

Article

# Grating Coupler Design for Low-Cost Fabrication in Amorphous Silicon Photonic Integrated Circuits

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**Abstract:** Photonic circuits find applications in biomedicine, manufacturing, quantum computing and communications. Photonic waveguides are crucial components, typically having cross-section orders of magnitude inferior when compared with other photonic components (e.g., optical fibers, light sources and photodetectors). Several light-coupling methods exist, consisting of either on-plane (e.g., adiabatic and end-fire coupling) or off-plane methods (e.g., grating and vertical couplers). The grating coupler is a versatile light-transference technique which can be tested at wafer level, not requiring specific fiber terminations or additional optical components, like lenses, polarizers or prisms. This study focuses on fully-etched grating couplers without a bottom reflector, made from hydrogenated amorphous silicon (a-Si:H), deposited over a silica substrate. Different coupler designs were tested, and of these we highlight two: the superimposition of two lithographic masks with different periods and an offset between them to create a random distribution and a technique based on the quadratic refractive-index variation along the device's length. Results were obtained by 2D-FDTD simulation. The designed grating couplers achieve coupling efficiencies for the TE-like mode over  $-8$  dB (mask overlap) and  $-3$  dB (quadratic variation), at a wavelength of 1550 nm. The coupling scheme considers a 220 nm a-Si:H waveguide and an SMF-28 optical fiber.

**Keywords:** grating coupler; light coupling; photonic waveguide; silicon-on-insulator; amorphous silicon; off-plane coupling



**Citation:** Almeida, D.; Lourenço, P.; Fantoni, A.; Costa, J.; Vieira, M. Grating Coupler Design for Low-Cost Fabrication in Amorphous Silicon Photonic Integrated Circuits. *Photonics* **2024**, *11*, 783. <https://doi.org/10.3390/photonics11090783>

Received: 6 July 2024

Revised: 6 August 2024

Accepted: 19 August 2024

Published: 23 August 2024



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## 1. Introduction

Photonic devices and circuits find applications in various fields of science, including communications [1,2], quantum computing [3–5], neural networks [6–9] and sensing, in biosensors [10–14], temperature sensors [15] and gas detectors [16–18]. Amorphous silicon in its hydrogenated form (a-Si:H) and hydrogenated amorphous silicon nitride (a-SiN:H) are suitable materials for photonic components; the former is usually preferred for applications in the near-infrared (NIR) wavelength [19] and the latter is used for near-infrared [20–22] and visible-light [22–25] wavelengths. Hydrogenated amorphous silicon compounds have the advantage of being compatible with CMOS manufacturing technology [23] and with low-temperature deposition methods, such as Plasma-Enhanced Chemical Vapor Deposition (PECVD) [26,27] and Hot-Wire Chemical Vapor Deposition (HWCVD) [28], contributing to a significant reduction in manufacturing costs. The relatively low extinction coefficient of hydrogenated amorphous silicon in the optical C-band (1530–1565) makes it a suitable material for near-infrared photonics. Waveguides, having a-Si:H as core material, can achieve remarkably low propagation losses, of 1 dB/cm [29] or less [30,31].

There are several methods to couple light from optical fibers to photonic waveguides and vice-versa, and the more common are based on [32,33] end fire-coupling, grating couplers or adiabatic coupling. We have included a brief review of light-coupling methods,

highlighting the advantages and disadvantages of several approaches commonly found in the literature. Prism light coupling is a possible alternative which is usually employed with waveguides made of materials featuring lower refractive indexes (under 2.5 at the wavelength of interest), such as silicon nitride and polymer. Silicon, both in crystalline and amorphous form, has a high refractive index (approximately 3.5 for a wavelength of 1550 nm). Prisms suitable for light coupling must have a higher refractive index than the refractive index of the waveguide core material and a low extinction coefficient at the wavelength of interest. There are no commercially available optical prisms with refractive indexes higher than that of hydrogenated amorphous silicon (a-Si:H) and with a low extinction coefficient at near-infrared (NIR) wavelengths. Grating couplers are attractive due to their off-plane coupling capability, which increases integration with other photonic devices, such as photodetectors and optical fibers. When used in conjunction with high-efficiency compact broadband tapers [34], coupling schemes featuring grating couplers can theoretically achieve efficiencies over  $-3$  dB. Light couplers based on diffraction gratings offer some advantages over other coupling methods, such as waveguide tapers and adiabatic couplers, in the sense that the former do not require application-specific fiber terminations [35] or anisotropic etching [36]. The avoidance of complex and tailor-made processes increases repeatability and lowers fabrication complexity.

In this study, we compare different a-Si:H grating coupler (GC) designs, and for this purpose bi-dimensional finite-difference time-domain (2D-FDTD) analyses were performed. Simulations were performed with the RSoft Photonics CAD Suite software package version 2023.12-SP1 (Synopsys, Inc., Sunnyvale, CA, USA), via University Donation Program [37]. Grating couplers featuring apodization (variation of the period and fill factor along the coupler's length) are compared with GCs without apodization. Several approaches are taken to optimize coupling efficiency, such as quadratic and linear variation of the effective refractive index over the GCs' length on apodized designs, and lithography-mask superimposition with offset and fill factor greater than 50% on non-apodized GCs. Due to the fact that the diffracted field decays in an exponential-like fashion in non-apodized gratings (over the length), these have lower efficiencies (under  $-4$  dB); in comparison, the diffracted field of apodized grating couplers has a distribution closer to a Gaussian, resulting in higher efficiencies (over  $-3.2$  dB), even for designs (such as those presented in this work) which do not feature any kind of buried layer acting as a reflector.

Hydrogenated amorphous silicon (a-Si:H), the material employed in the proposed devices, can be deposited by Plasma-Enhanced Chemical Vapor Deposition (PECVD) at low temperatures ( $\leq 250$  °C), contributing to a production-cost reduction. The a-Si:H layer can be deposited over a substrate of silica, or alternatively, above a glass surface, not requiring a silicon wafer. The grating couplers presented in this study can be fabricated with one or two lithographic masks, lowering the probability of imperfections caused by misalignment. Photolithography can be used to pattern the designs, employing either deep- or extreme-ultraviolet (UV) technology, thus avoiding the expensive and time-consuming electron beam (E-beam) process.

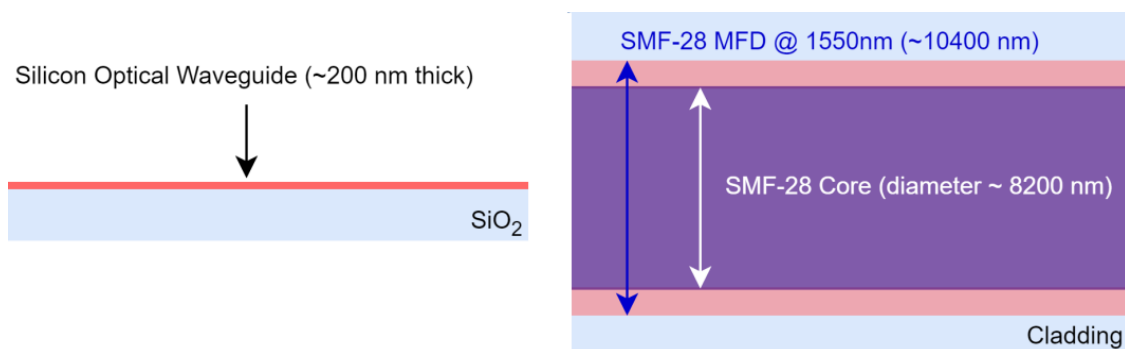
The main purpose of this study is to prove that it is possible to conceive alternative solutions to grating couplers featuring bottom reflector layers and to demonstrate that low-cost production techniques (i.e., based on low-temperature deposition, without the need for a silicon wafer and compatible with CMOS technology) can be employed to fabricate devices with a subjectively good performance. We were able to achieve a coupling efficiency of  $-4.3$  dB for a non-apodized design and  $-3.1$  dB and  $-2.8$  dB for designs featuring linear and quadratic apodization, respectively. The simplicity of the designs featured in this study also makes the devices less demanding in terms of fabrication complexity and less prone to misalignment issues.

This article is an extended version of the paper published in the *Proceedings of the SPIE 12880, Physics and Simulation of Optoelectronic Devices XXXII* (11 March 2024) [38]. This manuscript is organized into seven main chapters. The first chapter, Introduction, gives the reader an overview of the presented work, about the main research challenges

and advantages of the proposed grating-coupler designs. In the second chapter, a Brief Review of Light-Coupling Methods, the purpose of light coupling between photonic devices and optical fibers is briefly explained, some of the most common and best-performing light-coupler methods found in the literature are presented, and its main advantages and disadvantages are explained. A subchapter devoted to more advanced grating-coupler designs is also included, to highlight some of the latest progress in this type of technology and to emphasize the challenges posed by such techniques. The third chapter, Materials and Methods, is devoted to the explanation of the theoretical background and techniques proposed in our work; the type of optical fiber, waveguide and grating-coupler design methods are explained in this section. The fourth chapter, Results, is devoted to the presentation of 2D-FDTD and FEM simulation results of the proposed grating-coupler designs. The fifth chapter, Discussion, compares the proposed grating couplers with some of the most advanced designs, highlighting the advantages and drawbacks of the different methods. In the sixth chapter, Conclusions, the main findings of this study are presented and discussed. Finally, we have the seventh and last chapter, devoted to Future Work and Improvement, and in this section we highlight future perspectives in grating-coupler designs and possible methods for further improvement.

## 2. Brief Review of Light-Coupling Methods

Size mismatch between the optical fiber (e.g., SMF-28) and photonic waveguide is significant, Figure 1. In our specific setup, the former have a core diameter of 8.2  $\mu\text{m}$  and the latter a core height (thickness) of 0.22  $\mu\text{m}$ . The optical-fiber mode field diameter is approximately 10.4  $\mu\text{m}$ .



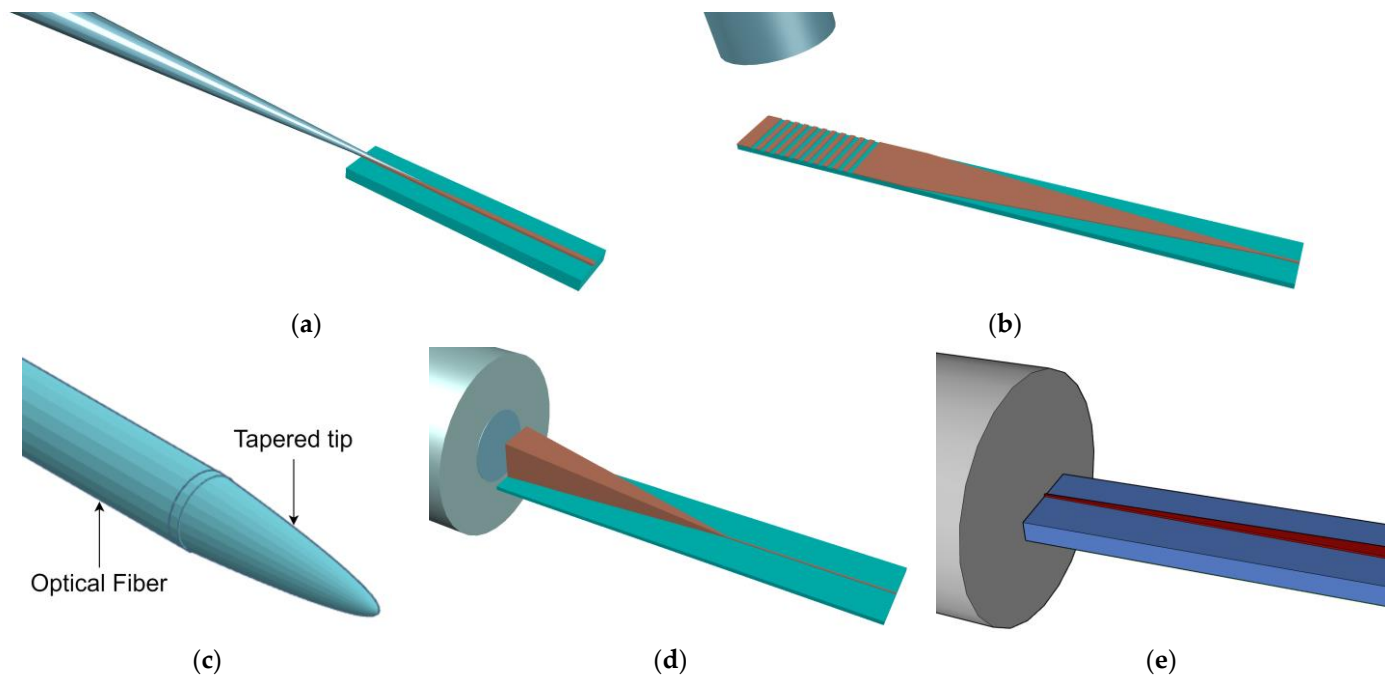
**Figure 1.** Size mismatch between a silicon optical waveguide ( $\approx 200$  nm thick) and an SMF-28 optical fiber, with core width  $\approx 8200$  nm and a mode field diameter (MFD) of about 10,400 nm. The system operates at a wavelength of 1550 nm. This illustration represents end-fire coupling without the assistance of a 3D taper.

