

## Article

# A Compact Fourth-Order Tunable Bandpass Filter Based on Varactor-Loaded Step-Impedance Resonators

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**Abstract:** In this paper, a compact high-selectivity frequency tunable bandpass filter (BPF) with constant absolute bandwidth (ABW) based on varactor-loaded step-impedance resonators (SIRs) is presented. By introducing cross coupling between resonators, a pair of transmission zeros (TZs) close to the passband are produced and the selectivity of the filter is enhanced significantly. Another pair of TZs are generated to improve the out-of-band rejection by using source-load coupling. The varactor-loaded SIRs are utilized to design the compact fourth-order tunable BPF in order to realize wide tuning range and compact size. In addition, the frequency-dependent coupling feeding structures are employed instead of lumped capacitors used in conventional feeding structures, as a result, the insertion-loss performance is improved. The simulated and measured results are presented and show good agreement. The measured results exhibit a tuning range from 0.8 to 1.14 GHz with a 3 dB constant ABW of about  $47 \pm 5$  MHz, the return loss of the filter is greater than 13.9 dB, and the insertion loss is about 2.7–3.1 dB. Moreover, four TZs are generated, and the proposed tunable filter shows high selectivity with a rectangular coefficient of 2.3–3.1.

**Keywords:** constant absolute bandwidth; step impedance resonator; tunable bandpass filter; varactor diode



**Citation:** Li, S.; Li, S.; Yuan, J. A Compact Fourth-Order Tunable Bandpass Filter Based on Varactor-Loaded Step-Impedance Resonators. *Electronics* **2023**, *12*, 2539. <https://doi.org/10.3390/electronics12112539>

Academic Editors: Tonino Pisanu, Paolo Maxia and Giacomo Muntoni

Received: 17 May 2023

Revised: 30 May 2023

Accepted: 2 June 2023

Published: 5 June 2023



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## 1. Introduction

With the development of modern wireless communication systems, more frequency bands are utilized, and RF front-ends need to operate at multiple frequency bands [1]. The traditional RF front-end solution termed filter bank to achieve multi-band compatibility generally occupy large area on circuit board and have high costs [2]. The electrically tunable filter which can support numerous frequency bands with a single filter structure may meet the requirements of miniaturization, integration and low cost, is one of the promising microwave components for future wireless communications [3].

Recently, numbers of electrically tunable bandpass filters have been proposed. Two-pole tunable bandpass filters based on uniform-impedance resonators (UIRs) are introduced in [4,5]. Two tunable transmission zeros are generated due to the source-load coupling architecture. In [6,7], two-pole frequency tunable filters based on mixed coupling structures are proposed, by adjusting the mixed coupling structures, the constant bandwidth characteristic is achieved. In addition to these tunable filters based on coupled resonators, dual-mode resonators are also used to design frequency tunable filters in [8] and [9]. In [10], a frequency tunable bandpass filter based on coupled line is proposed, the proposed tunable filter has a simple structure and wide tuning range. However, these aforementioned low-order tunable filters (with two poles) usually suffer from poor selectivity. In order to improve the selectivity of filters, several high-order tunable filters are presented in [11–16]. In [11], a miniature tunable four-pole bandpass filter is demonstrated, the selectivity of the proposed filters is improved significantly compared with two-pole filters. Nevertheless, the absolute bandwidth of the filter cannot remain constant throughout the tuning range. In [12], defected ground structures (DGSs) are applied in the coupling schemes of the

presented planar tunable filter to achieve constant ABW, while the selectivity of the filter is inadequate. In [13], a cross-coupled tunable bandpass filter based on full-wavelength tunable resonators is proposed. The proposed tunable filter is comprised of two cascaded-trisection structures, and two transmission zeros are generated for improving the frequency selectivity; however, the structure of the proposed filter is complicated, and the insertion loss is large. In [14–16], fourth-order tunable filters with constant ABW are demonstrated; however, in order to maintain constant ABW, lumped capacitors are used to connect the feeding microstrip lines to the filter, which may increase the insertion loss of the tunable filters. In [17,18], fourth-order tunable BPFs with high selectivity are designed. The proposed tunable BPFs are manufactured by low temperature co-fired ceramic (LTCC) technology, as a result, the integration of the filters is enhanced, and the filters are compact, but the tuning range of the filters is not good. In [19], an electrical tunable microstrip bandpass filter with cross coupling is proposed. The proposed tunable BPF cannot maintain constant ABW during the tuning range. In [20], a tunable fourth-order BPF with constant ABW is presented. The filter is based on varactor-loaded UIRs, and the selectivity of the proposed filter is enhanced significantly due to the cross-coupling structures; however, the tuning range of the filter also need to be improved. Thus, it is still a challenge to design a compact fourth-order tunable filter with high selectivity, wide tuning range and constant ABW.

In this paper, a compact frequency tunable bandpass filter with high selectivity, wide tuning range and constant ABW based on varactor-loaded SIRs is proposed. The proposed tunable filter is comprised of four varactor-loaded SIRs. To enhance selectivity of the filter, the cascaded quadruplet (CQ) filter topology is utilized, and two transmission zeros located close to the passband are generated. Meanwhile, cross coupling between the input and output feeding microstrip lines is introduced to generate two additional transmission zeros. By regulating the ratio between electric coupling and magnetic coupling and making the overall coupling coefficient inversely proportional to frequency, constant ABW of the proposed filter is realized. In addition, the proposed tunable filter is compact as a result of the application of the varactor-loaded SIR. The insertion loss of the proposed tunable BPF is improved due to the application of the frequency-dependent coupling-feeding structures.

## 2. Theory of Tunable BPF with Constant ABW

According to the classical filter synthesis theory, transmission zeros at finite frequencies can be generated by introducing cross coupling between resonators [21], and the selectivity of bandpass filter can be improved effectively by arranging the location of the transmission zeros properly. In this paper, the typical CQ filter topology is employed to enhance the selectivity of the tunable filter. In addition, the source-load coupling is introduced to produce additional two transmission zeros. Figure 1 illustrates the topology of the proposed filter. Each green node represents a resonator, the resonant frequencies of resonators are electronically controlled by varactor diodes. The S and L node represent the source and load of the tunable filter. The solid lines indicate the main path couplings, and the dashed line denotes the cross couplings. E and M represent electric and magnetic coupling, respectively.

