0.6$ mm, $d_p = 0.3$ mm, $d_p = 0.15$ mm, and $d_p = 0.1$ mm) under 5% solid volume fraction are shown in Figure [16.](#page-16-0) The particle diameter shows limited effects of on the head and efficiency. Because the particle volume fraction keeps as a constant of 5%, the total mass of the particles is still same. In addition, the increase in particle diameter led to the reduction of the particle number for the same volume fraction. The negative impact on increasing particle diameter on the pump head and efficiency is offset. Therefore, under the same small volume fraction, particle diameter had relatively limited influence on the external characteristics of the pump.

Figure 15. Wear distribution of each part at different volume flow rates: Impeller inlet: (a) $0.6Q_{BEP}$ **(b)** 1.0Q_{BEP}; **(c)** 1.4Q_{BEP}; Blades: **(d)** 0.6Q_{BEP}; **(e)** 1.0Q_{BEP}; **(f)** 1.4Q_{BEP}; Volute: **(g)** 0.6Q_{BEP}; **(h)** 1.0Q_{BEP} $-\infty$. Effects of Particle Diameters of Particle Diameters of Particle Diameters of Particle Diameters of P $(i) 1.4Q_{BFP}.$

Figure 17 shows the wear distribution of each part under different particle diameters, with volume flow rate of 26 m³/h and solid volume fraction of 5%. When the solid particles entered the impeller, they were affected by the inertia and blade displacement, which caused the collision between the particles and the hub at the impeller inlet. With the increase in particle diameters, the wear range of the impeller inlet became significantly

smaller, but the wear degree was markedly larger. Moreover, there was an increase in shaner, but the wear degree was markedifful anger. Moreover, there was an increase in
degree of wear on the blade leading edge. The wear at the pressure surface was changed from closer to the shroud to closer to the hub. Because of the increase in particle diameter, the particle mass increased, resulting in a stronger centrifugal force on the particles. It led the particle mass increased, resulting in a shortger centrified force on the particles. It is determined to a larger wear range and more severe wear on the pressure surface and the hub. There was more wear on the entire volute wall surface near the shroud. As the particle diameter $\frac{1}{2}$ became smaller, the entire volute wall wear area moved closer to the hub. The reason might be that as the particle diameter decreased, the inertia of the particles is smaller and the influence of the liquid centrifugal force on the solid particles was reduced. the particle mass increased, resulting in a stronger centrifugal force on the particles. It led

Figure 16. Pump performance characteristics curve at different particle diameters (1.0*QBEP*, *CV* = 5%). **Figure 16.** Pump performance characteristics curve at different particle diameters (1.0 Q_{BEP} , C_V = 5%).

4.2.3. Effects of Solid Volume Fraction

As can be seen in Figure [18,](#page-17-1) the performance curves of the pump under different solid volume fractions show remarkable differences compared with the pure water condition. The increase in solid volume fraction led to a bigger decrease in pump head and efficiency. The effect of smaller solid volume fraction on the flow field was weak. When the particle volume fraction $C_V = 15\%$, the head and efficiency drop reached their maximum. Due to the addition of particles, the density of the mixed liquid is higher than the density of pure water. Under the same conditions, the impeller should expend more energy to transport the solid–liquid two-phase flow. Therefore, the increase in solid volume fraction would make the head and the efficiency drop significantly.

Figure [19](#page-18-0) shows the wear distribution of each part under different solid volume fractions, with volume flow rate of 26 $\rm m^3/h$ and particle diameter of 0.6 mm. With the increase in solid volume fraction, the wear range at impeller inlet had tiny changes, but the wear degree became larger. Moreover, the wear degree on the blade leading edge and the shroud also increased. On the blade surface, with the increase in solid volume fraction, the wear basically showed similar distribution characteristics, but the wear degree gradually became larger. There was severe wear on the entire volute wall surface near the shroud. At the eighth section of the volute, the wall surface was worn entirely.

Figure 17. Wear distribution of each part under different particle diameters: Impeller inlet: (a) $d_p = 0.6$ mm; (b) $d_p = 0.3$ mm; (c) $d_p = 0.15$ mm; Blades: (d) $d_p = 0.6$ mm; (e) $d_p = 0.3$ mm; (f) $d_p = 0.15$ mm; Volute: (g) $d_p = 0.6$ mm; (h) $d_p = 0.3$ mm; (i) $d_p = 0.15$ mm.

18. Pump performance characteristics curve under different solid phase volume fractions **Figure 18.** Pump performance characteristics curve under different solid phase volume fractions
(O = 1 0O_{nnn} d = 0.6 mm) $t_{\rm F}$ basic showed similar distribution characteristics, but the wear degree grad-term ϵ $F_1(t) = 100$ per $f_2(t) = 0.6$ mm) $(Q = 1.0Q_{BEP}, d_p = 0.6$ mm).

Figure 19. Wear distribution of each fraction at different solid volume fractions: Impeller inlet: (a) $C_V = 5\%$; (b) $C_V = 10\%$; (c) $C_V = 15\%$; Blades: (d) $C_V = 5\%$; (e) $C_V = 10\%$; (f) $C_V = 15\%$; Volute: $5. - v$ control v control v (**g**) $C_V = 5\%$; (**h**) $C_V = 10\%$; (**i**) $C_V = 15\%$.

$\sum_{i=1}^n$ this study, we are $\sum_{i=1}^n$ type centrified simulations on a 500 MZ type centrified simulations on a 500 **5. Conclusions**

In this study, wear experiments and numerical simulations on a 500 MZ type centrifugal slurry pump were carried out. The flow characteristics and wall wear under pure water and solid–liquid two-phase conditions were studied and compared. Two-way coupling discrete phase model with the standard *k-ε* turbulence model were used. The numerical results were validated by experimental results. Most importantly, the effects of particle parameters and pump operating conditions on the internal flow field and wear were investigated. The main results can be concluded as follows:

(1) The head and efficiency based on numerical predictions agrees well with the experimental values. At the design volume flow rate, the difference between the simulation and test values of head and efficiency is less than 2.25%. Thus, the reliability of the numerical calculation method is verified.

(2) As solid volume fraction increases, the prewhirl becomes more severe at the impeller inlet. The flow separation is more serious at the downstream of the volute at a larger solid volume fraction. Compared with the pure water conditions, the intensity of the

vortex in the impeller passage has a certain degree of decreasing trend, which is caused by the particle inhibition. In addition, the particle movement is affected by the vortex. The wear on the pressure surface and suction surface is more serious than the wear on the impeller shrouds.

(3) With the increase in volume flow rate or solid volume fraction, the head and efficiency decline. Under the same solid volume fraction, particle diameter has a relatively limited impact on the pump performance. As the volume flow rate increases, the degree and range of wear at the impeller inlet become significantly larger. As the particle diameter decreases, the impeller inlet wear range becomes considerably larger, but the wear degree changes only slightly. The wear shows remarkable increase as the solid volume fraction increases.

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