Several coupling schemes can be employed to overcome this issue. We highlight the following [33]:

- Adiabatic couplers [33,35];
- Grating couplers [39–42];
- Lensed fibers [43];
- Spot size converters [33,44–47].

The first and third methods, as shown in Figure 2a,c, require tailor-made optical fiber terminations which demand fiber manipulation techniques, needing expensive equipment that is not commonly found in most optical laboratory setups. The first, fourth and fifth schemes, Figure 2a,d,e typically require a very good matching between the optical fiber termination and the waveguide's coupling region(s). These designs require very high-precision lithographic equipment (e.g., based on Electron Beam), which is very expensive and not compatible with CMOS technology. Some spot size converters also rely on vertical structures which can be very complex and expensive to manufacture, like, for example,

3D tapers [45–47] (Figure 2d) and vertical couplers [43]. Unlike its crystalline counterparts, amorphous silicon cannot be subjected to anisotropic wet etching [36], making the development of 3D structures more complex. Some grating couplers can have minimum feature sizes of the orders of dozens of nanometers or even a few hundred nanometers, making these devices feasible using photolithographic processes such as deep UV (DUV) lithography [39,42], employing excimer lasers.



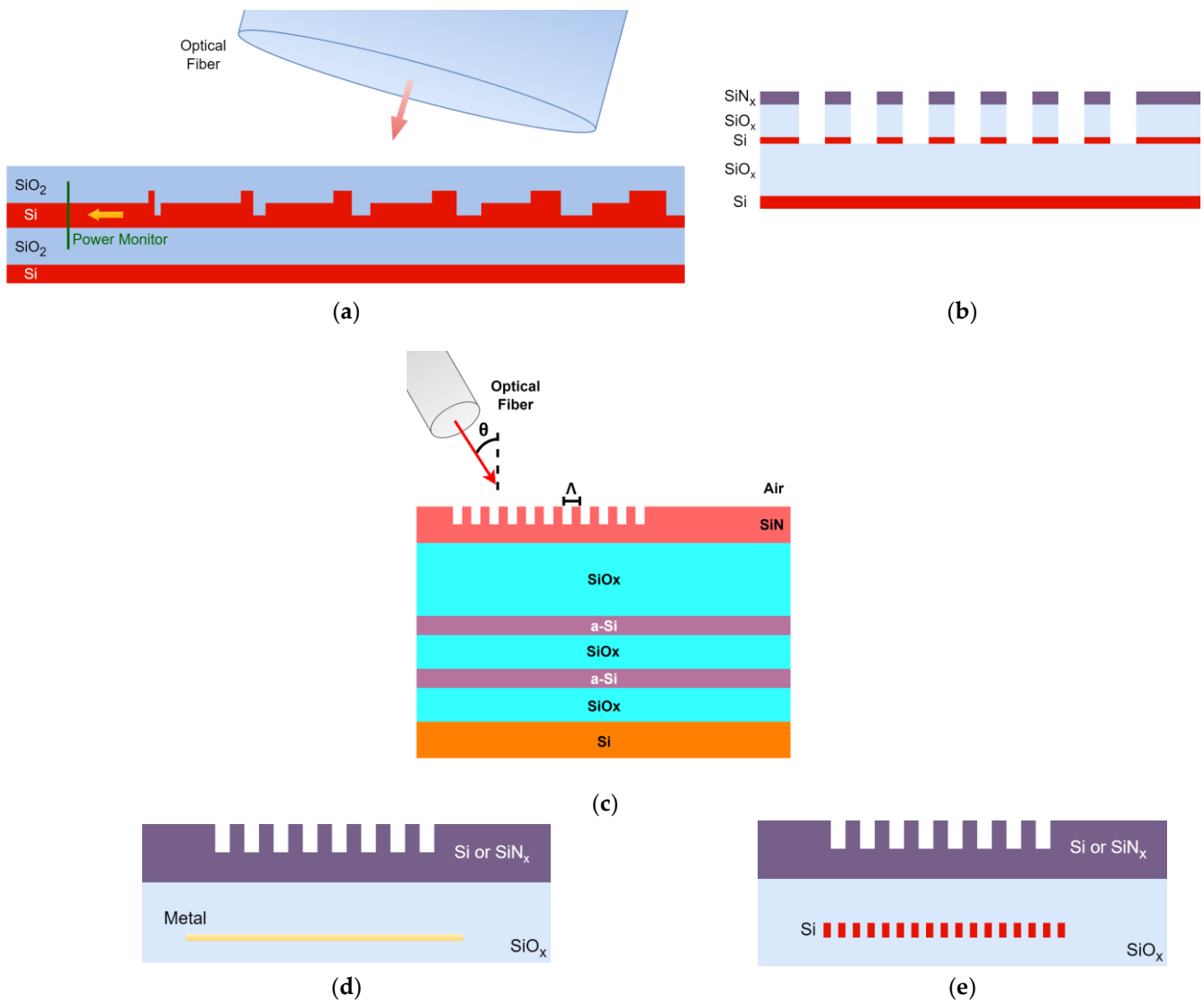
**Figure 2.** Photonic waveguide light-coupling techniques: (a) adiabatic coupler, consisting of tapered optical-fiber termination and waveguide taper; (b) non-apodized grating coupler, including optical fiber (with only the core and reduced cladding diameter represented), grating and waveguide taper; (c) optical fiber with tapered termination; (d) spot size converter based on 3D taper (optical fiber coating not represented, cladding diameter reduced); (e) spot size converter based on inverted taper (only the fiber core is represented). Representations not up to scale and not depicting functional devices, designed for demonstration purposes only.

### Grating Couplers

There are several variations of grating couplers [48,49] specifically designed to improve coupling efficiency [39,40,50–52] and bandwidth [52–54], such as the following:

- Dual-level grating couplers [40];
- Gratings featuring stacked layers [49,55];
- Grating couplers with distributed Bragg reflectors (DBR) [53,54,56];
- Grating couplers based on metamaterials [41];
- Gratings with bottom metal reflector [41,57–59];
- Grating couplers with bottom grating reflector [52].

Designs presented in Figure 3a,b,d,e can achieve remarkable performance; however, some of these (Figure 3b–e) require multiple deposition steps, increasing fabrication complexity and cost. Gratings based on metamaterials typically require a specific set of equipment that is not found in most typical semiconductor laboratory setups.



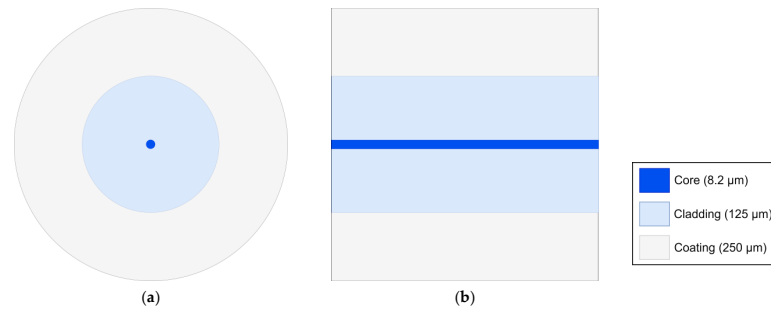
**Figure 3.** Some of the techniques employed to improve grating-coupler efficiency and/or bandwidth. (a) Dual-level grating coupler [40]; (b) grating coupler featuring stacked layers (SiN<sub>x</sub> on SOI) [49]; (c) silicon nitride grating coupler featuring a DBR stack composed of two layers of amorphous silicon (a-Si) [56]; (d) grating coupler featuring a bottom metal reflector layer [57,58]; (e) grating coupler with silicon bottom grating reflector [52].

To simplify the etching process, we designed a fully etched grating coupler (i.e., the silicon layer is etched to the substrate/under cladding) as an alternative to the more common partially etched designs, which require two lithographic processes/masks [39,54,60] or even more [40,42,52]. With the increased number of lithographic masks, misalignment impacts must be studied, since these can have an impact on the grating coupler’s performance [40].

### 3. Materials and Methods

#### 3.1. Single-Mode Optical Fiber

Considering a design operating wavelength of 1550 nm for the waveguide and light coupler, it is necessary to consider the effects of the coupling optical fiber. The fiber of choice is the SMF-28, due to its ultra-low propagation loss in the optical C-band (1530–1565 nm) and also for its ubiquity, being a very common type of optical fiber employed in optical communications, Figure 4.



**Figure 4.** Two-dimensional representations of the SMF-28 optical fiber. (a) Transversal section; (b) longitudinal cut. Both figures show (from center to periphery), core, cladding and coating. Coating is represented with a +5 μm deviation from the nominal value.

A summary of the typical optical fiber parameters at a wavelength of 1550 nm is presented in Table 1. Since the refractive index of the optical fiber’s core and refractive index difference are known, it is possible to determine the average cladding refractive index, which is about 1.4629. The refractive index values are similar to the figures presented by Saktioto et al. for an SMF28e optical fiber [61].

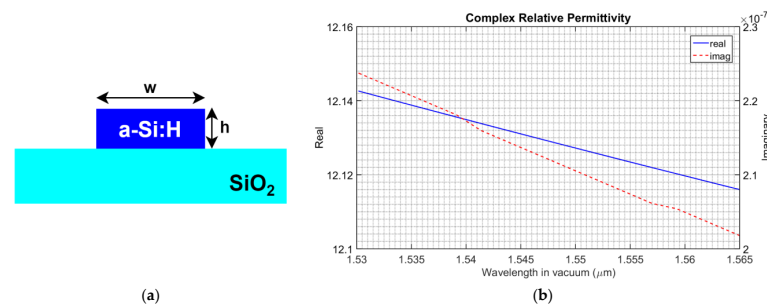
**Table 1.** SMF-28 Typical Specifications and Performance Characteristics. Data from [62].

Parameter	Value
Attenuation	≤0.22 dB/km
Mode Field Diameter (MFD)	10.4 ± 0.8 μm
Core Diameter	8.2 μm
Cladding Diameter	125 ± 0.7 μm
Coating Diameter	245 ± 5 μm
Effective Refractive Index ( $N_{eff}$ at rated MFD)	1.4682
Refractive Index Difference	0.36%

### 3.2. Single-Mode Waveguide

A 220 nm-thick hydrogenated amorphous silicon (a-Si:H) strip waveguide is used to couple light into the diffraction grating coupler. The model of the photonic waveguide material is the GUTL (Gauss–Urbach–Tauc–Lorentz) [63], which is based on the optical characteristics of a-Si:H. The complex refractive index at the operating wavelength (1550 nm) is approximately  $3.4824 + 3.023i \times 10^{-8}$ .

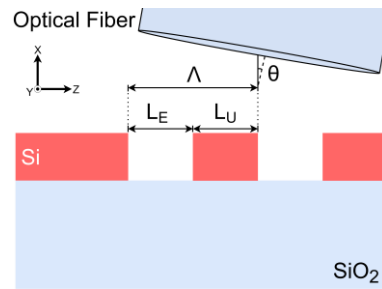
The 220 nm a-Si:H layer is deposited over silica, and the environment is air, Figure 5. Attenuation was not calculated, as it depends on the width of the waveguide, and simulations were performed using a 2D model.



**Figure 5.** Transversal cut of the proposed strip waveguide (a); the background is air, and a 220 nm hydrogenated amorphous silicon (a-Si:H) layer (represented in blue) is deposited over silica (represented in light blue). The SiO<sub>2</sub> is much thicker (by at least one order of magnitude) than the amorphous silicon layer (figure not up to scale). The waveguide’s height (thickness) is h, and the width is w, respectively. Since we are assuming a 2D model, the influence of width was not studied. Complex relative permittivity of hydrogenated amorphous silicon (b), given for the wavelengths of the optical C-band, values from the GUTL model [63].

### 3.3. Grating Coupler Design

The main design parameters of a grating coupler are period ( $\Lambda$ ), coupling angle ( $\theta$ ) and fill factor ( $F$ ), Figure 6. A grating coupler can have a constant effective refractive index (fill factor and period) throughout its length, in a non-apodized design, or have a variation of the effective refractive index (fill factor, period or both) over its span in an apodized variant.



**Figure 6.** Simplified representation of an optical fiber to silicon waveguide light-coupling scheme, showing some of the grating-coupler design parameters: coupling angle ( $\theta$ ), grating period ( $\Lambda$ ), etched length ( $L_E$ ) and unetched length ( $L_U$ ). This figure does not represent a functional coupling scheme and is not up to scale (e.g., the size mismatch between the optical fiber and the grating coupler is significantly greater).

