



# *Article* **Geary's** *c* **and Spectral Graph Theory: A Complement**

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**Abstract:** Spatial autocorrelation, which describes the similarity between signals on adjacent vertices, is central to spatial science, and Geary's *c* is one of the most-prominent numerical measures of it. Using concepts from spectral graph theory, this paper documents new theoretical results on the measure. MATLAB/GNU Octave user-defined functions are also provided.

**Keywords:** spatial autocorrelation; Geary's *c*; spectral graph theory; graph Laplacian; graph Fourier transform

**MSC:** 62H11; 05C50

## **1. Introduction**

Using concepts from spectral graph theory, in this paper, we document new theoretical results on Geary's *c*, which is one of the most-prominent numerical measures of spatial autocorrelation. (Moran's *I* is another prominent numerical measure. For the measure, see, e.g., [\[1](#page-6-0)[,2\]](#page-6-1).) Here, spatial autocorrelation describes the similarity between signals on adjacent vertices and is central to spatial science ([\[3\]](#page-6-2)). Therefore, the results about the measure contribute to the development of spatial science.

More specifically, we provide a new representation of Geary's *c*. It is an expansion of it into a linear combination of variables with different degrees of spatial autocorrelation. By using the distribution of the coefficients, we can characterize spatial data. It is somewhat similar to Fourier series. Subsequently, we develop a way to compute the graph Laplacian eigenvectors needed for the graph Fourier transform. MATLAB/GNU Octave user-defined functions are also provided.

This paper can be considered complementary to [\[4\]](#page-6-3). As in this paper, using concepts from spectral graph theory, [\[4\]](#page-6-3) provided three types of representations for it: (a) graph Laplacian representation, (b) graph Fourier transform representation, and (c) Pearson's correlation coefficient representation. Our new representation can be regarded as an addition to them. Moreover, the way to compute the graph Laplacian eigenvectors is useful not only for this paper, but also for [\[4\]](#page-6-3).

We make two remarks on Geary's *c*. First, Geary's *c* was developed by [\[5\]](#page-6-4) and modified by  $[6-9]$  $[6-9]$ ) into what it is today. It is a spatial generalization of the von Neumann ratio ([\[10\]](#page-6-7)). Second, unlike Pearson's correlation coefficient, there exists the following:



See, e.g., [\[11\]](#page-6-8) (Equation (6)).

This paper is organized as follows. In Section [2,](#page-1-0) we provide some preliminaries for the following two sections. In Section [3,](#page-2-0) we present a new representation of Geary's *c*. In Section [4,](#page-3-0) we develop a way to compute the graph Laplacian eigenvectors needed for the



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graph Fourier transform. Section [5](#page-5-0) concludes. In Appendix [A,](#page-5-1) we provide MATLAB/GNU Octave user-defined functions.

#### *Some Notations*

Let  $y = [y_1, \ldots, y_n]^\top$ ,  $I_n$  be the identity matrix of order *n* and *ι* be the *n*-dimensional column vector of ones, i.e.,  $\boldsymbol{\iota} = [1, \dots, 1]^{\top}$ . For an  $n \times m$  full column rank matrix A, denote the column space of  $A$  and its orthogonal complement by  $\mathbb{S}(A)$  and  $\mathbb{S}^{\perp}(A)$ , respectively. For square matrices  $A_1, \ldots, A_p$ , diag $(A_1, \ldots, A_p)$  denotes the block diagonal matrix, whose diagonals are  $A_1, \ldots, A_p$ .

#### <span id="page-1-0"></span>**2. Preliminaries**

Let  $y_i$  denote the realization of a variable on a spatial unit *i* for  $i = 1, \ldots, n$ . Here, we exclude the case where  $y_1 = \cdots = y_n$ , i.e.,  $y \notin \mathbb{S}(\iota)$ . Accordingly,  $\sum_{i=1}^n (y_i - \bar{y})^2 > 0$ , where  $\bar{y}=\frac{1}{n}\sum_{i=1}^n y_i$ . In addition, let  $w_{ij}$  be the nonnegative spatial weight between the spatial units *i* and *j*. Here, we suppose that  $w_{ii} = 0$  and  $w_{ii} = w_{ii}$  for  $i, j = 1, \ldots, n$ . Accordingly, the spatial weights matrix  $W = [w_{ij}]$  is an  $n \times n$  symmetric hollow matrix. In addition, let  $\Omega = \sum_{i=1}^n \sum_{j=1}^n w_{ij}$ , which is assumed positive.