The general normalized coupling matrix denoted by  $[m]$  is employed to represent the proposed tunable filter topology. The coupling coefficient between the resonator  $i$  and  $j$  ( $i, j = 1, 2, 3, 4$ ) is denoted by  $m_{ij}$ ,  $m_{S1}$  represents the input couplings between the source, and resonator 1,  $m_{4L}$  represents the output coupling between the load and resonator 4, and  $m_{SL}$  represents the input couplings between the source and load. For the sake of simplicity, the diagonal entries of the general coupling matrix are all zeros assuming all resonators resonating at the same frequency here. In addition,  $m_{13}$  and  $m_{24}$  are also zeros, which indicates that the coupling between the nonadjacent resonators is zero, with the exception of the cross coupling between resonators 1 and 4.

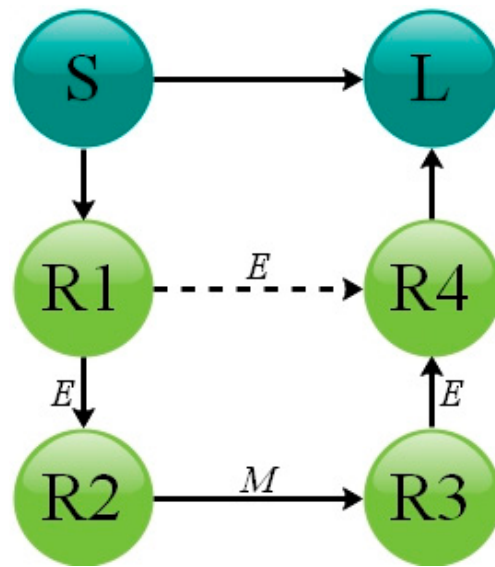


Figure 1. Coupling topology of the proposed filter.

$$[m] = \begin{bmatrix} 0 & m_{S1} & 0 & 0 & 0 & m_{SL} \\ m_{S1} & 0 & m_{12} & 0 & m_{14} & 0 \\ 0 & m_{12} & 0 & m_{23} & 0 & 0 \\ 0 & 0 & m_{23} & 0 & m_{34} & 0 \\ 0 & m_{14} & 0 & m_{34} & 0 & m_{4L} \\ m_{SL} & 0 & 0 & 0 & m_{4L} & 0 \end{bmatrix} \quad (1)$$

then the normalized coupling matrix [m] can be mapped to the center frequency  $f_c$  and scaled by fractional bandwidth (FBW). Thus, the denormalized coupling matrix [M] can be obtained from the following equation:

$$[M] = [m] \times \frac{ABW}{f_c} \quad (2)$$

$Q_{eS}$  and  $Q_{eL}$  are the external quality factors in association with the input and output couplings, respectively.  $Q_{eS}$  and  $Q_{eL}$  can be expressed as follows:

$$\begin{aligned} Q_{eS} &= \frac{f_c}{m_{S1}^2 \times ABW} \\ Q_{eL} &= \frac{f_c}{m_{4L}^2 \times ABW} \end{aligned} \quad (3)$$

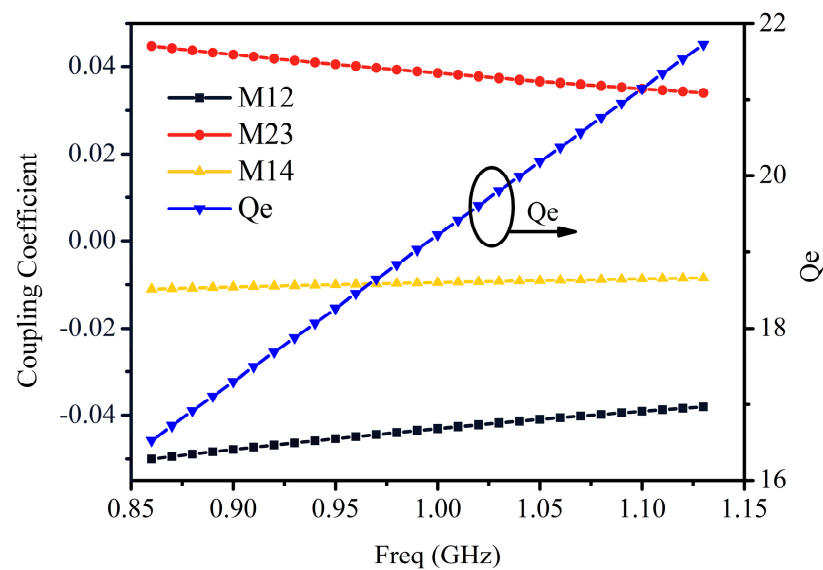
for frequency tunable filters with constant ABW. In order to make filter response of the tunable filter remain stable, the normalized coupling matrix [m] of the filter should remain unchanged [22]. Thus, from Equations (1)–(3), it can be shown that the coupling coefficient  $M_{ij}$  between resonators is inversely proportional to the frequency, and the external quality factor  $Q_{eS}$  and  $Q_{eL}$  is directly proportional to the frequency.

In order to illustrate the synthesis procedure for frequency tunable filter with constant ABW, a design example of a fourth-order tunable filter is demonstrated here. For a frequency tunable filter with ABW 50 MHz, the tuning range is 0.8–1.1 GHz, maximum in-band return loss of the filter is 20 dB. To enhance the selectivity and stopband performance of the filter, two pairs of transmission zeros at the normalized resonant frequency  $\pm 1.7$  and  $\pm 25$  are designed. Based on the method introduced in [23], the normalized coupling matrix

[m] can be obtained using the predefined return loss and transmission zero distribution. The matrix [m] is expressed as follows:

$$[m] = \begin{bmatrix} 0 & 1.02 & 0 & 0 & 0 & 0.0005 \\ 1.02 & 0 & -0.86 & 0 & -0.19 & 0 \\ 0 & -0.86 & 0 & 0.77 & 0 & 0 \\ 0 & 0 & 0.77 & 0 & -0.86 & 0 \\ 0 & -0.19 & 0 & -0.86 & 0 & 1.02 \\ 0.0005 & 0 & 0 & 0 & 1.02 & 0 \end{bmatrix} \quad (4)$$

The de-normalized coupling matrix and external quality can be obtained by Equations (1)–(3). Thus, the  $M_{12}$ ,  $M_{23}$ ,  $M_{14}$  and  $Q_e$  can be expressed as shown in Figure 2. Because of the symmetric distribution of the filter,  $M_{34} = M_{12}$ ,  $Q_e = Q_{eS} = Q_{eL}$ .