The fill factor corresponds to the ratio of the unetched length ( $L_U$ ) to the grating period ( $\Lambda$ ), which can be expressed as a percentage (1):

$$F = \frac{L_U}{\Lambda}. \tag{1}$$

The effective refractive index of a diffraction grating coupler ( $n_{eff}$ ) can be obtained from its fill factor ( $F$ ), the refractive index of the etched ( $n_E$ ) and unetched segments ( $n_U$ ) of the grating, as described in expression (2) [50]:

$$n_{eff} = F \times n_U + (1 - F) \times n_E, \tag{2}$$

in this equation,  $n_E$  corresponds to the refractive index of air and  $n_U$  to the effective refractive index of the waveguide’s fundamental quasi-transverse electric (TE) mode of propagation. Having the effective refractive index,  $n_{eff}$ , it is possible to calculate the required grating-coupler period ( $\Lambda$ ), for a given coupling angle ( $\theta$ ), typically represented in degrees and free-space wavelength ( $\lambda_0$ ) as described in Equation (3) [64]:

$$\Lambda = \frac{\lambda_0}{n_{eff} - \sin \theta'} \tag{3}$$

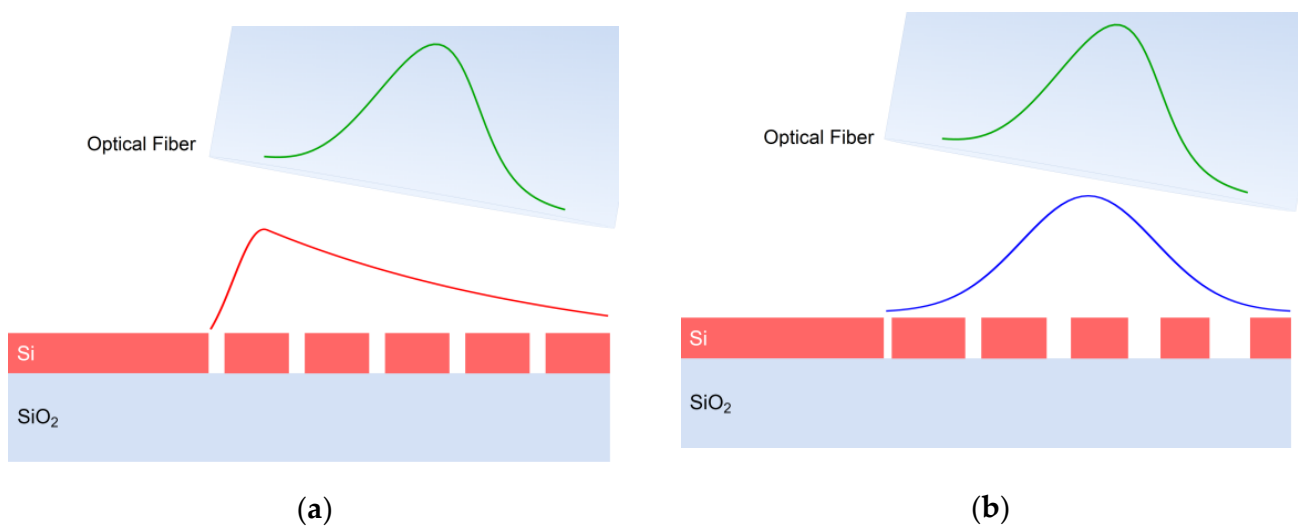
the resulting grating period has the same units as the free-space wavelength; both are typically represented in units of nanometer (nm).

Having the main design expressions, it is possible to begin the design of a grating coupler in its simplest form. By defining the fill factor ( $F$ ) as 0.5 (50%), while trying to keep the value of the denominator of expression (3) as low as possible, to achieve a higher period and, consequently, a large feature size, these measures can significantly decrease production costs if successfully implemented. The grating’s refractive index,  $n_{eff}$ , cannot be varied without changing the materials composing the grating coupler, the waveguide’s height or width, the fill factor or etch depth, so this value is regarded as a constant. The coupling angle,  $\theta$ , was set to a very high value of 22.5°. The operating wavelength is 1550 nm, a typical value used in optical fiber communications, situated close to the middle of the optical C-band. Since the grating is fully etched, without upper cladding, and simulated in an environment filled with air,  $n_E$  can be assumed as 1;  $n_U$  was obtained using the Finite

Element Method (FEM), having a value of approximately 2.8372. The effective refractive index of the grating coupler can now be obtained from Equation (2), having a value of about 1.9186. In the presence of all variables, it is now possible to calculate the grating’s period ( $\Lambda$ ), Equation (3), which is approximately 1009.2 nm (1.0092  $\mu\text{m}$ ), resulting in a feature size of about 505 nm, a value close to the required width of a silicon single-mode strip waveguide operating at 1550 nm [30,31]. A grating coupler with these characteristics has a relatively narrow diffracted field. The grating coupler was simulated using the 2D-FDTD method and the coupling efficiency (SMF-28 optical fiber) measured under  $-9$  dB. There was some discrepancy in the value of the coupling angle, which was found (in simulation) to be about  $16.5^\circ$  for optimum coupling. The values for coupling efficiency and coupling angle were obtained from the 2D-FDTD simulation tool (FullWAVE) provided with the RSoft software package (Synopsys, Inc., Sunnyvale, CA, USA). Comparing with uniform grating couplers reported in the literature [65,66], the simulated efficiency,  $-9.7$  dB, is relatively low when contrasted with  $-7$  dB [66] and  $-5.8$  dB [65]. Since we are trying to maximize the feature size (in order to reduce production costs), the theoretical diffraction angle is also much higher,  $22.5^\circ$ , when compared to uniform grating coupler designs reported in the literature, of  $8^\circ$  [65] and  $10^\circ$  [66].

### 3.4. Grating-Coupler Optimization

Without apodization, grating couplers exhibit an exponential decay of the diffracted field over its length [55] (Figure 7a); due to this fact, most authors tend to prefer designs employing some form of apodization [39,40,51]. With this technique, the diffracted field distribution becomes more like a Gaussian throughout the grating coupler’s length [55], i.e., with the field distribution commonly found in optical fibers (Figure 7b) resulting in more efficient coupling.



**Figure 7.** Hypothetical diffracted field-intensity distributions of (a) non-apodized (red line) and (b) apodized (blue line) grating couplers; both figures display a theoretical fiber model with a Gaussian field distribution (green lines). The figures are not up to scale and do not represent functional designs.

In the most common approach, gratings feature linear apodization (i.e., the fill factor varies linearly throughout the grating’s length), Equation (4) [50].

$$F(z) = F_0 - R \times z, \tag{4}$$

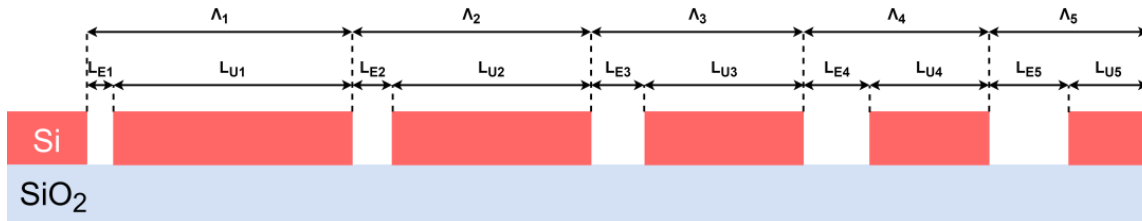
where  $F_0$  corresponds to the grating’s initial fill factor (refer to expression (1)),  $R$  is the linear apodization factor, typically expressed in units of  $\mu\text{m}^{-1}$  or  $\text{nm}^{-1}$ , and  $z$  corresponds



to the distance from the beginning of the grating, usually in units of  $\mu\text{m}$  (if  $R$  is represented in  $\mu\text{m}^{-1}$ ) or  $\text{nm}$  (when  $R$  is represented in  $\text{nm}^{-1}$ ), respectively.

### 3.5. Linear Refractive-Index Variation

A linear variation of the grating’s refractive index can be achieved by varying both the period and fill factor over the grating coupler’s length, Figure 8, while maintaining the coupling angle constant.



**Figure 8.** Hypothetical apodized grating coupler with varying refractive index over length; both period and fill factor vary. This is not a functional design and was not based on any distribution of the refractive index, linear, polynomial, or other. Typically, grating couplers have dozens of periods. Representation not up to scale.

The expression for the linear variation of the effective refractive index over a grating’s length can be written as (5)

$$n_{eff}(z) = n_{eff}(0) + m \times z, \tag{5}$$

where  $n_{eff}(0)$  is the effective refractive index at the start of the grating coupler (the effective refractive index of the first grating period),  $m$  corresponds to the slope of the refractive index variation, in  $\mu\text{m}^{-1}$ , and  $z$  is the position measured from the start of the grating.

Following optimization in 2D-FDTD software, expression (6) was obtained for a 220 nm-thick fully etched grating coupler of hydrogenated amorphous silicon (a-Si:H), operating at a wavelength of 1550 nm:

$$n_{eff}(z) = 2.65348 - 0.0288z. \tag{6}$$

The value of  $n_{eff}(0)$ , 2.65348, was obtained considering a maximum fill factor of 0.9 (90%), an etched-section refractive index,  $n_E$ , of 1 and an unetched-section effective refractive index,  $n_U$ , of 2.8372 (refer to expression (2)). The slope was adjusted considering a correlation which was observed between the length of the diffracted field distribution (Gaussian-like function) and the total variation of the refractive index over the grating coupler’s length. It was found on a 2D-FDTD simulation that for maximum coupling efficiency, the variation of the refractive index,  $\Delta n_{eff}$ , over a length of one optical fiber MFD ( $10.4 \mu\text{m}$ ),  $\Delta z$ , had to be of about  $-0.3$ , resulting in a slope,  $m$ , of approximately  $-0.0288 \mu\text{m}^{-1}$ . (Refer to expression (7) and plot of Figure 9.)

$$m = \frac{\Delta n_{eff}}{\Delta z} \left[ \mu\text{m}^{-1} \right]. \tag{7}$$

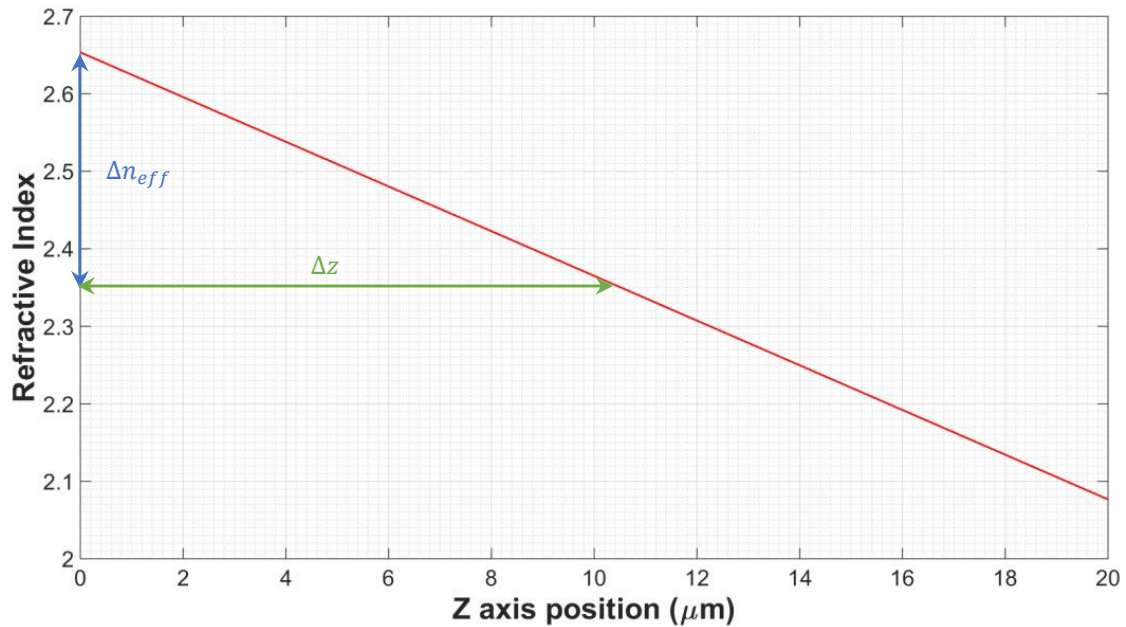
The plot of function (6) is represented in Figure 9, for a  $20 \mu\text{m}$  grating coupler.