Geary's  $c$  for  $y$ , denoted by  $c(y)$ , is defined by

$$
c(\mathbf{y}) = \frac{n-1}{2\Omega} \frac{\sum_{i=1}^{n} \sum_{j=1}^{n} w_{ij} (y_i - y_j)^2}{\sum_{i=1}^{n} (y_i - \bar{y})^2}.
$$
 (1)

Let  $D = \text{diag}(d_1, \ldots, d_n)$ , where  $d_i = \sum_{j=1}^n w_{ij}$  for  $i = 1, \ldots, n$ . Then, the graph Laplacian in spectral graph theory (see, e.g.,  $[12-14]$  $[12-14]$ ) is defined as

<span id="page-1-1"></span>
$$
L = D - W.\t\t(2)
$$

Accordingly, as shown in, e.g., [\[4](#page-6-3)[,14\]](#page-6-10), *L* is a nonnegative definite matrix such that *Lι* = **0**. Ref. [\[4\]](#page-6-3) (Proposition 3.1) showed that *c*(*y*) can be represented using *L* as

$$
c(y) = \frac{n-1}{\Omega} \frac{y^{\top}Ly}{y^{\top}Q_iy'},
$$
\n(3)

where  $Q_i = I_n - \iota(\iota^\top \iota)^{-1} \iota^\top$ , which is a symmetric idempotent matrix, i.e.,  $Q_i^\top = Q_i$  and  $Q_{i}^{2} = Q_{i}$ .

Given that *L* is a real symmetric matrix, it can be spectrally decomposed as

<span id="page-1-2"></span>
$$
L = U \Lambda U^{\top}, \tag{4}
$$

where  $\Lambda = diag(\lambda_1, ..., \lambda_n)$  and  $\mathbf{U} = [\mathbf{u}_1, ..., \mathbf{u}_n]$  is an orthogonal matrix. Here,  $(\lambda_k, \mathbf{u}_k)$ denotes an eigenpair of *L* for  $k = 1, ..., n$ , and the eigenvalues,  $\lambda_1, ..., \lambda_n$ , are in ascending order. Given that *L* is a nonnegative definite matrix and  $L\ell = 0$ , we can suppose that  $(\lambda_1, u_1) = \left(0, \frac{1}{\sqrt{2}}\right)$  $\left(\frac{1}{n} \iota\right)$ . Let *m* denote the number of connected components. Then, it is known that  $0 = \lambda_1 = \cdots = \lambda_m < \lambda_{m+1} \leq \cdots \leq \lambda_n$ . See, e.g., [\[14\]](#page-6-10). Let  $\Lambda_2 = \text{diag}(\lambda_2, \ldots, \lambda_n)$  and  $U_2 = [u_2, \dots, u_n]$ . We show how to obtain  $U_2$ , as well as  $\Lambda_2$  from *L* in Section [4.](#page-3-0)

In spectral graph theory, the linear transformation given by  $U^{\dagger}y$  is referred to as the graph Fourier transform of *y* ([\[15\]](#page-6-11)). In addition,  $\lambda_k$  and  $u_k$  for  $k = 1, ..., n$  are referred to as graph Laplacian eigenvalues and graph Laplacian eigenvectors, respectively. Let  $[\alpha_1, \ldots, \alpha_n]^\top = \mathbf{U}^\top \mathbf{y}.$ 

Given that  $u_1 \in \mathbb{S}(t)$ ,  $u_k \in \mathbb{S}^{\perp}(t)$  for  $k = 2, \ldots, n$ , and  $Q_t$  is an orthogonal projection matrix onto  $\mathbb{S}^{\perp}(\iota)$ , it follows that  $Q_{\iota}\boldsymbol{U}=[\boldsymbol{0},\boldsymbol{u}_2,\dots,\boldsymbol{u}_n]$ , which yields

$$
\mathbf{y}^\top \mathbf{Q}_l \mathbf{U} = [0, \alpha_2, \dots, \alpha_n]. \tag{5}
$$