**Figure 2.** Theoretical coupling coefficient and external quality for the predefined tunable filter based on the synthesis method introduced before.

Based on the coupling matrix [M] and external quality  $Q_e$  calculated before, the filter response  $S_{11}$  and  $S_{12}$  can be obtained by Equations (5)–(7).

$$S_{11} = 1 + 2j[A]_{11}^{-1} \quad (5)$$

$$S_{21} = -2j[A]_{n+2,1}^{-1}$$

$$A = [m] + \Omega[U] - j[q] \quad (6)$$

$$\Omega = \frac{1}{FBW} \left( \frac{w}{w_0} - \frac{w_0}{w} \right) \quad (7)$$

where [U] is similar to the  $(n + 2) \times (n + 2)$  identity matrix, except that  $[U]_{11} = [U]_{n+2, n+2} = 0$ , [q] is the  $(n + 2) \times (n + 2)$  matrix with all entries zeros, except for  $[q]_{11} = [q]_{n+2, n+2} = 1$ .

The filter response of the tunable filter is shown in Figure 3. It can be seen that during the whole tuning range, and the  $S_{11}$  and  $S_{12}$  curves remain stable, the absolute bandwidth remains 50 MHz, two transmission zeros close to the passband are generated, and the selectivity of the proposed filter is improved significantly, besides, another two transmission zeros located far from the passband are also produced, and this two transmission zeros are very useful to enhance the rejection level of stopband.

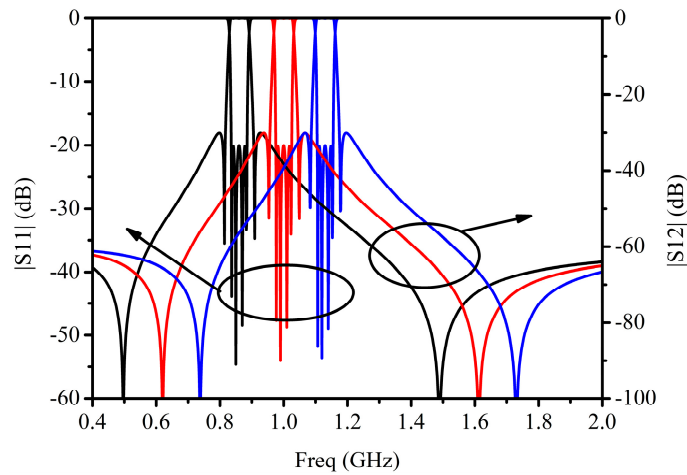


Figure 3. Frequency response of the tunable filter during the whole tuning range.

### 3. Analysis of the Varactor-Loaded SIR

In compared to conventional UIR, filters based on SIRs usually have a smaller size and superior stopband performance. Therefore, varactor-loaded SIR is employed to design tunable filters. Figure 4 shows the layout of the proposed varactor-loaded SIR, the SIR consists of two transmission lines, whose electrical length and characteristic admittance are  $\theta_1, \theta_2$  and  $Y_1, Y_2$ , respectively. The varactor diode is loaded at the position  $\theta_c$  away from the short end of the SIR.  $C_v$  is the effective capacitance of varactor diode.

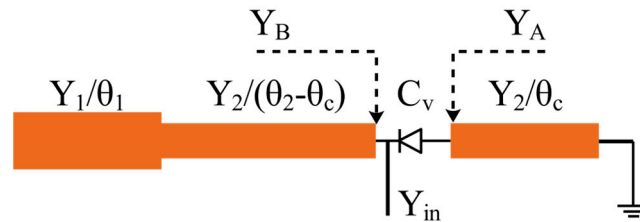


Figure 4. Layout of the proposed tunable SIR.

Based on transmission line theory in [23], the following equations can be used to estimate the input admittance  $Y_{in}$ .

$$Y_A = \frac{Y_2}{j \tan \theta_c} \tag{8}$$

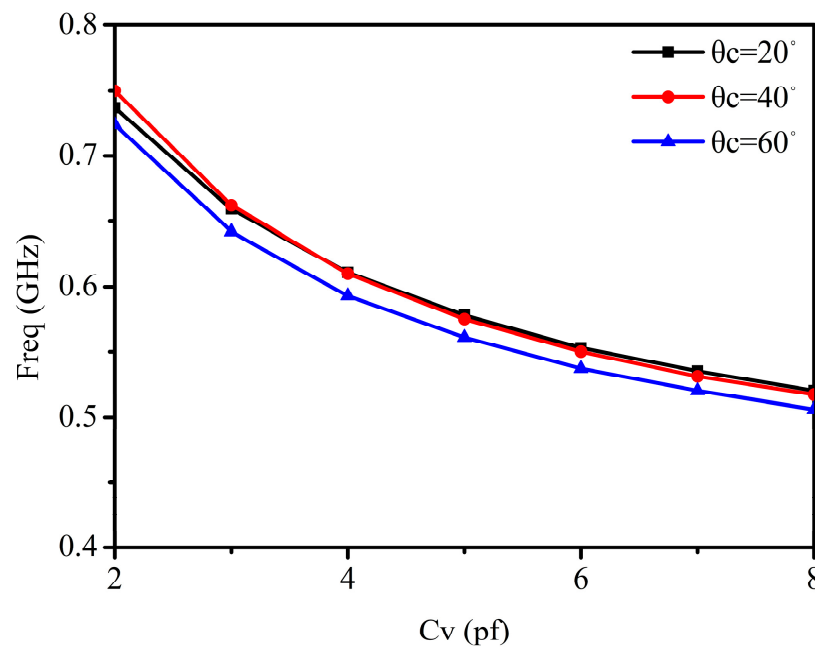
$$Y_B = \frac{jY_1Y_2 \tan \theta_1 + jY_2^2 \tan(\theta_2 - \theta_c)}{Y_2 - jY_1 \tan \theta_1 \tan(\theta_2 - \theta_c)} \tag{9}$$

$$Y_{in} = \frac{j\omega C_v Y_A}{Y_A + j\omega C_v} + Y_B \tag{10}$$

The resonance condition that the imaginary part of the input admittance must be zero is shown by Equation (11). Thus Equations (8)–(11) can be used to calculate the resonance frequency of tunable SIR.