Accounting for the variation of the refractive index over the length of the grating coupler, it is possible to rearrange (2) to obtain the fill factor for each grating position (8):

$$F(z) = \frac{n_{eff}(z) - n_E}{n_U - n_E}. \tag{8}$$

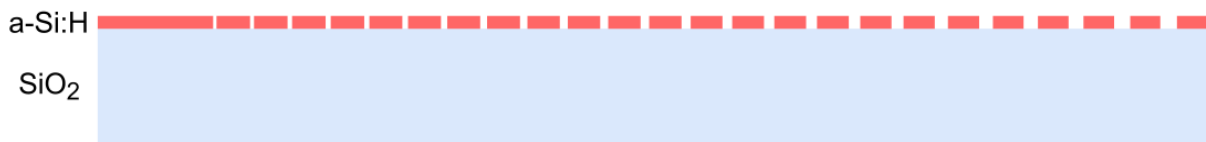
The expression which gives the grating period (3) can also be rewritten to account for the variation of the effective refractive index over the grating’s length (9):

$$\Lambda(z) = \frac{\lambda_0[\text{m}]}{n_{eff}(z) - \sin \theta} [\text{m}]. \tag{9}$$



**Figure 9.** Proposed linear variation of the refractive index over the grating coupler’s length (represented in red).

A computational algorithm, implemented in MATLAB (The MathWorks, Inc., La-Grange, GA, USA) obtains the effective refractive-index values at the begin of each grating period, using expression (6), and calculates the fill factor,  $F$  from function (8), and period,  $\Lambda$ , using expression (9) at each iteration, until the end of the grating coupler. The diffraction angle was defined as  $15^\circ$ . The pseudo-code algorithm is presented in Appendix A: this code sketch is used to represent the grating coupler in CAD, based on its geometric design parameters (fill factor and period) for each grating-segment iteration. The resulting grating coupler is presented in Figure 10. The length can overshoot the predefined value; in this case, for example, it is approximately  $20.6 \mu\text{m}$ , instead of  $20 \mu\text{m}$ . It was decided that it is preferable to have a slightly longer grating coupler than the opposite.



**Figure 10.** Grating coupler with a linear variation of the refractive index: the fill factor and period vary over the grating’s length, the refractive index is the highest close to the access waveguide (left) and the lowest at the end of the grating coupler (right). In this figure, only 24 segments are represented. Representation is accurate to 1/100th of a micrometer ( $\pm 10 \text{ nm}$  tolerance per segment).

The fill factor ( $F$ ) and period ( $\Lambda$ ) for each grating segment, numbered from the beginning to the end of the grating coupler (or from left to right in Figure 10), considering a total of 28 segments, are presented in Table 2.

**Table 2.** Period ( $\Lambda$ ) and fill factor ( $F$ ) for each segment of the grating coupler with linear variation of the refractive index.

Segment Number	$\Lambda$ [nm]	$F (L_E/\Lambda)$	Segment Number	$\Lambda$ [nm]	$F (L_E/\Lambda)$
1	647	0.900	15	732	0.750
2	652	0.890	16	739	0.738
3	658	0.880	17	747	0.727
4	663	0.869	18	754	0.715
5	668	0.859	19	763	0.703
6	674	0.848	20	771	0.691
7	680	0.838	21	779	0.679
8	686	0.827	22	788	0.667
9	692	0.816	23	798	0.654
10	698	0.805	24	807	0.642
11	704	0.795	25	817	0.629
12	711	0.783	26	827	0.616
13	718	0.772	27	838	0.603
14	725	0.761	28	849	0.590

### 3.6. Quadratic Refractive-Index Variation

The refractive index of the grating coupler can also have a quadratic variation; in this case, the function which gives the refractive index over the grating’s length is defined as follows (10):

$$n_{eff}(z) = n_{eff}(0) + a \times z^2 + b \times z, \tag{10}$$

following the same approach as in the previous subchapter, the initial fill factor, if defined as 0.9, results in an  $n_{eff}(0)$  of 2.65348,  $z$  corresponds to the grating position relative to its starting point, and  $a$  and  $b$  are the coefficients of the quadratic polynomial, with units  $\mu\text{m}^{-2}$  and  $\mu\text{m}^{-1}$ , respectively. As in the linear function, a  $\Delta n_{eff}$  of  $-0.3$  is pretended for a  $\Delta z$  of  $10.4 \mu\text{m}$  (measured from the grating’s starting point,  $z = 0$ ). The vertex of the parabola is set to occur at or near the end of the grating coupler, at  $z_{max}$ ,  $20 \mu\text{m}$ . With known values for  $\Delta n_{eff}$  ( $-0.3$ ),  $\Delta z$  ( $10.4 \mu\text{m}$ ) and  $z_{max}$  ( $20 \mu\text{m}$ ), it is possible to calculate the values of the coefficients  $a$  and  $b$ , recurring to a system of linear equations (11):

$$\begin{cases} \Delta n_{eff} = n_{eff}(\Delta z) - n_{eff}(0) \\ -\frac{b}{2a} = z_{max} \end{cases}, \tag{11}$$

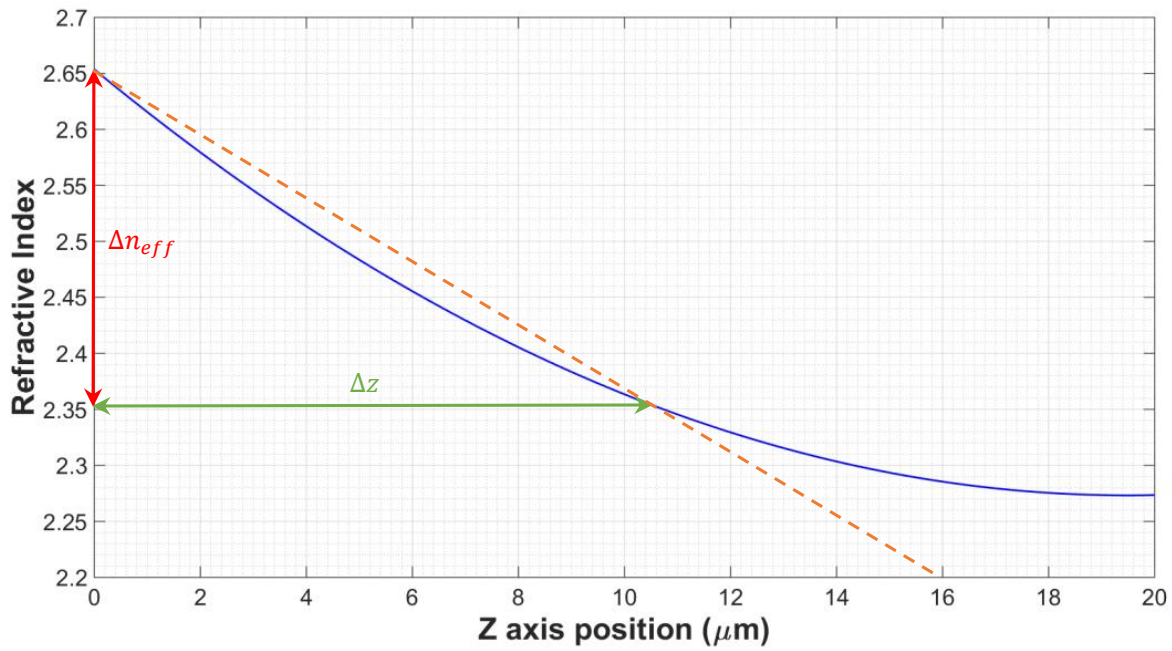
which can be rearranged as (12)

$$\begin{cases} \Delta n_{eff} = a \times \Delta z^2 + b \times \Delta z \\ -\frac{b}{2a} = z_{max} \end{cases}. \tag{12}$$

By solving the linear system, the values obtained for  $a$  and  $b$  are of approximately  $9.745 \times 10^{-4} \mu\text{m}^{-2}$  and  $-3.898 \times 10^{-2} \mu\text{m}^{-1}$ , respectively. Equation (10) can now be rewritten with the approximated values of the quadratic function coefficients, which were rounded to the 3rd decimal place (13):

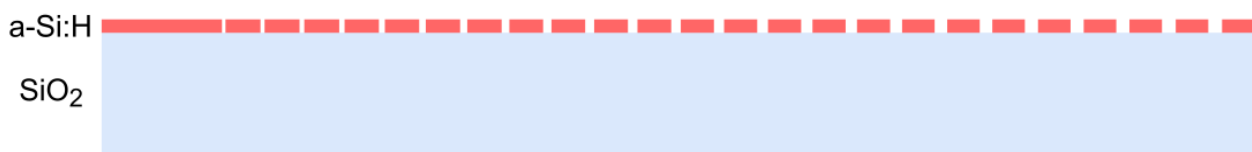
$$n_{eff}(z) = 2.65348 + 0.001z^2 - 0.039z. \tag{13}$$

The plot of the quadratic function which gives the effective refractive index for each position along the grating’s length is represented in Figure 11.



**Figure 11.** Proposed quadratic variation of the refractive index over the grating coupler’s length, represented in blue. The orange dashed line corresponds to the linear regression obtained from the quadratic-function values at position 0 and position 10.4  $\mu\text{m}$ .

As in the linear variation of the refractive index, a computational algorithm was written in MATLAB (The MathWorks, Inc., LaGrange, GA, USA), to calculate the fill factor and period for each segment of the grating coupler, using expressions (8), (9) and (13). The pseudo-code is presented in Appendix B, this code sketch is used to represent the grating coupler in CAD based on its geometric design parameters (fill factor and period) for each grating-segment iteration. The grating coupler of Figure 12 was obtained after running the algorithm. As in the case of the grating coupler designed for linear refractive-index variation, the length of the grating slightly overshoots the preset value of 20  $\mu\text{m}$ , achieving a length of about 20.25  $\mu\text{m}$  over 28 periods (segments).



**Figure 12.** Grating coupler with a quadratic variation of the refractive index: the fill factor and period vary over the grating’s length. The refractive index is the highest close to the waveguide (left) and the lowest at the end of the grating coupler (right). Only 24 segments are represented. Representation accurate to 1/100th of a micrometer ( $\pm 10$  nm tolerance per segment).

The period ( $\Lambda$ ) and fill factor ( $F$ ) of the quadratic refractive-index variation grating coupler is presented in Table 3, for a total of 28 segments.

**Table 3.** Period ( $\Lambda$ ) and fill factor ( $F$ ) for each segment of the grating coupler with quadratic variation of the effective refractive index.

Segment Number	$\Lambda$ [nm]	$F (L_E/\Lambda)$	Segment Number	$\Lambda$ [nm]	$F (L_E/\Lambda)$
1	647	0.900	15	734	0.746
2	654	0.886	16	739	0.738

Table 3. Cont.

Segment Number	$\Lambda$ [nm]	$F (L_E/\Lambda)$	Segment Number	$\Lambda$ [nm]	$F (L_E/\Lambda)$
3	661	0.873	17	744	0.731
4	668	0.860	18	748	0.725
5	674	0.848	19	752	0.719
6	681	0.836	20	755	0.713
7	687	0.824	21	759	0.709
8	694	0.813	22	761	0.705
9	700	0.802	23	764	0.701
10	706	0.791	24	766	0.698
11	712	0.781	25	767	0.696
12	718	0.772	26	769	0.694
13	724	0.763	27	769	0.693
14	729	0.754	28	769	0.693

### 3.7. Fill Factor Greater than 50%

An alternative to apodization is to use a fill factor different to 0.5 (50%); the fill factor was increased by over  $\frac{1}{2}$  in an attempt to decrease the decay rate of the diffracted-field intensity over the grating coupler’s length. In our observations of simulations using the 2D-FDTD method, there was evidence that a correlation exists between the fill factor and the exponential decay of the diffracted field. The higher the fill factor, the slower the decrease in the field over the grating’s length; however, there is also a drawback, since the field intensity (per unit area) decreases as the fill factor increases. The best result of our optimization process was obtained for a grating coupler with a period of 690 nm and a fill factor of approximately 82.6%, with a coupling angle ( $\theta$ ) of about  $13.3^\circ$ , Figure 13.

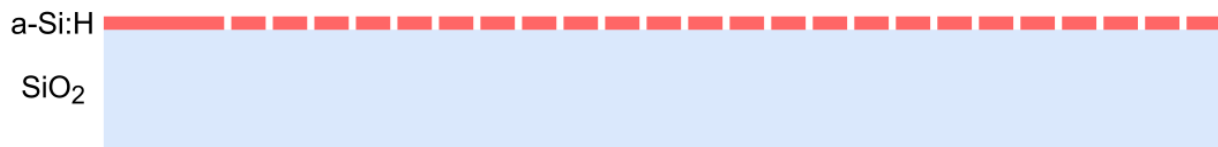


Figure 13. Grating coupler with a period of 690 nm and a fill factor of 82.6%. First 24 periods represented, coupling angle not represented.

### 3.8. Overlapped-Grating Design

Another interesting method consists in the superposition (overlap) of two grating couplers with a fill factor of 0.5 (50%); using this method, it is possible to fabricate gratings using two lithography processes (and/or masks). This method enables the fabrication of grating couplers with fill factors different from  $\frac{1}{2}$ , from designs with periods over 1000 nm, allowing the usage of masks with larger feature sizes and consequently lowering production costs. The coupling efficiency which results from this technique achieves values over  $-8$  dB, as found in 2D-FDTD simulations (FullWAVE) performed with the RSoft software package (Synopsys, Inc., Sunnyvale, CA, USA) (refer to Sections 4.7 and 5). The coupling efficiency obtained with this technique,  $-7.5$  dB, is more similar to the value reported by Laere et al. [66], of  $-7$  dB.