In addition, given that *L* is symmetric and *Lι* = **0**, it follows that

<span id="page-2-7"></span><span id="page-2-2"></span><span id="page-2-1"></span>
$$
L = Q_{\iota} L Q_{\iota}.
$$
 (6)

Moreover,  $U$  is an orthogonal matrix. By combining these results,  $c(y)$  can be represented as

$$
c(y) = \frac{n-1}{\Omega} \frac{y^{\top} Q_i U \Lambda U^{\top} Q_i y}{y^{\top} Q_i U U^{\top} Q_i y} = \frac{n-1}{\Omega} \frac{\sum_{k=2}^{n} \lambda_k \alpha_k^2}{\sum_{j=2}^{n} \alpha_j^2}.
$$
 (7)

Here, given that  $y \notin \mathbb{S}(\iota)$ , it follows that  $\sum_{j=2}^n \alpha_j^2 = y^\top Q_i y > 0$ . Finally, we note that [\(7\)](#page-2-1) is a part of [\[4\]](#page-6-3) (Proposition 3.3).

## <span id="page-2-0"></span>**3. A New Representation of Geary's** *c*

Given that  $u_k \notin \mathbb{S}(i)$  for  $k = 2, ..., n$ , we can consider

$$
c(\boldsymbol{u}_k) = \frac{n-1}{\Omega} \frac{\boldsymbol{u}_k^\top \boldsymbol{L} \boldsymbol{u}_k}{\boldsymbol{u}_k^\top \boldsymbol{Q}_i \boldsymbol{u}_k}, \quad k = 2, \dots, n,
$$
\n
$$
(8)
$$

which can be regarded as Geary's *c* when  $y = u_k$ . We note that, given  $u_1 \in \mathbb{S}(\iota)$ , Geary's *c* when  $y = u_1$  is excluded. (Actually, it cannot be defined. This is because  $u_1^\top Q_i u_1 = 0$ .)

For  $k = 2, ..., n$ , it follows that  $u_k^{\top} L u_k = u_k^{\top} U \Lambda U^{\top} u_k = e_k^{\top} \Lambda e_k = \lambda_k$  and  $u_k^{\top} Q_i u_k =$  $||u_k||^2 = 1$ , where  $e_k$  is the *k*-th column of  $I_n$ . Thus,  $c(u_k)$  in [\(8\)](#page-2-2) can be represented as

$$
c(\mathbf{u}_k) = \frac{n-1}{\Omega} \lambda_k, \quad k = 2, \dots, n. \tag{9}
$$

Then, from the inequalities,  $0 \leq \lambda_2 \leq \cdots \leq \lambda_n$ , it follows that

<span id="page-2-5"></span><span id="page-2-3"></span>
$$
0 \leq c(\mathbf{u}_2) \leq \cdots \leq c(\mathbf{u}_n). \tag{10}
$$

Moreover, given that  $\sum_{k=2}^{n} \lambda_k = \sum_{k=1}^{n} \lambda_k = \text{tr}(\bm{L}) = \Omega$ , it follows that

$$
\frac{1}{n-1}\sum_{k=2}^{n}c(u_k) = \frac{1}{n-1}\sum_{k=2}^{n}\frac{n-1}{\Omega}\lambda_k = 1.
$$
 (11)

By combining  $(7)$  and  $(9)$ , we obtain

$$
c(\mathbf{y}) = \frac{n-1}{\Omega} \frac{\sum_{k=2}^{n} \lambda_k \alpha_k^2}{\sum_{j=2}^{n} \alpha_j^2} = \frac{\sum_{k=2}^{n} \left(\frac{n-1}{\Omega} \lambda_k\right) \alpha_k^2}{\sum_{j=2}^{n} \alpha_j^2} = \frac{\sum_{k=2}^{n} c(u_k) \alpha_k^2}{\sum_{j=2}^{n} \alpha_j^2}
$$

$$
= \sum_{k=2}^{n} \left(\frac{\alpha_k^2}{\sum_{j=2}^{n} \alpha_k^2}\right) c(u_k) = \sum_{k=2}^{n} \psi_k c(u_k), \tag{12}
$$

where

<span id="page-2-4"></span>
$$
\psi_k = \frac{\alpha_k^2}{\sum_{j=2}^n \alpha_j^2}, \quad k = 2, \dots, n. \tag{13}
$$

Note that  $\psi_k$  in [\(13\)](#page-2-4) satisfy  $\psi_k \geq 0$  and  $\sum_{k=2}^n \psi_k = 1$ .