$$\text{Im}\{Y_{in}\} = 0 \tag{11}$$

To investigate the impact of locating position  $\theta_c$  on resonance characteristics of the proposed tunable SIR, the following analysis is conducted. Figure 5 shows resonance frequency versus the capacitance  $C_v$  of varactor diode with different values of locating position  $\theta_c$ . As the capacitance  $C_v$  increases, the resonance frequency of the proposed SIR decreases. However, for the same  $C_v$ , the resonance frequency will not change significantly with different values of loading position  $\theta_c$ .



**Figure 5.** Resonance frequency versus capacitance  $C_v$  with different values of  $\theta_c$ ,  $Y_1 = 0.03$  S,  $Y_2 = 0.01$  S,  $\theta_1 = \theta_2 = 90^\circ$ , where the reference frequency is 1GHz.

Tuning range (TR) is important for frequency tunable bandpass filters, which can be defined as Equation (12). In (12),  $f(C_{v\_min})$  and  $f(C_{v\_max})$  represent the resonance frequency of the SIR when the efficient capacitance  $C_v$  is the minimum and maximum value, respectively.

$$TR = \frac{f(C_{v\_min}) - f(C_{v\_max})}{(f(C_{v\_min}) + f(C_{v\_max}))/2} \tag{12}$$

Equation (13) defines the electrical length ratio  $\alpha$  of the proposed tunable SIR, and Equation (14) defines the characteristic admittance ratio  $\beta$  of the proposed tunable SIR. It can be seen that the resonance characteristic can be affected by  $\alpha$  and  $\beta$  from Equations (8)–(11).

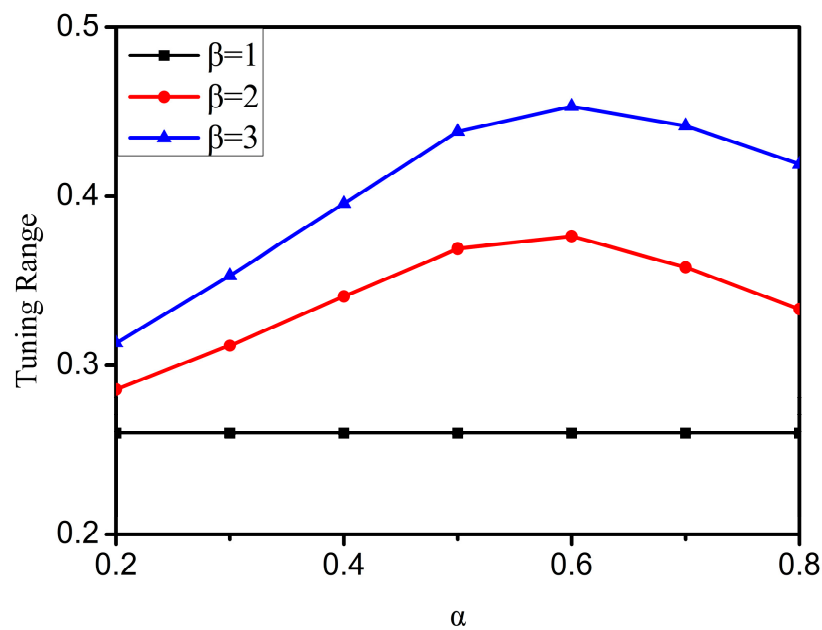
$$\alpha = \frac{\theta_2}{\theta_1 + \theta_2} \tag{13}$$

$$\beta = \frac{Y_1}{Y_2} \tag{14}$$

for specific electrical length ratio  $\alpha$  and characteristic admittance ratio  $\beta$ ,  $Y_2$ ,  $\theta_1$  and  $\theta_2$  as marked in Figure 4. These values can be calculated by Equations (13) and (14) when assuming that  $\theta_1 + \theta_2 = 180^\circ$  and  $Y_1 = 0.03$  S. Then, resonance frequency of the proposed step-impedance resonator can be calculated by Equations (8)–(11). Finally, the tuning range can be obtained by Equation (12), here,  $C_{v\_min} = 1.77$ pf  $C_{v\_max} = 9.24$  pf (SMV1413).

The relationships between TR and  $\alpha$  under different values of  $\beta$  is shown in Figure 6. If  $\beta$  is chosen to be 1, the proposed tunable resonator will become a UIR, and the TR will remain unchanged when  $\alpha$  varies. In addition, it can be seen that the TR will be widened when  $\beta$  increases, which means that the varactor-loaded SIR will have a wider TR compared with UIR. As  $\alpha$  increases, TR increases firstly and then decreases, and when  $\alpha = 0.6$ , the maximum of TR is realized.

Based on the above analysis, compared with conventional UIRs, the proposed tunable SIRs can achieve wider tuning range. Therefore, the proposed tunable SIR is adopted to implement the tunable BPF with high selectivity and constant ABW.



**Figure 6.** Tuning range of the tunable SIR versus electronic length ratio  $\alpha$  with different values of characteristic impedance  $\beta$ ,  $Y_1 = 0.03 \text{ S}$ ,  $\theta_c = 20^\circ$ ,  $\theta_1 + \theta_2 = 180^\circ$ , where the reference frequency is 1 GHz.

#### 4. Frequency-Dependent Coupling Structure

According to the conclusion of Section 2, the coupling coefficient  $M_{ij}$  must be inversely proportional to the frequency, and the external quality factor  $Q_e$  must be directly proportional to the frequency in order to realize constant ABW characteristic. Consequently, the design of frequency-dependent coupling structures is crucial to design a tunable filter with constant ABW. Some simulations are carried out to obtain the frequency-dependent coupling characteristic. The dielectric substrate Rogers 5880 with relative permittivity  $\epsilon_r = 2.2$  and thickness  $h = 0.787 \text{ mm}$  is used for simulation.