During the optimization process, the highest coupling efficiency (considering an SMF-28 optical fiber), was obtained for the overlap of two 1500 nm-period diffraction gratings, with a  $-400$  nm  $z$  axis offset, Figure 14, resulting in a fill factor of 0.7(6) and an optimum coupling angle of  $22^\circ$ . The minimum feature size is 750 nm. The same grating coupler, if fabricated using a single lithography procedure, would have a feature size of 350 nm, requiring higher precision in mask fabrication.



**Figure 14.** Grating coupler made from the superposition of two diffraction gratings with a period of 1500 nm; the result is a grating with the same period and a fill factor of 76.(6)%. Grating segments with and without offset are represented by the colors green and red, respectively.

### 3.9. Random-Distribution Overlapped Grating Design

Following a similar approach as in the previous subchapter, overlapped grating couplers with chaotic (random) distribution were also simulated. To decrease production costs, the period of the gratings was set between 1  $\mu\text{m}$  and 1.5  $\mu\text{m}$ , and varied by steps of 0.1  $\mu\text{m}$ . The combination which results in the highest efficiency consists in the overlap of 1.3  $\mu\text{m}$  over a 1.5  $\mu\text{m}$  grating, with a z-axis offset of  $-0.3 \mu\text{m}$  applied to the grating with the smallest feature size, Figure 15. For simplification purposes, both gratings have a fill factor of 0.5 (50%). The optimized coupling angle is about  $20.9^\circ$ .

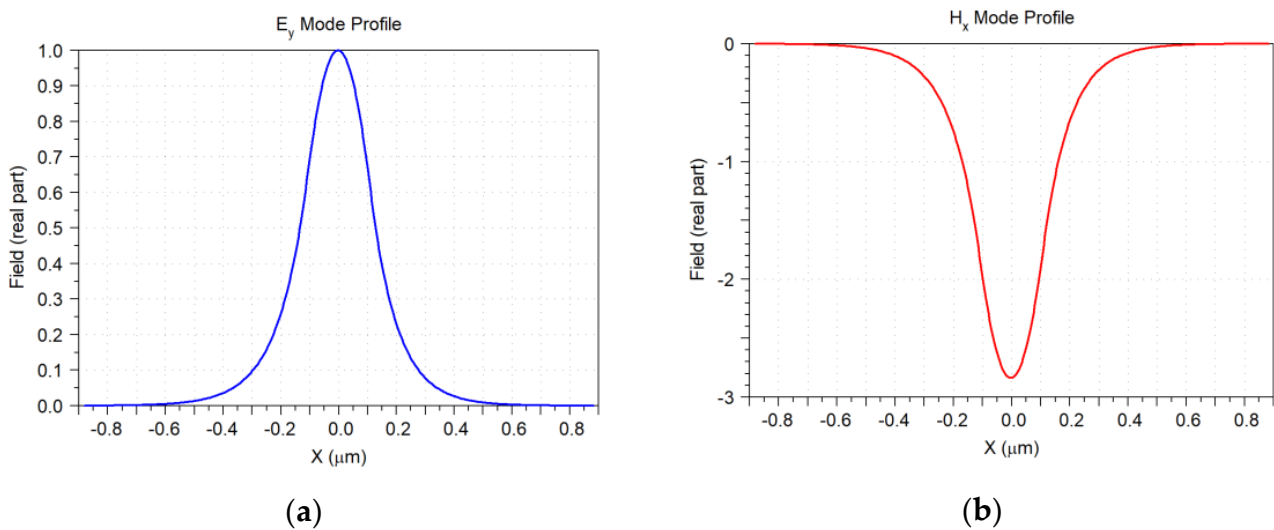


**Figure 15.** Grating coupler made from the superposition of two diffraction gratings with periods of 1500 nm (represented in red) and 1300 nm (represented in blue). The result is a grating with a chaotic (random) distribution.

## 4. Results

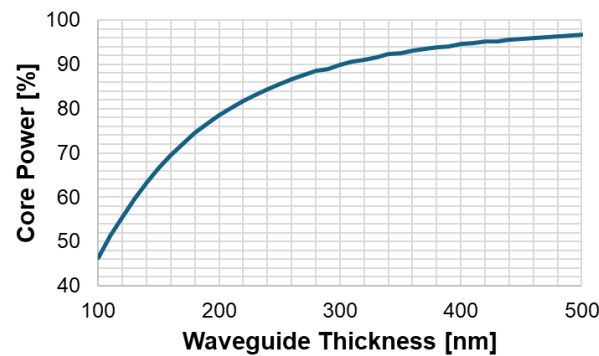
### 4.1. Fundamental Mode of the Waveguide

All grating couplers covered in this work are optimized for the fundamental quasi-transverse electric mode (quasi- $\text{TE}_{00}$ ). The fundamental mode of the 220 nm-strip waveguide was analyzed using the 2D-FEM (two-dimensional finite element method). The effective refractive index of the quasi- $\text{TE}_{00}$  mode of the waveguide is approximately 2.8372, with an extinction coefficient of  $3.037 \times 10^{-8}$  for a 1550 nm operating wavelength. The intensity of the electric ( $E_y$ ) and magnetic ( $H_x$ ) major fields along the  $x$  axis is represented in Figures 16a and 16b, respectively.



**Figure 16.** Electric,  $E_y$  (a) and magnetic,  $H_x$  (b) field distributions of the fundamental quasi-transverse electric mode (quasi- $\text{TE}_{00}$ ) of the waveguide.

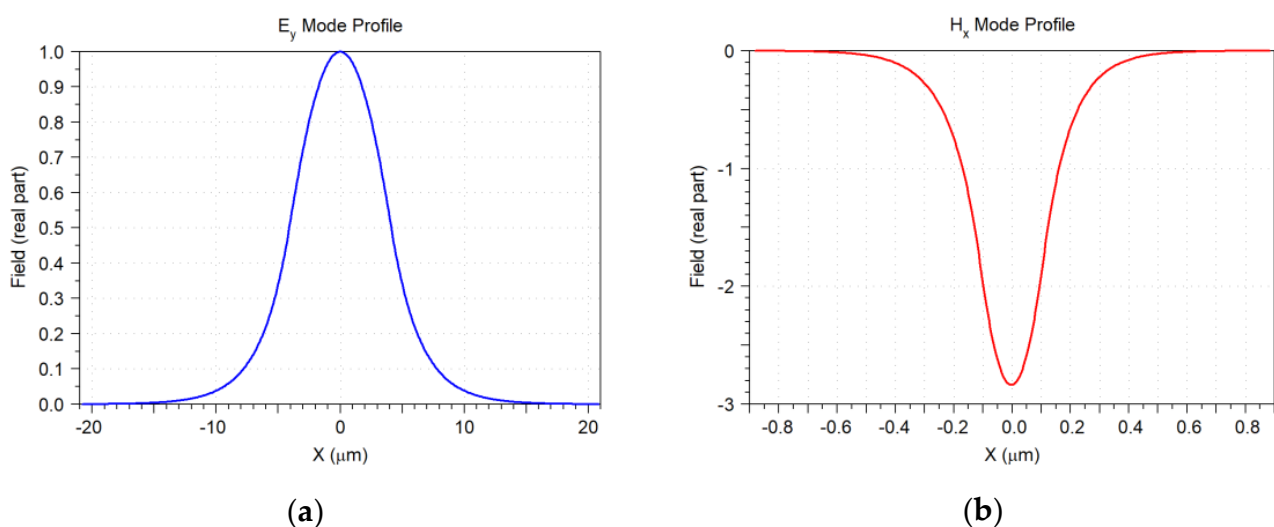
Variation of waveguide light confinement was simulated in 2D-FDTD for the quasi-TE<sub>00</sub> mode for different thickness values of the a-Si:H layer thickness (height), considering an operating wavelength in vacuum of 1550 nm. With a 220 nm-thick layer, approximately 81.8% of the total waveguide power is inside the a-Si:H core. The waveguide light confinement follows a logarithmic-like distribution, Figure 17; lower a-Si:H thickness values result in more power being transferred to the substrate and cover (air), while the opposite is verified for high thickness values, where most of the launch power is carried in the waveguide's core. Sensing applications typically explore the evanescent field effect, and in such cases, a tradeoff must be made between power loss and light confinement.



**Figure 17.** Waveguide light confinement for the fundamental quasi-TE mode versus waveguide thickness (height) for a wavelength of 1550 nm.

#### 4.2. Simulated Single-Mode-Fiber Fundamental Mode

To estimate the coupling efficiency, it is necessary to obtain the field distribution for the typical parameters of the single-mode optical fiber (SMF-28) of the coupling setup (refer to Table 1). The optical fiber was simulated using 2D-FDTD, like the single-mode waveguide. The CAD design is identical to that presented in Figure 4b; for simplicity and algorithm-efficiency reasons, the effects of the coating layer were not considered. The results of the simulation of the fiber model are presented in Figure 18, showing the electric,  $E_y$  (a) and magnetic,  $H_x$  (b) field distributions along the  $x$  axis, for the TE<sub>00</sub> mode, at a wavelength of 1550 nm.



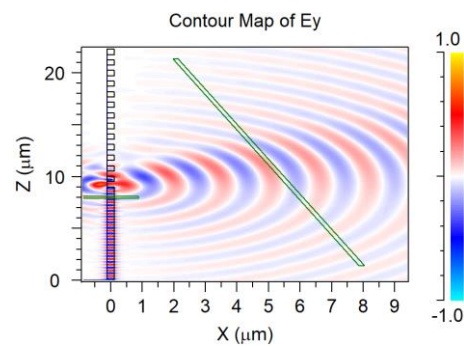
**Figure 18.** Electric (a) and magnetic (b) field distributions of the fundamental transverse electric mode (TE<sub>00</sub>) of the single-mode optical fiber.

### 4.3. Non-Optimized Grating-Coupler Performance

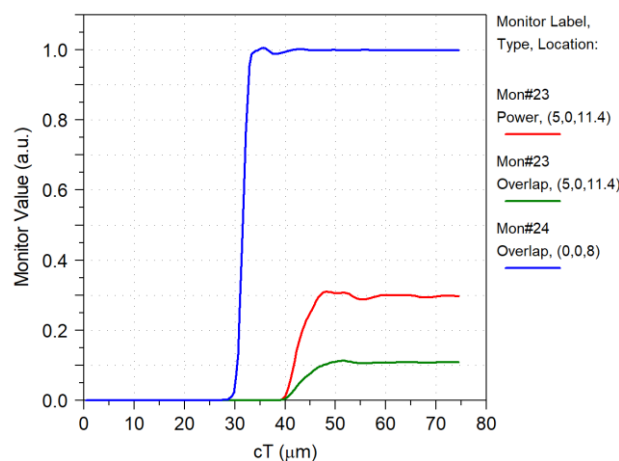
As previously mentioned, without apodization and with a fill factor of 0.5 (50%), the coupling efficiency is typically modest. In the simulation that we performed, a value slightly above  $-10$  dB was obtained for this parameter. It is thought to be of imperative importance that some additional measures are taken to improve grating-coupler performance.

Two monitors were used in the calculation of coupling efficiency, one centered at position  $(x, z) = (0, 8)$ , measuring the power near the output of the waveguide, and another centered at position  $(x, z) = (5, 11.4)$ , to measure the overlap between the diffracted field and the fundamental mode of the optical fiber. This second monitor is placed at an angle of  $16.5^\circ$  (optimized coupling angle): its center is at a distance of  $5 \mu\text{m}$  from the grating center axis and at a distance of  $2.4 \mu\text{m}$  from the grating coupler’s starting point, measured on the vertical  $z$  axis. The access waveguide is  $9 \mu\text{m}$  long.

In Figure 19, it is noticeable that the diffracted electric-field intensity decays at a very high rate (exponentially), following the two first periods of the grating coupler; due to this fact, coupling efficiency, Figure 20, is relatively low. In Figure 20, and in subsequent figures, the abscissa label “ $cT$ ” corresponds to a time-related unit, which is given as the product of the velocity of light in vacuum ( $c$ ) and time ( $\tau$ ). Since the velocity of light is constant, this unit represents the space travelled by light for a given time interval, in units of micrometer ( $\mu\text{m}$ ).



**Figure 19.** Representation on the XZ plane of the diffracted electric field: the green (transparent filled) horizontal bar is the waveguide’s monitor and the green (transparent filled) slanted bar is the diffracted field monitor. Image obtained from 2D-FDTD.