The next proposition summarizes the above-mentioned results.

<span id="page-2-6"></span>**Proposition 1.** (a)  $c(y)$  *in* [\(1\)](#page-1-1) *can be represented as*  $\sum_{k=2}^{n} \psi_k c(u_k)$ )*, where*  $\psi_k = \frac{\alpha_k^2}{\sum_{j=2}^n \alpha_j^2}$  for  $k = 2, \ldots, n$ . Here,  $\psi_k$  for  $k = 2, \ldots, n$  are nonnegative and sum to unity. *(b)*  $c(\textit{\textbf{u}}_k)$  for  $k=2,\ldots$  ,  $n$  satisfy the inequalities given by [\(10\)](#page-2-5), and their simple average equals unity. **Remark 1.** *Concerning Proposition [1,](#page-2-6) we make three remarks:*

- *(i)* Proposition [1\(](#page-2-6)*a*) implies that Geary's *c* can be represented as a weighted average of  $c(u_2), \ldots, c(u_n)$ . Concerning the weight,  $\psi_k$ , the larger  $|\alpha_k| = |u_k^{\top}y| = |(u_k^{\top}u_k)^{-1}u_k^{\top}y|$  is, the larger  $\psi_k$  is. We note that  $\alpha_k = arg\min_{\phi_k} \|y - \phi_1u_1 - \cdots - \phi_nu_n\|^2 = arg\min_{\phi_k} \|y - \phi_ku_k\|^2$ .
- *(ii)* Proposition [1\(](#page-2-6)b) implies that the graph Laplacian eigenvectors,  $u_2, \ldots, u_n$ , can be sorted in the *spatial autocorrelation measured by Geary's c. In other words, u<sup>k</sup> is more positively spatially* autocorrelated than or equal to  $u_{k+1}$  for  $k=2,\ldots$  ,  $n-1.$  Accordingly,  $\bm{y}$  is positively (respec*tively negatively) spatially autocorrelated if* {*ψk*} *is a monotonically decreasing (respectively increasing) sequence. y can be characterized by the distribution of the coefficients,*  $\psi_2, \dots, \psi_n$ *. It is somewhat similar to the Fourier series.*

(iii) If 
$$
c(u_2) = \cdots = c(u_n) = \mu
$$
, then  $c(y) = \sum_{k=2}^n \psi_k c(u_k) = \mu \sum_{k=2}^n \psi_k = \mu$  regardless of **y**.

Let  $\eta = \eta_1 + \eta_2$ , where  $\eta_1 \in \mathbb{S}(\iota)$  and  $\eta_2 \in \mathbb{S}(\iota) \setminus \{0\}$ . Then,  $c(\eta)$  equals  $c(\iota)$ . That is, for all  $\gamma_1 \in \mathbb{R}$  and  $\gamma_2 \in \mathbb{R} \setminus \{0\}$ , it follows that

<span id="page-3-1"></span>
$$
c(\gamma_1 \iota + \gamma_2 y) = c(y). \tag{14}
$$

Given that  $Q_i y = y - \bar{y} \iota$ ,  $c(Q_i y) = c(y)$  is an example of [\(14\)](#page-3-1). From (14) and Proposi-tion [1\(](#page-2-6)a), we obtain  $c(\gamma_1 \mathbf{i} + \gamma_2 \mathbf{y}) = \sum_{k=2}^n \psi_k c(\mathbf{u}_k)$ .

The next corollary summarizes the above result.