The coupling coefficient  $M_{ij}$  between varactor-loaded SIRs can be extracted by the equation in [23] as follows:

$$M_{ij} = \frac{f_o^2 - f_e^2}{f_o^2 + f_e^2} \quad (15)$$

where  $f_o$  and  $f_e$  represent the odd-mode and even-mode resonance frequency of the coupling structures, and  $f_o$  and  $f_e$  can be obtained by the commercial simulation software Advanced Design system (ADS) (v 2019).

Similarly, the external quality  $Q_e$  can be extracted by the following equation in [23] as follows:

$$Q_e = \frac{w_0 \tau_{S11}(w_0)}{4} \quad (16)$$

where  $w_0$  and  $\tau_{S11}$  represent the center frequency and group delay of S11, and  $w_0$  and  $\tau_{S11}$  can be obtained by ADS simulation.

##### 4.1. Design of Frequency-Dependent Electric Coupling Structure

The proposed frequency-dependent electric coupling structure between two varactor-loaded SIRs is shown in Figure 7. For a varactor-loaded SIR, the maximum position of the electric field is at the open end of the resonator. As the distance from the open end of the SIR increases, the electric field decreases, and the magnetic field increases; hence, the proposed coupling structure is electric coupling dominant.

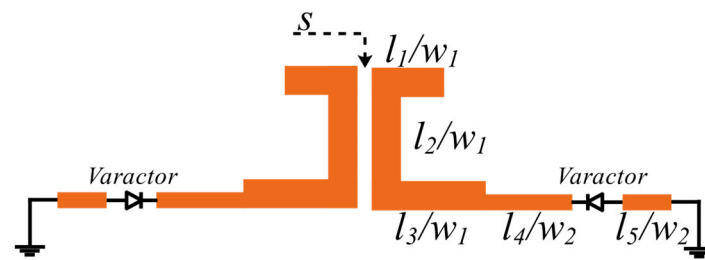


Figure 7. The proposed frequency-dependent electric coupling structure.

The extracted electric coupling coefficient versus efficient capacitance  $C_v$  with different values of coupling slot  $s$  and locating position  $l_5$  is shown in Figure 8. As shown, the coupling coefficient goes higher as  $C_v$  increases, while the resonant frequency decreases when the capacitance  $C_v$  increases. As a result, the coupling coefficient is inversely proportional to the frequency. For different values of coupling slot  $s$ , the coupling coefficient drops when the coupling slot  $s$  is larger. For different locating positions  $l_5$ , the slope of the coupling coefficient increases as the locating position  $l_5$  of the varactor increases. Thus, the magnitude and slope of coupling coefficient can be adjusted flexibly by altering the coupling slot  $s$  and locating position  $l_5$ , respectively.

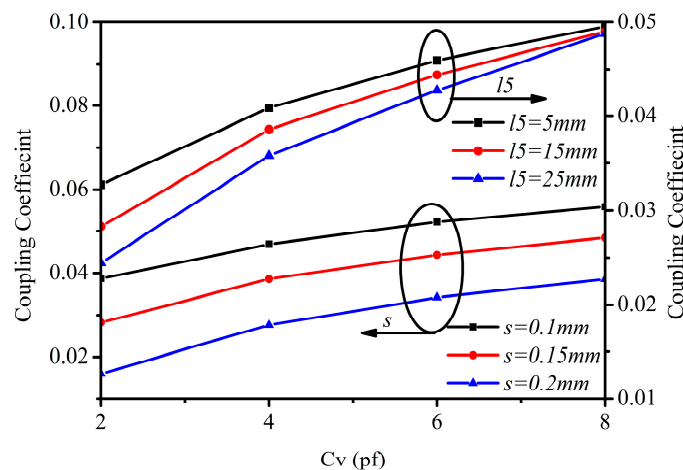


Figure 8. Electric coupling coefficient versus efficient capacitance  $C_v$  with different values of coupling slot  $s$  and locating position  $l_5$ ,  $w_1 = 2$  mm,  $w_2 = 1$  mm,  $l_1 = 5$  mm,  $l_2 = 25$  mm,  $l_3 = 20.5$  mm,  $l_4 + l_5 = 37.5$  mm.

#### 4.2. Design of Frequency-Dependent Magnetic Coupling Structure

Figure 9 shows the proposed frequency-dependent magnetic coupling structure between two varactor-loaded SIRs. As the distance from the short end of a varactor-loaded SIR rises, the magnetic field will decrease, and the electric field will increase. The maximum position of the magnetic field is at the short end of resonator; hence, the proposed coupling structure in this section is magnetic coupling dominant.

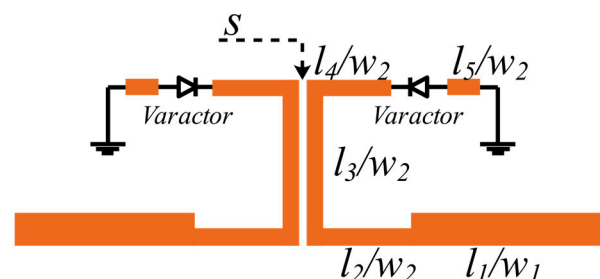
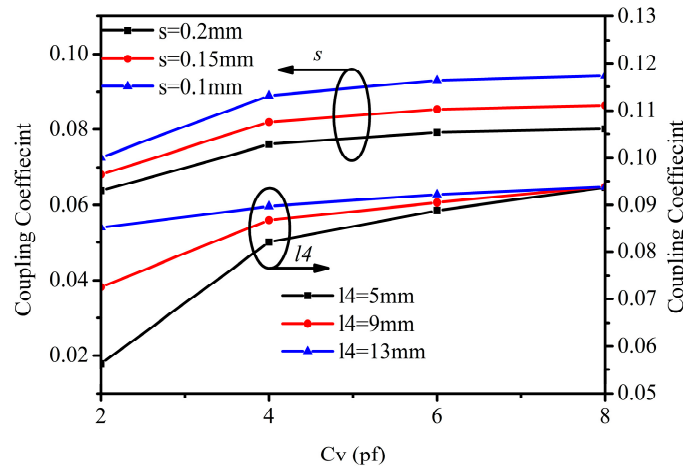


Figure 9. The proposed frequency-dependent magnetic coupling structure.