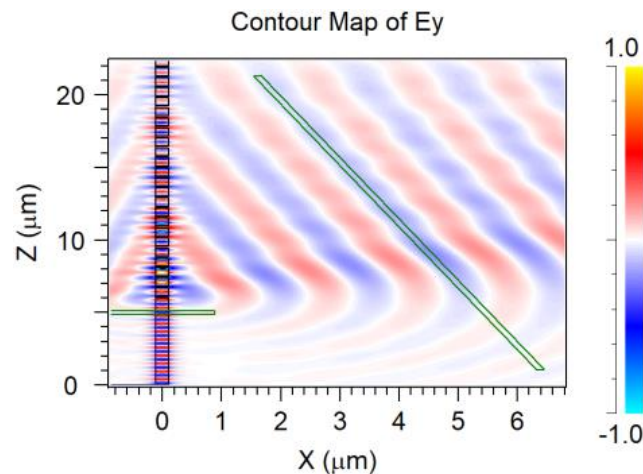
**Figure 20.** Efficiency test: from top to bottom, the blue line represents the overlap value (with the quasi- $\text{TE}_{00}$  waveguide’s mode) measured at  $1 \mu\text{m}$  before the waveguide’s end. The red and green lines represent the power and overlap (with the  $\text{TE}_{00}$  mode of a simulated SMF-28 optical fiber) of the diffraction field, respectively, measured by a monitor whose center is at a distance of  $5 \mu\text{m}$  from the grating’s center.



Looking at the results presented in Figure 20, it is possible to assert that the waveguide-monitor overlap (with the launch field, or fundamental quasi-TE<sub>00</sub> mode) remains nearly constant throughout the integration time. The coupling-efficiency value achieved at the time of convergence is slightly over  $-9.7$  dB ( $\sim 10.8\%$ ). The overlap with the diffracted field (i.e., the percentage of power of the diffracted field passing through the monitor which is converted to the SMF-28 fundamental mode) is about 36.4%. The ratio of the diffracted power reaching the monitor to the waveguide’s power is slightly over  $-5.3$  dB ( $\sim 29.7\%$ ).

#### 4.4. Constant Refractive Index and Fill Factor Greater than 50%

The results of the grating coupler featuring an 82% fill factor and a period of 690 nm are presented herein. In Figure 21, the contour map of the diffracted electric field is presented. Two monitors are included, one centered at position  $(x, z) = (0, 5)$ , measuring the power near the output of the waveguide, and another centered at position  $(x, z) = (4, 11.2)$ , which measures the overlap between the diffracted field and the fundamental mode of the optical fiber. This second monitor is placed at an angle of  $13.3^\circ$  (optimized coupling angle); its center is at a distance of  $4 \mu\text{m}$  from the grating center axis and at a distance of  $5.2 \mu\text{m}$ , measured from the grating coupler’s starting point. The access waveguide is  $6 \mu\text{m}$  long.

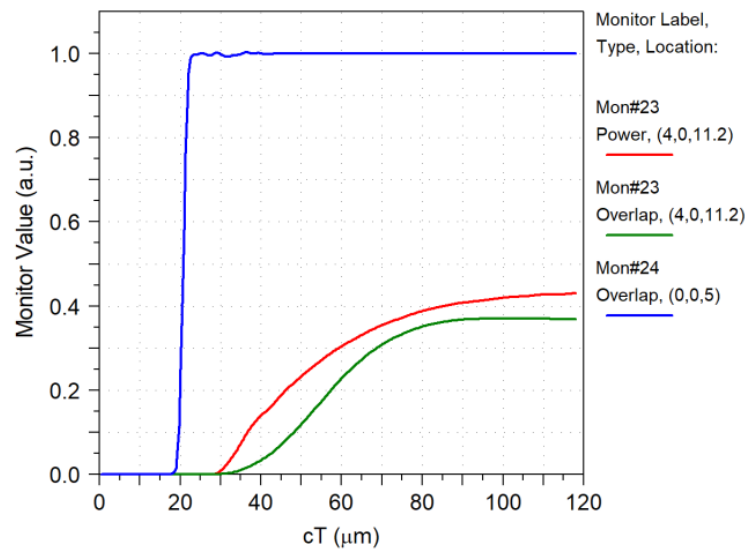


**Figure 21.** Representation on the XZ plane of the diffracted electric field: the green (transparent filled) horizontal bar is the waveguide’s monitor and the green (transparent filled) slanted bar is the diffracted field monitor. Image obtained from 2D-FDTD.

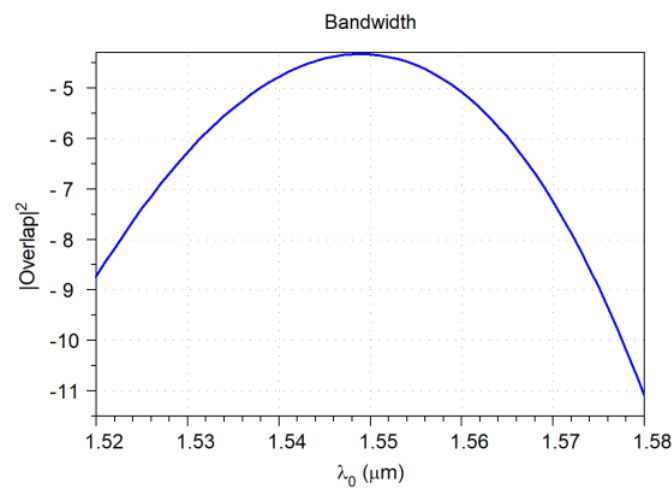
In Figure 21, it is possible to notice that the diffracted field intensity decays (exponentially) along the coupler’s length, a characteristic of non-apodized grating couplers. By optimizing the fill factor, monitor angle and position, it was possible to maximize the coupling efficiency.

Efficiency test results are shown in Figure 22. It is possible to verify that the waveguide-monitor overlap (with the launch field, or fundamental quasi-TE<sub>00</sub> mode) remains nearly constant throughout the integration time. The coupling-efficiency value achieved at the time of convergence is slightly under  $-4.3$  dB ( $\sim 36.9\%$ ). The overlap with the diffracted field (i.e., the percentage of power of the diffracted field passing through the monitor which is used to excite the SMF-28 fundamental mode) is about 85.6%. The ratio of the diffracted power reaching the monitor to the waveguide’s power is slightly over  $-3.7$  dB ( $\sim 43.1\%$ ).

Coupling efficiency was measured for wavelengths in vacuum between 1520 nm and 1580 nm, Figure 23. The  $-1$  dB bandwidth is about 25 nm and the  $-3$  dB bandwidth is 44 nm. The center wavelength (peak efficiency) occurs at 1549 nm. The optical C-band (1530–1565 nm) is covered with a maximum 2 dB penalty.



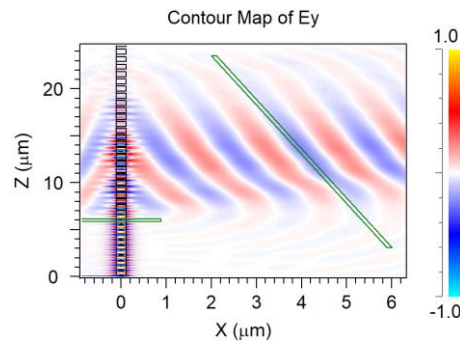
**Figure 22.** Efficiency test: from top to bottom, the blue line represents the overlap value (with the quasi-TE<sub>00</sub> waveguide’s mode) measured at 1 μm before the waveguide’s end. The red and green lines represent the power and overlap (with the TE<sub>00</sub> mode of a simulated SMF-28 optical fiber) of the diffraction field, respectively, measured by a monitor whose center is at a distance of 4 μm from the grating’s center.



**Figure 23.** Overlap with the SMF-28 optical fiber TE<sub>00</sub> mode at the diffracted field monitor, obtained for different wavelengths. The  $|\text{Overlap}|^2$  is given in logarithmic scale, decibel (dB), and wavelength in vacuum ( $\lambda_0$ ) is represented on a linear scale.

#### 4.5. Linear Effective-Refractive-Index Variation

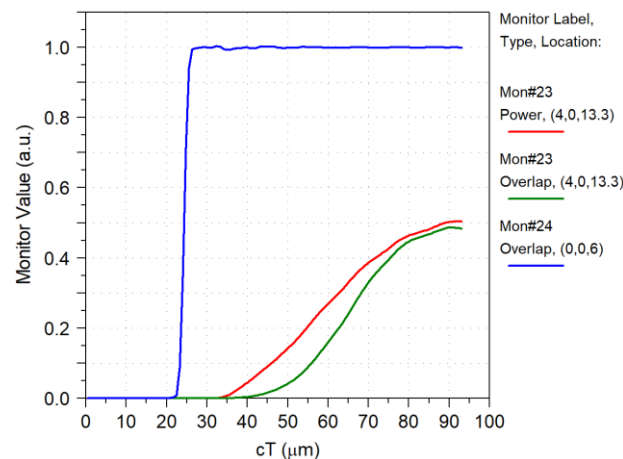
The results of the apodized grating coupler with linear variation of the effective refractive index are presented herein. In Figure 24, the distribution of the diffracted electric field is shown. Two monitors are included, one centered at position  $(x, z) = (0, 6)$ , measuring the power near the output of the waveguide, and another centered at position  $(x, z) = (4, 13.3)$ , measuring the overlap between the diffracted field and the fundamental mode of the optical fiber. The second monitor is positioned at an angle of  $10.8^\circ$  (optimized coupling angle); its central position is at a distance of 4 μm from the grating-center axis relative to the  $x$  axis and at a distance of 6.3 μm, measured from the grating coupler’s starting point on the  $z$  axis. The access waveguide is 7 μm long.



**Figure 24.** Representation on the XZ plane of the diffracted electric field: the green (transparent filled) horizontal bar is the waveguide’s monitor and the green (transparent filled) slanted bar is the diffracted field monitor. Image obtained from 2D-FDTD.

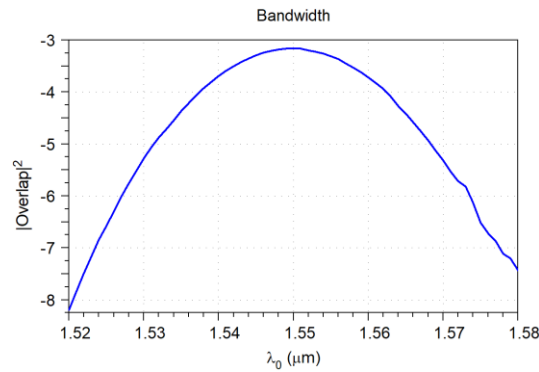
In Figure 24, it can be noticed that the diffracted field intensity has a distribution closer to a Gaussian, having a higher intensity near the center of the monitor, a characteristic of apodized designs. By optimizing apodization, monitor angle and position, it was possible to maximize coupling efficiency. The values of period and fill factor presented in Table 2 were limited to two decimal points in simulations, since most optical lithography systems do not have precisions greater than 10 nm. The impact of this loss of precision in the simulation results was found to be minimal.

In Figure 25, efficiency results are presented. It is possible to confirm that the waveguide monitor overlap (with the launch field, or fundamental quasi-TE<sub>00</sub> mode) remains nearly constant throughout the integration time. The coupling efficiency achieved at the time of convergence is slightly under  $-3.1$  dB ( $\sim 48.4\%$ ). The overlap with the diffracted field (i.e., the percentage of power of the diffracted field passing through the monitor which is used to excite the SMF-28 fundamental mode) is approximately 96.1%. The ratio of the diffracted power reaching the monitor to the waveguide’s power is slightly over  $-3$  dB ( $\sim 50.4\%$ ).



**Figure 25.** Efficiency test: from top to bottom, the blue line represents the overlap value (with the quasi-TE<sub>00</sub> waveguide’s mode) measured at 1 μm before the waveguide’s end. The red and green lines represent the power and overlap (with the TE<sub>00</sub> mode of a simulated SMF-28 optical fiber) of the diffraction field, respectively, measured by a monitor centered at a distance of 4 μm from the grating.

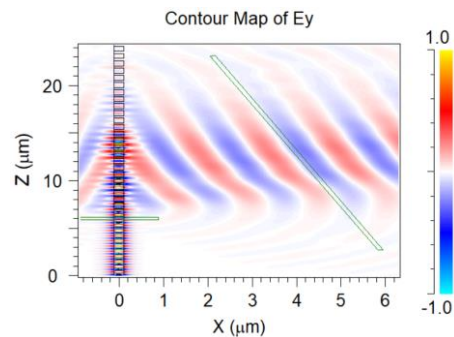
Coupling efficiency was measured for wavelengths in vacuum between 1520 nm and 1580 nm, Figure 26. The  $-1$  dB bandwidth is approximately 26 nm, and the  $-3$  dB bandwidth is 47 nm. Peak efficiency (center wavelength) occurs at 1550 nm. The optical C-band (1530–1565 nm) is covered with a maximum 2.1 dB penalty.