<span id="page-3-3"></span>**Corollary 1.** *For all*  $\gamma_1 \in \mathbb{R}$  *and*  $\gamma_2 \in \mathbb{R} \setminus \{0\}$ *, c*( $\gamma_1 \iota + \gamma_2 \iota$ *y*) *equals*  $\sum_{k=2}^n \psi_k c(\iota_k)$ *.* 

In addition, from Proposition [1,](#page-2-6) it immediately follows that

$$
c(y) = \sum_{k=2}^{n} \psi_k c(u_k) \leq \sum_{k=2}^{n} \psi_k c(u_n) = c(u_n) \sum_{k=2}^{n} \psi_k = c(u_n).
$$
 (15)

Likewise,  $c(u_2) \leq c(y)$  follows.

The next corollary summarizes the above result.

<span id="page-3-2"></span>**Corollary 2.**  $c(y)$  belongs to the closed interval given by  $[c(u_2), c(u_n)]$ .

**Remark 2.** *Concerning Corollary [2,](#page-3-2) we make two remarks:*

- *(i)* If  $c(u_2) = \cdots = c(u_n)$ , then the interval given by  $[c(u_2), c(u_n)]$  reduces to a singleton. For *example, if*  $W = \boldsymbol{\mu}^\top - \mathbf{I}_n$ , which is the binary adjacency matrix of the complete graph with n *vertices, then*  $L = (n-1)I_n - (u^\top - I_n) = nQ_t$  *and, accordingly, c*( $u_2$ ) = · · · = *c*( $u_n$ ) =  $\frac{n-1}{n(n-1)}\times n=1.$  Then, in this case,  $c(\bm{y})=\sum_{k=2}^n\psi_k c(\bm{u}_k)=\sum_{k=2}^n\psi_k=1$  regardless of  $\bm{y}.$
- *(ii)*  $Ref.$  [\[11\]](#page-6-8) showed that  $c(y)$  belongs to the closed interval given by  $\left\lceil \frac{n-1}{\Omega} \lambda_2, \frac{n-1}{\Omega} \lambda_n \right\rceil$ . Given [\(9\)](#page-2-3), *Corollary [2](#page-3-2) is its equivalent.*

## <span id="page-3-0"></span>**4. A Way to Compute the Eigenvectors in** [\(4\)](#page-1-2)

In this section, we develop a way to compute  $U_2 = [u_2, \ldots, u_n]$ , which also provides  $\Lambda_2 = \text{diag}(\lambda_2, ..., \lambda_n)$ . Here, we explain the reason why it is useful. When there is only one connected component, i.e.,  $m = 1$ , there is no problem. This is because *L* has single 0 eigenvalue and the corresponding normalized eigenvector is  $\frac{1}{\sqrt{2}}$  $\frac{1}{n}$ *ι*. However, when there is more than one connected component, i.e.,  $m \geq 2$ , L has multiple 0 eigenvalues, and then,

 $\frac{1}{\sqrt{2}}$  $\frac{1}{n}$ *ι* is not necessarily one of the eigenvectors returned from a computer program. For example, when

<span id="page-4-1"></span>
$$
W = \begin{bmatrix} 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 \end{bmatrix},
$$
(16)

there are two connected components and, accordingly,  $m = 2$  (Figure [1\)](#page-4-0). In this case, both  $(0, u_1^*)$  and  $(0, u_2^*)$ , where  $u_1^* = \frac{1}{\sqrt{2}}$  $\frac{1}{3}$ [1, 1, 1, 0, 0]<sup>⊤</sup> and  $u_2^* = \frac{1}{\sqrt{3}}$  $\mathbb{E}_{\overline{2}}[0, 0, 0, 1, 1]^\top$ , and  $(0, u_1)$ and  $(0, u_2)$ , where  $u_1 = \frac{1}{\sqrt{2}}$  $\frac{1}{\sqrt{5}}[1, 1, 1, 1, 1]^{\top}$  and  $u_2 = \left[\sqrt{\frac{2}{15}}, \sqrt{\frac{2}{15}}, \sqrt{\frac{2}{15}}, -\sqrt{\frac{3}{10}}, -\sqrt{\frac{3}{10}}\right]^{\top}$ , are eigenpairs of the corresponding graph Laplacian. The results below can handle such a situation.