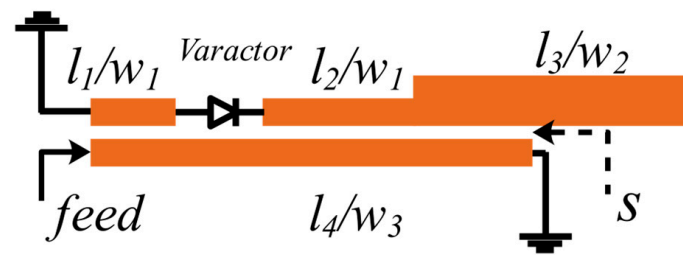
The extracted coupling coefficient versus capacitance  $C_v$  with different values of coupling slot  $s$  and  $l_4$  is shown in Figure 10. As shown, when the efficient capacitance  $C_v$  of varactor increases, the coupling increases, too. For different values of coupling slot  $s$ , coupling coefficient decreases when the coupling slot  $s$  increases. For different values of  $l_4$ , the slope of coupling coefficient will decrease when  $l_4$  increases; thus, the slope of magnetic coupling coefficient can be adjusted by changing  $l_4$ .



**Figure 10.** Magnetic coupling coefficient versus efficient capacitance  $C_v$  with different values of coupling slot  $s$  and  $l_4$ ,  $w_1 = 2$  mm,  $w_2 = 1$  mm,  $l_1 = 5$  mm,  $l_2 = 25$  mm,  $l_3 = 20.5$  mm,  $l_4 + l_5 = 37.5$  mm.

#### 4.3. Design of Frequency-Dependent Feeding Structure

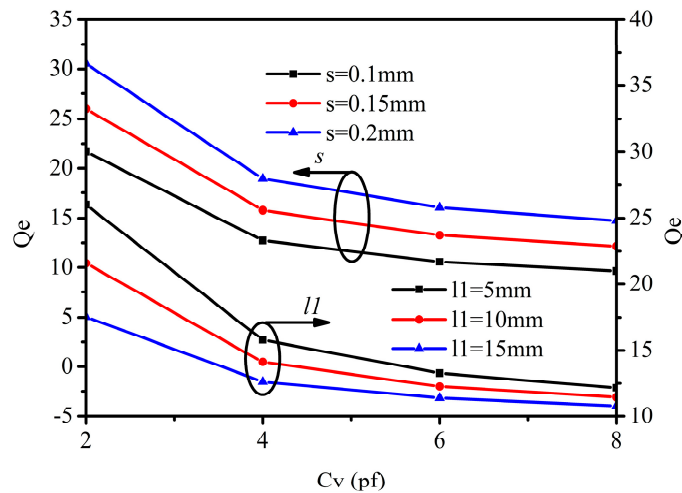
The proposed feeding structure for a varactor-loaded SIR is shown in Figure 11. The feeding microstrip line with short end is coupled to the varactor-loaded SIR. The width of the feeding microstrip line is  $w_3$ , and the coupling slot width between the resonator and feeding line is  $s$ .



**Figure 11.** The feeding structure of the varactor-loaded SIR.

Figure 12 presents the external quality  $Q_e$  versus efficient capacitance  $C_v$  of varactor with different values of coupling slot  $s$  and locating position  $l_1$ . As shown, the external quality  $Q_e$  goes lower as  $C_v$  increases, while the resonant frequency decreases when the capacitance  $C_v$  increases. As a result, the external quality  $Q_e$  is proportional to the frequency. For various values of coupling slot  $s$ , the larger coupling slot  $s$  means the larger external quality  $Q_e$ , and the slope of  $Q_e$  changes slightly when coupling slot  $s$  changes. However, for different values of locating position  $l_1$ , the slope of  $Q_e$  can be adjusted by changing the locating position  $l_1$  of varactor.

At this point, the frequency-dependent electric/magnetic coupling and feeding structures have been shown and analyzed. According the analysis before, the magnitude and slope of the coupling coefficient and external quality can be adjusted flexibly by changing the corresponding structural parameters, and this characteristic is crucial to realize constant ABW for tunable filters.

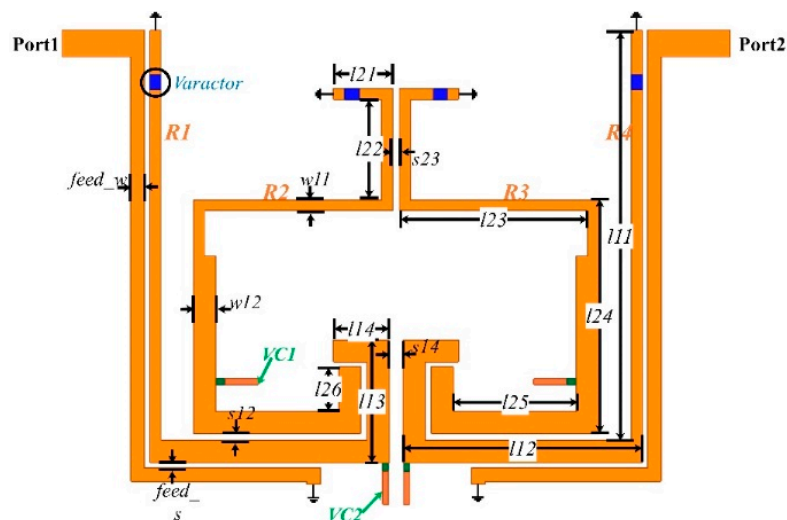


**Figure 12.** The external quality  $Q_e$  versus efficient capacitance  $c_v$  with different values of coupling slot  $s$  and locating position  $l_1$ .  $w_1 = 1$  mm,  $w_2 = 2$  mm,  $w_3 = 1.2$  mm,  $l_1 + l_2 = 37.5$  mm,  $l_3 = 37.5$  mm,  $l_4 = 53$  mm.