**Figure 26.** Overlap with the SMF-28 optical fiber TE<sub>00</sub> mode at the diffracted field monitor, obtained for different wavelengths. The  $|Overlap|^2$  is given in logarithmic scale, decibel (dB), and wavelength in vacuum ( $\lambda_0$ ) is represented on a linear scale.

#### 4.6. Quadratic Effective-Refractive-Index Variation

The performance results of the apodized grating coupler with quadratic (polynomial) variation of the effective refractive index are presented in this subchapter. In Figure 27, the distribution of the diffracted electric field is shown. Two monitors are included, one centered at position  $(x, z) = (0, 6)$ , measuring the power near the output of the waveguide, and another centered at position  $(x, z) = (4, 12.9)$ , measuring the overlap between the diffracted field and the fundamental mode of the optical fiber. The second monitor is positioned at an angle of  $10.5^\circ$  (optimized coupling angle); its center is at a distance of  $4 \mu\text{m}$  from the grating center axis, relative to the  $x$  axis, and at a distance of  $5.9 \mu\text{m}$ , measured from the grating coupler’s starting point, on the  $z$  axis. The waveguide has a length of  $7 \mu\text{m}$ .

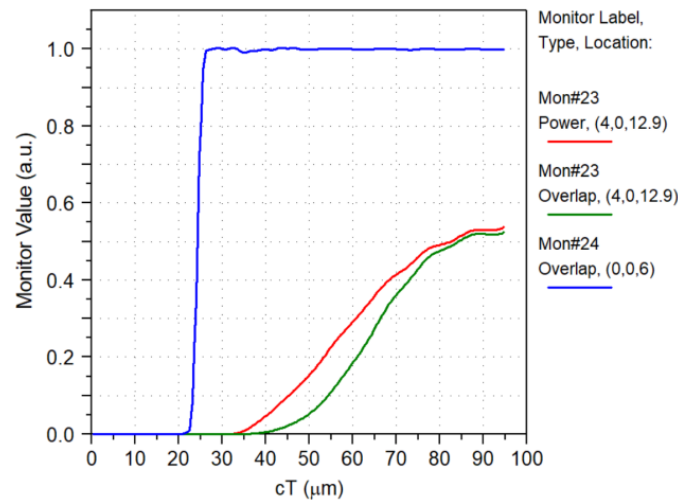


**Figure 27.** Representation on the XZ plane of the diffracted electric field: the green (transparent filled) horizontal bar is the waveguide’s monitor and the green (transparent filled) slanted bar is the diffracted field monitor. Image obtained from 2D-FDTD.

After looking carefully at Figure 27, it can be noticed that the diffracted field intensity has a Gaussian-like distribution, as in the previous apodized grating coupler, exhibiting higher intensity near the center of the monitor, a characteristic of apodized designs. By optimizing apodization, monitor angle and position, it was possible to maximize the coupling efficiency. The values of period and fill factor presented in Table 3 were limited to two decimal points in simulations, due to the fact that most optical lithography systems do not have precisions greater than  $10 \text{ nm}$ . The impact of this loss of precision was verified as having a negligible impact on the results.

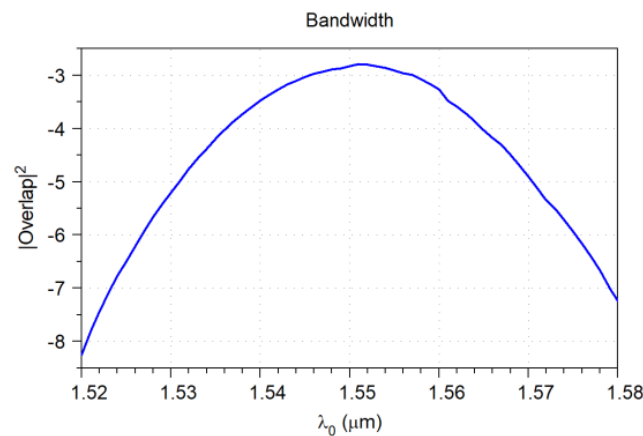
In Figure 28, coupling efficiency results are depicted. It is possible to verify that the waveguide monitor overlap (with the launch field, or fundamental quasi-TE<sub>00</sub> mode), remains nearly constant throughout the integration time. Coupling efficiency achieved at the time of convergence is slightly under  $-2.8 \text{ dB}$  ( $\sim 52.2\%$ ). The overlap with the diffracted field (i.e., the percentage of power of the diffracted field passing through the monitor which

is used to excite the SMF-28 fundamental mode) is around 97.4%. The ratio of the diffracted power reaching the monitor to the waveguide’s power is slightly under  $-2.7$  dB ( $\sim 53.6\%$ ).



**Figure 28.** Efficiency test: from top to bottom, the blue line represents the overlap value (with the quasi-TE<sub>00</sub> waveguide’s mode) measured at 1 μm before the waveguide’s end (diffraction-grating starting point). The red and green lines represent the power and overlap (with the TE<sub>00</sub> mode of a simulated SMF-28 optical fiber) of the diffraction field, respectively, measured by a monitor whose center is at a distance of 4 μm from the grating.

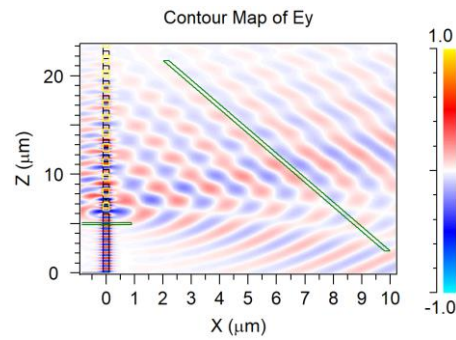
Coupling efficiency was measured for wavelengths in vacuum between 1520 nm and 1580 nm, Figure 29. The  $-1$  dB bandwidth is approximately 25 nm and the  $-3$  dB bandwidth is 46 nm. Peak efficiency (center wavelength) occurs at 1551 nm. The optical C-band (1530–1565 nm) is covered with a maximum 2.4 dB penalty.



**Figure 29.** Overlap with the SMF-28 optical fiber TE<sub>00</sub> mode at the diffracted field monitor, obtained for different wavelengths. The  $|\text{Overlap}|^2$  is given in logarithmic scale, decibel (dB), and wavelength in vacuum ( $\lambda_0$ ) is represented on a linear scale.

#### 4.7. Overlapped Micrometric Design

The coupling efficiency results for the overlapped micrometric period gratings (1500 nm) are presented herein. The two superimposed gratings with an offset of  $-400$  nm form a grating coupler with a period of 1500 nm and a fill factor of 0.7(6) (76(6)%). In Figure 30, the contour map of the diffracted electric field is presented.

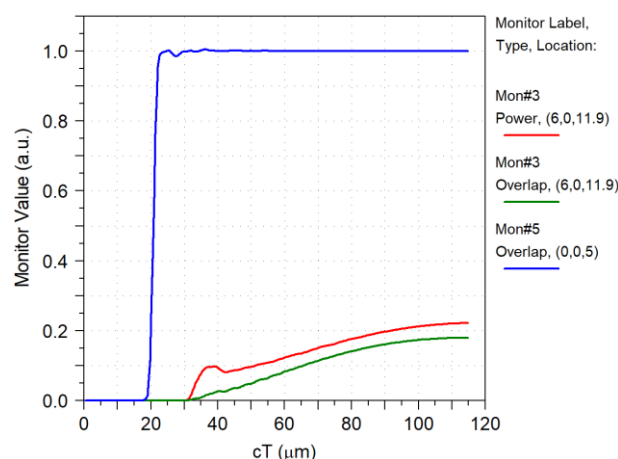


**Figure 30.** Representation on the XZ plane of the diffracted electric field: the green (transparent filled) horizontal bar is the waveguide’s monitor and the green (transparent filled) slanted bar is the diffracted field monitor. Image obtained from 2D-FDTD.

Two monitors are included, one centered at position  $(x, z) = (0, 5)$ , measuring the power near the output of the waveguide, and another centered at position  $(x, z) = (6, 11.9)$ , which measures the overlap between the diffracted field and the fundamental mode of the optical fiber. This second monitor is placed at an angle of  $22^\circ$  (optimized coupling angle); its center is at a distance of  $6 \mu\text{m}$  from the grating center axis (due to the higher coupling angle) and at a distance of  $5.9 \mu\text{m}$ , measured from the grating coupler’s starting point. The access waveguide is  $6 \mu\text{m}$  long.

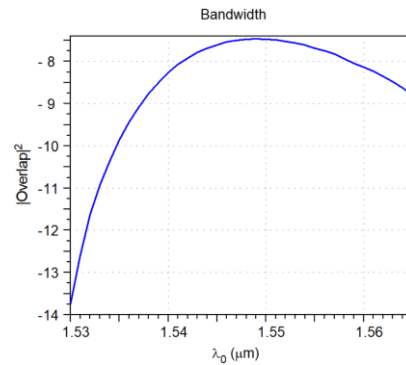
In Figure 30, it is possible to notice that the diffracted field intensity (exponentially) decays along the coupler’s length, a characteristic of non-apodized grating couplers. By optimizing the fill factor, monitor angle and position, it was possible to maximize the coupling efficiency.

In Figure 31, coupling efficiency results are shown. It is possible to attest that the waveguide-monitor overlap (with the launch field, or fundamental quasi-TE<sub>00</sub> mode), remains nearly constant throughout the integration time. Coupling efficiency achieved at the time of convergence is slightly over  $-7.5 \text{ dB}$  ( $\sim 20.0\%$ ). The overlap with the diffracted field (i.e., the percentage of power of the diffracted field passing through the monitor which is used to excite the SMF-28 fundamental mode) is approximately  $80.8\%$ . The ratio of the diffracted power reaching the monitor to the waveguide’s power is slightly under  $-6.5 \text{ dB}$  ( $\sim 22.3\%$ ).



**Figure 31.** Efficiency test: from top to bottom, the blue line represents the overlap value (with the quasi-TE<sub>00</sub> waveguide’s mode) measured at  $1 \mu\text{m}$  before the waveguide’s end (diffraction grating starting point). The red and green lines represent the power and overlap (with the TE<sub>00</sub> mode of a simulated SMF-28 optical fiber) of the diffraction field, respectively, measured by a monitor whose center is at a distance of  $6 \mu\text{m}$  from the grating.

Coupling efficiency was measured for wavelengths in vacuum between 1530 nm and 1565 nm, Figure 32. The  $-1$  dB bandwidth is approximately 22 nm. Peak efficiency (center wavelength) occurs at 1549 nm. The optical C-band (1530–1565 nm) is covered with a maximum 6.3 dB penalty.



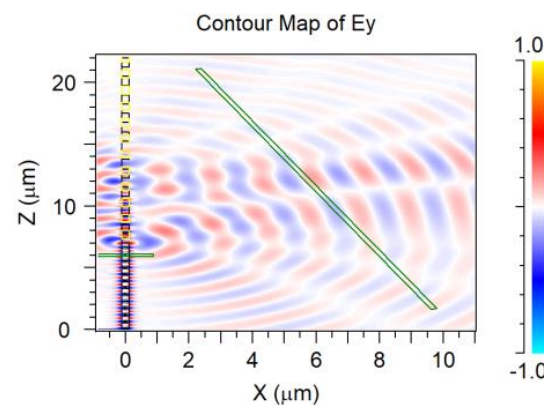
**Figure 32.** Overlap with the SMF-28 optical fiber  $\text{TE}_{00}$  mode at the diffracted field monitor, obtained for different wavelengths. The  $|\text{Overlap}|^2$  is given in logarithmic scale, decibel (dB), and wavelength in vacuum ( $\lambda_0$ ) is represented on a linear scale.

#### 4.8. Random Distribution

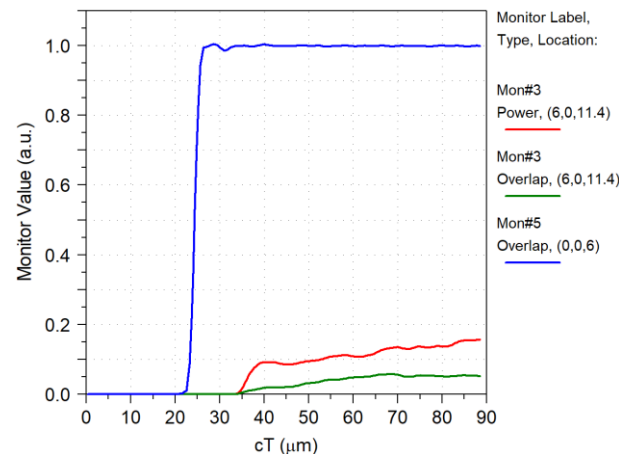
Results obtained for the random (chaotic)-distribution grating coupler are presented in this subchapter. Two gratings with periods of 1500 and 1300 nm were overlapped, with the second grating having a  $z$  axis offset of  $-300$  nm. This was the combination which held the best results for a grating coupler featuring a random distribution, achieved by the superposition of two gratings with periods over 1000 nm.