<span id="page-4-0"></span>
$$
\begin{array}{c}\n\begin{pmatrix}\n1 \\
\end{pmatrix}\n\end{array}\n\qquad\n\begin{array}{c}\n\begin{pmatrix}\n3 \\
\end{pmatrix}\n\end{array}\n\qquad\n\begin{array}{c}\n\begin{pmatrix}\n4 \\
\end{pmatrix}\n\end{array}\n\qquad\n\begin{array}{c}\n\begin{pmatrix}\n5 \\
\end{pmatrix}\n\end{array}
$$

**Figure 1.** Undirected graph whose binary adjacency matrix is *W* in [\(16\)](#page-4-1).

Let  $\{g_2, \ldots, g_n\}$  denote any orthonormal basis of  $\mathbb{S}^{\perp}(\iota)$  and  $G = [g_1, G_2]$ , where  $g_1 = \frac{1}{\sqrt{2}}$  $\mathbf{F}_n$ *ι* and  $G_2 = [g_2, \ldots, g_n]$ . Then, *G* is an  $n \times n$  orthogonal matrix. Accordingly, from  $(4)$  and  $(6)$ , it follows that

$$
G^{\top} Q_{\iota} L Q_{\iota} G V = V \Lambda, \qquad (17)
$$

where  $V = G^{\top}U$ . Here, from  $g_1^{\top}u_1 = 1$ ,  $g_1^{\top}U_2 = 0$ , and  $G_2^{\top}u_1 = 0$ , it follows that

<span id="page-4-4"></span><span id="page-4-3"></span><span id="page-4-2"></span>
$$
V = \text{diag}(1, V_2),\tag{18}
$$

where  $V_2 = G_2^{\top} U_2$ , which is an  $(n-1) \times (n-1)$  orthogonal matrix. This is because  $V^{\top}V = U^{\top}GG^{\top}U = I_n$  and  $V^{\top}V = \text{diag}(1, V_2^{\top}V_2)$ . In addition, given that  $Q_\iota G = [\mathbf{0}, G_2]$ , it follows that

$$
G^{\top} Q_{\iota} L Q_{\iota} G = \text{diag}(0, G_2^{\top} L G_2). \tag{19}
$$

By combining [\(18\)](#page-4-2) and [\(19\)](#page-4-3), [\(17\)](#page-4-4) becomes diag( $0, G_2^{\top}LG_2$ )  $\cdot$  diag( $1, V_2$ )  $=$  diag( $1, V_2$ )  $\cdot$ diag( $\lambda_1$ ,  $\Lambda_2$ ). Therefore, it follows that

$$
G_2^{\top} L G_2 V_2 = V_2 \Lambda_2. \tag{20}
$$

Here, recall that  $V_2$  is an orthogonal matrix. In addition, given that  $G$  is an orthogonal matrix, premultiplying [\(18\)](#page-4-2) by  $G = [g_1, G_2]$  yields  $U = [g_1, G_2] \cdot diag(0, V_2) = [g_1, G_2 V_2]$ . Therefore, it follows that

<span id="page-4-5"></span>
$$
U_2 = G_2 V_2. \tag{21}
$$

The next proposition summarizes the above-mentioned results.

<span id="page-4-6"></span>**Proposition 2.** Denote the k-th column of  $V_2$  by  $v_{k+1}$  for  $k = 1, \ldots, n-1$ , i.e.,  $V_2 = [v_2, \ldots, v_n].$ *Then,*  $(\lambda_k, v_k)$  for  $k = 2, ..., n$  are the eigenpairs of  $G_2^{\top}LG_2$ . In addition,  $U_2$  is obtainable from  $V_2$ *by* [\(21\)](#page-4-5)*.*

**Remark 3.** *Concerning Proposition [2,](#page-4-6) we make two remarks:*

*(i) The following n*  $\times$  *(n* − 1*) matrix*  $\mathbf{F}_2$  *is an example of*  $\mathbf{G}_2$ *:* 

$$
F_2 = \begin{bmatrix} 1 & \cdots & \cdots & 1 \\ -1 & \ddots & & \vdots \\ 0 & -2 & \ddots & \vdots \\ \vdots & \ddots & \ddots & 1 \\ 0 & \cdots & 0 & -(n-1) \end{bmatrix} r^{-1},
$$
(22)