### 5. High-Selectivity Tunable BPF with Constant ABW

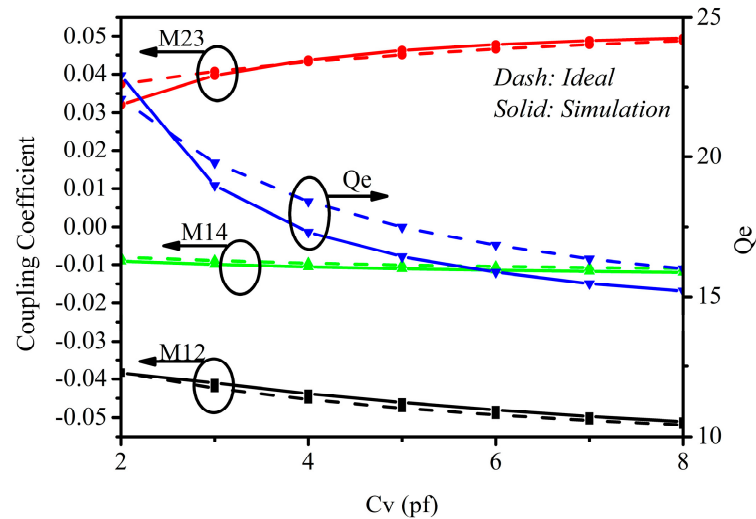
In order to validate the effectiveness of the frequency-dependent coupling structures, a fourth-order high-selectivity frequency tunable BPF with constant ABW based on varactor-loaded SIRs is designed in this section.

The layout of the proposed fourth-order tunable BPF is shown in Figure 13. The proposed tunable filter is comprised of four varactor-loaded SIRs, named R1, R2, R3 and R4, and two feeding microstrip lines, the resonators and feeding microstrip lines are connected to ground plane by metal vias with a diameter of 0.4 mm. The proposed tunable BPF is fabricated on Rogers 5880 substrate with a relative dielectric constant  $\epsilon_r$  of 2.2, height of 0.787 mm. The commercial varactor diode SMV1413 from Skyworks is chosen as the tuning element. The varactor diode SMV1413 can be equivalent to a variable capacitor  $C_v$ , inductor  $L_s$  and a resistor  $R_s$  in series. Variable capacitance  $C_v$  is from 1.77 pf (30V) to 9.24 pf (0V),  $L_s = 0.7$  nH, and  $R_s = 0.35$  Ohm. The values of capacitor, inductor and resistor are according to the datasheet of SMV1413. The effective capacitance of the varactors can be adjusted by DC voltage VC1 and VC2. The DC bias lines are connected to resonators by 100k ohm resistors to reduce the influence on resonators.



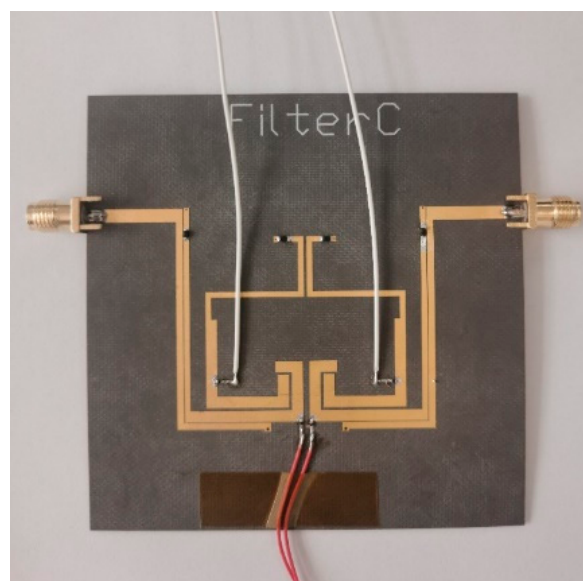
**Figure 13.** Layout of the proposed tunable filter with CABW.

The design targets are described in Section 2. Based on the synthesis procedure introduced in Section 2, the ideal coupling coefficient  $M_{12}$ ,  $M_{23}$ ,  $M_{14}$  and the external quality factor  $Q_e$  are shown in Figure 2. By adjusting the structure parameters of the frequency-dependent coupling structures described in Section 4, the simulation coupling coefficient  $M_{12}$ ,  $M_{23}$ ,  $M_{14}$  and the external quality factor  $Q_e$  can be fitted to the ideal values. The contrast between the final simulation values and the ideal values is illustrated in Figure 14.



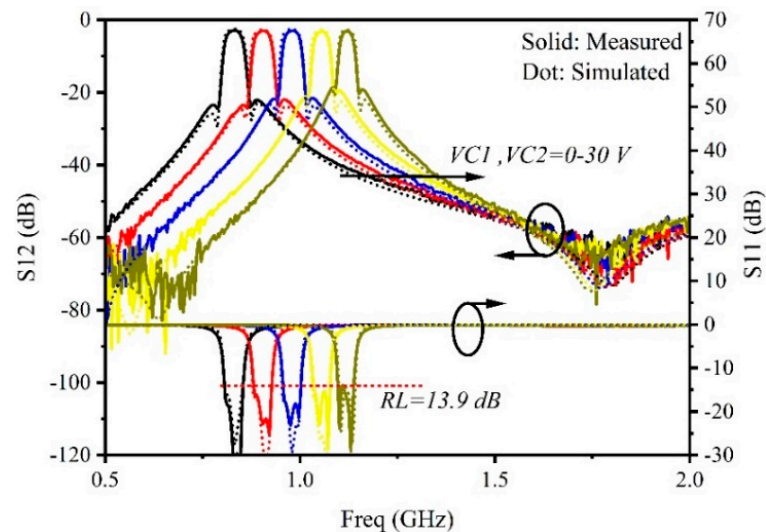
**Figure 14.** The ideal and simulated coupling coefficient and external quality  $Q_e$  for the predefined design goal.

The full-wave EM simulation software HFSS 15.0 (Driven Modal) is utilized to optimize the proposed filter. The final structure parameters are given as follows:  $w_{11} = 1$  mm,  $w_{12} = 2$  mm,  $feed_w = 1.2$  mm,  $l_{11} = 36.8$  mm,  $l_{12} = 17.9$  mm,  $l_{13} = 11$  mm,  $l_{14} = 5$  mm,  $l_{21} = 5.3$  mm,  $l_{22} = 9$  mm,  $l_{23} = 16.9$  mm,  $l_{24} = 21.1$  mm,  $l_{25} = 11$  mm,  $l_{26} = 4$  mm,  $s_{12} = 0.5$  mm,  $s_{23} = 0.62$  mm,  $s_{14} = 1.2$  mm,  $feed_s = 0.18$  mm. Finally, the overall size of the filter is  $39.3$  mm  $\times$   $45.2$  mm ( $0.14\lambda_g \times 0.17\lambda_g$  at  $0.8$  GHz). The photograph of the fabricated filter is shown in Figure 15.