Two monitors are included, one centered at position  $(x, z) = (0, 6)$ , measuring the power near the output of the waveguide, and another centered at position  $(x, z) = (6, 11.4)$ , which measures the overlap between the diffracted field and the fundamental mode of the optical fiber. This second monitor is placed at an angle of  $20.9^\circ$  (the optimized coupling angle); its center is at a distance of  $6 \mu\text{m}$  from the grating center axis (due to the higher coupling angle) and at a distance of  $4.4 \mu\text{m}$ , measured from the grating coupler’s starting point. The waveguide is  $7 \mu\text{m}$  long.

In Figure 33, it is noticeable that the diffracted electric field shows several diffraction angles, resulting in a poor coupling efficiency, Figure 34.



**Figure 33.** Representation on the XZ plane of the diffracted electric field: the green (transparent filled) horizontal bar is the waveguide’s monitor and the green (transparent filled) slanted bar is the diffracted field monitor. Image obtained from 2D-FDTD.



**Figure 34.** Efficiency test: from top to bottom, the blue line represents the overlap value (with the quasi-TE<sub>00</sub> waveguide's mode) measured at 1 μm before the waveguide's end. The red and green lines represent the power and overlap (with the TE<sub>00</sub> mode of a simulated SMF-28 optical fiber) of the diffraction field, respectively, measured by a monitor whose center is at a distance of 6 μm from the grating's center.

Looking at the results presented in Figure 34, it is possible to assert that the waveguide-monitor overlap (with the launch field, or fundamental quasi-TE<sub>00</sub> mode), remains nearly constant throughout the integration time. The coupling-efficiency value achieved at the time of convergence is slightly under −12.8 dB (~5.2%). The overlap with the diffracted field (i.e., the percentage of power of the diffracted field passing through the monitor which is converted into the SMF-28 fundamental mode) is about 33.3%. The ratio of diffracted power reaching the monitor to the waveguide's power is slightly over 8.1 dB (~15.7%).

## 5. Discussion

To summarize, all results obtained in this work are presented and compared with state-of-the-art (SoA) and high-efficiency grating-coupler designs, Table 4. All grating couplers presented herein were projected to operate in the optical C-band (1530–1565 nm), and the coupling efficiency of the proposed designs is given for a wavelength of 1550 nm. In Table 4, GC stands for grating coupler, N/A means information is not available, BW corresponds to bandwidth and DBR is distributed Bragg reflector.

Despite having inferior performance when compared to the state of the art in grating-coupler designs, the proposed GCs have the possibility of being manufactured at a lower cost. The devices, made of a-Si:H, can be deposited at low temperature by Plasma-Enhanced Chemical Vapor Deposition (PECVD), over a substrate of silica. Since silica is one of the main constituents of glass, there is also the possibility of adapting the devices to be fabricated over a substrate of this material.

Due to the fully etched grating design, all devices can be fabricated using one or two lithographic masks, making misalignment less of an issue. Designs requiring three lithographic processes are more affected by misalignment issues, and a study about the impact of misalignment must be conducted [40]. Using the method covered in this work, no silicon wafers are required, and the devices work without any bottom silicon layer(s) or upper cladding. Without a bottom reflector layer, the proposed grating couplers were designed to shape the diffracted field to enhance the overlap with the coupling profile of the single-mode optical fiber. The fill factor (and period, in apodized designs) was engineered to achieve the intended diffracted-field distribution. All presented designs can be patterned by photolithography [68], requiring either deep-UV (DUV) or extreme-UV (EUV) lithographic processes.



**Table 4.** Comparison between the devices designed in this work and several high-performance state-of-the-art grating couplers. The devices operate in the optical C-band (1530–1565 nm).

GC Design	GC Material	GC Feature Size *	Coupling Efficiency	−1 dB Bandwidth	Bottom Reflector	Required Masks **	Reference
Silicon Nitride Top Layer	Si	266 nm	−1.7 dB	64 nm	No	3	[67]
Fully Etched Apodized	Si	100 nm	−0.6 dB	(71 nm, −3 dB)	Yes, Al layer	1	[57]
Shift-pattern Overlay	Si/Poly-Si	171 nm	−0.9 dB	35 nm	No	3	[39]
Dual-level GC	Si	60 nm	−0.8 dB	31.3 nm	No	3	[40]
Bilayer GC	Si <sub>3</sub> N <sub>4</sub>	N/A	−1.0 dB	117 nm	Yes, DBR	2	[60]
Dual-level GC	Si <sub>3</sub> N <sub>4</sub> /Si	200 nm	−1.3 dB	80 nm	Yes, GC	2	[52]
Chirped GC	Si	26 nm	−0.1 dB	(35 nm, −3 dB)	Yes, DBR	2	[54]
Multilayer Bottom Reflector	SiN <sub>x</sub>	86 nm	−1.8 dB	52.5 nm	Yes, DBR	1	[53]
Bilayer GC	SiN <sub>x</sub>	N/A	−2.7 dB	47.9 nm	No	2	[42]
Compact Focusing GC	Si	315 nm	−7.0 dB ***	41 nm	No	2	[66]
Uniform GC	Si	270 nm	−5.8 dB	(48.5 nm, −3 dB)	No	2	[65]
Segmented GC	Si	190 nm	−2.9 dB	(71.4 nm, −3 dB)	No	2	[65]
Non-optimized GC	a-Si:H	510 nm	−9.7 dB	N/A	No	1	This work
Fill factor > 50%	a-Si:H	120 nm	−4.3 dB	25 nm	No	1	This work
Linear R.I. Variation	a-Si:H	60 nm	−3.1 dB	26 nm	No	1	This work
Quadratic R.I. Variation	a-Si:H	60 nm	−2.8 dB	25 nm	No	1	This work
Overlapped Micrometric	a-Si:H	750 nm	−7.5 dB	22 nm	No	2	This work
Random Distribution	a-Si:H	650 nm	−12.8 dB	N/A	No	2	This work

\* Lithographic mask minimum feature size. \*\* or etch steps. \*\*\* Result obtained for a wavelength of 1575 nm.

Regarding the performance of the presented devices, it is clear that apodization plays a significant role in performance enhancement, at the cost of requiring more expensive lithographic procedures, such as DUV, EUV or E-beam. The results suggest that quadratic refractive-index variation might present a small improvement in coupling efficiency when compared with linear variation. By optimizing the fill factor or opting for a grating made from the overlap of two lithographic masks, it is possible to make a compromise between performance and manufacturing costs. We were unable to achieve acceptable results from pseudo-random distributions; nevertheless, these could be the basis for more sophisticated optimization, such as, for example, genetic algorithms, by creating seed samples. In addition, it might be necessary to decrease the feature size of such grating designs, to improve efficiency.

## 6. Conclusions

Five different grating couplers were successfully designed and simulated, achieving coupling efficiencies between −12.8 dB and −2.8 dB. These results were obtained considering a hydrogenated amorphous silicon (a-Si:H) over silica platform, without any silicon bottom layer or bottom reflector. The technology is compatible with low-temperature PECVD and HWCVD, significantly reducing production costs.

Since the designs are fully etched, only one or two lithographic masks are required, reducing the dependency on mask alignment.

The simulations were conducted using the Finite-Difference Time-Domain (FDTD) analysis and Finite Element Method (FEM). The optical-fiber model was based on a commercial type (SMF-28), typically used in near-infrared communications, being characterized by typical design parameters.

The grating coupler featuring quadratic variation of the effective refractive index achieved the highest efficiency of the devices studied in this work (−2.8 dB), opening

the way for studies of different arithmetic functions or combination of functions (e.g., polynomial, exponential, logarithmic, hyperbolic) in the variation of the refractive index through the grating's length.

The superposition of two masks to produce a different grating coupler from two gratings is also a possible method, which allows a more complex design, with smaller feature sizes to be fabricated using photolithography techniques, thus avoiding electron beam (e-Beam).

## 7. Future Work and Improvement

From our perspective, this work contributes to open ways to future research in the optoelectronics and photonics fields, in which several paths undeniably still need to be explored. Since the analysis covered in this work was performed using a bidimensional model, further analysis should be conducted using a tridimensional (3D) model. With one more dimension, simulation of adiabatic tapers for grating coupler (GC) connection and focusing grating designs can be studied, and its geometry optimized. Following the tridimensional analysis, it is time to proceed to the production of real-world prototypes, followed by characterization and testing.

Despite the increase in costs and complexity, the inclusion of one or more bottom layers of silicon (arranged as a DBR stack) and under-cladding thickness adjustment could provide possibilities for further improvements in performance, opening the way for a performance comparison analysis of such designs, with the same GC devices featuring metal bottom reflectors.

It is also important to study the impacts of the arithmetic functions used in the apodization of the grating couplers to tailor the shape, distribution and length of the diffracted field, and to analyze its importance in matching the optical fiber's field profile and mean field diameter (MFD). In addition, it is relevant to study the effects on the diffracted field, of the slope and initial fill factor in linear apodization and of the coefficients of polynomial functions, like the quadratic distribution (where concavity can play a significant role).

Of course, there is also space for the utilization of machine learning (e.g., neural networks) and other artificial intelligence techniques for the optimization of apodization parameters. The study of arithmetic functions alone or combined opens an interesting perspective in the optimization of diffraction grating couplers.

**Author Contributions:** Conceptualization, D.A., J.C. and A.F.; methodology, D.A., J.C. and A.F.; software, D.A. and P.L.; validation, J.C. and A.F.; writing—original draft preparation, D.A.; writing—review and editing, J.C., A.F. and M.V.; supervision, J.C., A.F. and M.V.; project administration, A.F.; funding acquisition, A.F. and M.V. All authors have read and agreed to the published version of the manuscript.

**Funding:** This research was supported by Portuguese national funds provided by FCT—Fundação para a Ciência e a Tecnologia, through grant SFRH/BD/07792/2021 (<https://doi.org/10.54499/2021.07792.BD>) and funded by the Portuguese FCT program, Center of Technology and Systems (CTS) UIDB/00066/2020/UIDP/00066/2020 and by FCT project ASER-META 2022.07694.PTDC.

**Institutional Review Board Statement:** Not applicable.

**Informed Consent Statement:** Not applicable.

**Data Availability Statement:** The datasets presented in this article are not readily available because the data are part of an ongoing study. Requests to access the datasets should be directed to [afantoni@deetc.isel.ipl.pt](mailto:afantoni@deetc.isel.ipl.pt) or [dalmeida@deetc.isel.ipl.pt](mailto:dalmeida@deetc.isel.ipl.pt).

**Conflicts of Interest:** The authors declare no conflicts of interest.

## Appendix A

The pseudo-code algorithm employed to obtain the grating period ( $\Lambda$ ) and fill factor ( $F$ ), for each segment of grating coupler featuring linear refractive-index variation along its length, is presented herein:

```

gratingLength = 20;//micrometer
delta_neff = -0.3;//1/micrometer
delta_z = 10.4;//micrometer (SMF-28 MFD)
m = delta_neff/delta_z;
neff0 = 2.65348;//F = 0.9
theta = deg2rad(15);//15 deg
ne = 1;//air
nu = 2.8372;//220 nm aSiH strip waveguide (2D-FEM)
lambda0 = 1.55;//micrometer

z = 0;
neff = 0;
GP = 0;
F = 0;
num = 1;

while (z < gratingLength){
neff = m*z + neff0;
GP = lambda0/(neff - sin(theta));//GP = Grating period (micrometer)
F = (neff - ne)/(nu - ne);

print("GP"+num2str(num)+" = "+num2str(GP,3)+"\r\n");//round to nm
print("F"+num2str(num)+" = "+num2str(F,3) + "\r\n");//round to nm

z = z + GP;
num = num + 1;
}

```

## Appendix B

The pseudo-code used to obtain the period ( $\Lambda$ ) and fill factor ( $F$ ), for each segment of the grating coupler, featuring quadratic variation of the refractive index through its length, is presented below:

```

a = 0.001;
b = -0.039;
gratingLength = 20;//micrometer
neff0 = 2.65348;//F = 0.9
theta = deg2rad(15);//15 deg
ne = 1;//air
nu = 2.8372;//220 nm aSiH strip waveguide (2D-FEM)
lambda0 = 1.55;//micrometer

z = 0;
neff = 0;

```

```

GP = 0;
F = 0;
num = 1;

while (z < gratingLength){
neff = a*z^2 + b*z + neff0;
GP = lambda0/(neff - sin(theta)); //GP = Grating period (micrometer)
F = (neff - ne)/(nu - ne);

print(["GP"+num2str(num)+" = "+num2str(GP,3)+"\r\n"]); //round to
nm
print(["F"+num2str(num)+" = "+num2str(F,3)+"\r\n"]); //round to nm

z = z + GP;
num = num + 1;
}

```

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