*where*  $\Gamma = \text{diag}($ √  $\overline{1\cdot 2}, \ldots, \sqrt{(n-1)\cdot n}$ . (The use of  $\overline{F_2}$  is inspired by [\[16\]](#page-6-12).) Here,  $F = [f_1, F_2]$ , where  $f_1 = \frac{1}{\sqrt{2}}$ *n ι, is a Helmert orthogonal matrix ([\[17\]](#page-6-13)). Instead of F*2*, we may use*  $H_2 = \Delta^\top (\Delta \Delta^\top)^{-\frac{1}{2}}$ , where  $\Delta$  is the  $(n-1)\times n$  matrix such that  $\Delta\zeta = [\zeta_2 - \zeta_1]$  $\zeta_1,\ldots,\zeta_n-\zeta_{n-1}]^{\top}$  for an *n*-dimensional vector  $\zeta=[\zeta_1,\ldots,\zeta_n]^{\top}$ . This is because  $H_2$  $satisfies$  that  $H_2^{\top} \iota = (\Delta \Delta^{\top})^{-\frac{1}{2}} \Delta \iota = 0$  and  $H_2^{\top} H_2 = (\Delta \Delta^{\top})^{-\frac{1}{2}} \Delta \Delta^{\top} (\Delta \Delta^{\top})^{-\frac{1}{2}} = I_{n-1}$ . *Here, ∆∆*<sup>&</sup>gt; *is a positive definite matrix.*

*(ii) MATLAB/GNU Octave user-defined functions required for the calculation of* **Λ**2*, U*2*, and the bounds of Geary's c are provided in Appendix [A.](#page-5-1)*

### <span id="page-5-0"></span>**5. Concluding Remarks**

In this paper, we showed new theoretical results on Geary's *c*, which included (i) a new representation of Geary's *c* and (ii) a way to compute the graph Laplacian eigenvectors. The obtained results are summarized in Propositions [1](#page-2-6) and [2](#page-4-6) and Corollaries [1](#page-3-3) and [2.](#page-3-2) The required MATLAB/GNU Octave user-defined functions are also provided. Finally, as stated, this paper can be considered complementary to [\[4\]](#page-6-3).

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**Conflicts of Interest:** The authors declare no conflict of interest.

### <span id="page-5-1"></span>**Appendix A. MATLAB/GNU Octave User-Defined Functions**

In this section, we provide three MATLAB/GNU Octave user-defined functions. Among them, Lam2U2 is a function for calculating  $U_2$ , as well as  $\Lambda_2$  from L. Gearycbounds is a function for calculating the bounds of Geary's *c* corresponding to *L*. Note that (A+A')/2 in the functions is to ensure symmetry. Finally, these two functions depend on Fmat, which is a function to make *F*.

```
1
2 function [Lam2, U2] = Lam2U2(W)\text{ }3 \text{ } n=size (W, 1);
_4 L=diag (sum (W, 2)) -W;
5 F = Fmat(n); F2 = F(:,2:n); A = F2' * L * F2;6 [X, E] = eig ((A+A')/2);
      [e, ind] = sort ( diag ( E ), 'ascend');
8 Lam2=diag(e); V2 = X(:,ind);9 \text{ } U2 = F2 * V2;10 end
1 function [c_lb, c_ub] = Gearycbounds (W)
n = size(W, 1);3 Omega=sum (sum (W));
4 L = diag(sum(W, 2)) - W;
```

```
5 F = Fmat(n); F2 = F(:,2:n); A = F2' * L * F2;6 eigv=sort \left(eig((A+A')/2), 'ascend');
7 \t c_1 b = ((n - 1) / 0mega) * eigv(1);
s \t c_ub = ((n-1) / 0 \text{mega}) * eigv(n-1) ;9 end
1 function [F] = Fmat(n)P = zeros(n, n):
F(:, 1) = ones(n, 1) / sqrt(n);4 for k=2:nF(:, k) =[ones (k-1, 1); (k-1); zeros (n-k, 1)]/sqrt ((k-1)*k);6 end
  end
```
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