**Figure 15.** The photograph of the fabricated filter.

The measured frequency responses are obtained by using Keysight N5244A vector network analyzer. The measured and simulated frequency response is shown in Figure 16. The result shows that the center frequency of the proposed filter can be tuned from 0.8 to 1.14 GHz while varying the bias voltage  $VC1/VC2$  from 0 to 30 V, the tuning range of the tunable BPF is about 35.1%. The return loss of the filter is better than 13.9 dB, and the insertion loss (IL) is 2.7–3.1 dB during the whole tuning range. The 3 dB ABW of the tunable filter is  $47 \pm 5$  MHz; obviously, the constant ABW characteristic is realized. Two transmission zeros on both sides of the passband are generated, and the selectivity of the proposed filter is improved, and the rectangular coefficient (20 dB:3 dB) is 2.3–3.1. Another two transmission zeros located far from the passband are produced to improve the stopband performance. It can be seen that there are some deviations of  $S_{11}$  and  $S_{12}$  between the measured and simulated EM results. The deviation may be caused by mechanical tolerances and discontinuity of connectors. In addition, the series resistance of the varactor diode and mechanical tolerances make measured insertion losses higher than the simulated ones.



**Figure 16.** Simulated and measured frequency response of the proposed tunable BPF.

The performance comparison with some prior works is listed in Table 1. It can be seen that many fourth-order tunable filters achieve two or three TZs due to the cross-coupling structures between resonators; however, the tuning range (normalized by center frequency, %) of the filters are not good. The filter in [7] based on SIR realizes a 50% tuning range, but the selectivity of the filter is poor, and ABW of the filter can remain constant. The proposed tunable filter based on varactor-loaded SIR has a wide tuning range and small layout size compared with other high-order tunable filters. Moreover, four transmission zeros make the proposed tunable filter own high frequency selectivity and wide stopband performance. Additionally, the constant ABW characteristic is realized, and the characteristic is very useful in some practical applications.

Although the proposed tunable filter realizes good performance in terms of tuning range, number of TZs, selectivity, constant ABW and layout size, the insertion loss of the filter is about 3.1 dB, and it is still a little large, which may limit its applications in some practical scenarios. The large insertion loss is mainly caused by the equivalent resistance of the varactor diodes. Thus, how to reduce the insertion loss of fourth-order tunable filters may be an important work in the future. In addition, high-order filters are more likely to realize better frequency selectivity and wide stopbands, so it is also an interesting work to apply varactor-loaded SIRs to the design of sixth- or eight-order frequency tunable filters.

**Table 1.** Typical performance comparison with prior works.

REF.	Freq (GHz)	TR (%)	BW (MHz)	IL (dB)	Order	NTZ <sup>3</sup>	Resonator	CABW <sup>2</sup>	Size ( $\lambda_g \times \lambda_g$ )
[4]	1.5–1.8	18.1	119	2.5	2	0	UIR	YES	0.18 × 0.13
[7]	0.6–1	50	93–155	2.8	2	0	SIR	NO	0.19 × 0.19
[13]	1.48–1.88	23	92	5.6	5	2	UIR	YES	NG
[15]	0.89–1.13	23.8	46.8	4.3	4	2	UIR	YES	0.23 × 0.19
[17]	1.26–1.61	24.4	176–225	2.3	4	2	UIR	NO	0.09 × 0.08
[18]	2.86–3.2	13	170–320	6.85	4	3	NG <sup>1</sup>	NO	0.3 × 0.17
[19]	2–2.45	20.2	104–135	6.85	4	2	LC	NO	NG <sup>1</sup>
[20]	1.36–1.78	26.6	93	5	4	2	UIR	YES	NG <sup>1</sup>
[22]	1.21–1.58	26.8	133	3	4	4	UIR	YES	0.73 × 0.11
[24]	2.24–2.64	16.3	240	3.9	4	1	SIW <sup>4</sup>	YES	NG <sup>1</sup>
Proposed	0.8–1.14	35.1	47	3.1	4	4	SIR	YES	0.14 × 0.17

<sup>1</sup> Not Given; <sup>2</sup> constant ABW; <sup>3</sup> numbers of TZ; <sup>4</sup> substrate-integrated waveguide.

## 6. Conclusions

In this paper, a compact high-selectivity fourth-order tunable bandpass filter based on varactor-loaded step-impedance resonators with constant ABW is proposed. By establishing cross coupling between resonators, a pair of transmission zeros located close to the passband are generated, and the selectivity of the filter is enhanced significantly. Another pair of transmission zeros are produced to improve the out-of-band rejection by using source-load coupling. Meanwhile, by adopting a frequency-dependent coupling structure between resonators, the ABW of the proposed tunable filter remains constant during the whole tuning range. In addition, the proposed tunable filter has a wide tuning range and compact size due to the use of the varactor-loaded SIR. The good performance of the proposed tunable filter makes it attractive in modern wireless communication systems.

**Author Contributions:** The paper is completed by S.L. (Shuang Li), S.L. (Shengxian Li) and J.Y. Conceptualization, S.L. (Shuang Li) and S.L. (Shengxian Li); methodology, S.L. (Shuang Li); validation, S.L. (Shuang Li), S.L. (Shengxian Li) and J.Y.; formal analysis, S.L. (Shuang Li); investigation, S.L. (Shuang Li); writing—original draft preparation, S.L. (Shuang Li); writing—review and editing, S.L. (Shengxian Li) and J.Y.; project administration, S.L. (Shengxian Li). All authors have read and agreed to the published version of the manuscript.

**Funding:** The APC was funded by Xi'an Institute of Space Radio Technology.

**Institutional Review Board Statement:** Not applicable.

**Informed Consent Statement:** Not applicable.

**Data Availability Statement:** The data presented in this study are available on request from the corresponding author. The data are not publicly available due to privacy.

**Conflicts of Interest:** The authors declare no conflict of interest.

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