



Article Improved Single-Crystal Growth of Sr₂RuO₄

Jake S. Bobowski ^{1,2,*}, Naoki Kikugawa ^{3,*}, Takuto Miyoshi ¹, Haruki Suwa ¹, Han-shu Xu ^{1,4}, Shingo Yonezawa ¹, Dmitry A. Sokolov ⁵, Andrew P. Mackenzie ⁵ and Yoshiteru Maeno ^{1,*}

- ¹ Department of Physics, Kyoto University, Kyoto 606-8502, Japan; miyoshi.takuto.82x@st.kyoto-u.ac.jp (T.M.); suwa.haruki.28w@st.kyoto-u.ac.jp (H.S.); xu.shu.63r@st.kyoto-u.ac.jp (H.-s.X.); yonezawa@scphys.kyoto-u.ac.jp (S.Y.)
- ² Department of Physics, University of British Colombia, Kelowna, BC V1V 1V7, Canada
- ³ National Institute for Materials Science, Tsukuba 305-0003, Japan
- ⁴ Hefei National Laboratory for Physical Sciences at the Microscale, University of Science and Technology of China, Hefei 230026, China
- ⁵ Max Plank Institute for Chemical Physics of Solids, D-01187 Dresden, Germany; Dmitry.Sokolov@cpfs.mpg.de (D.A.S.); Andy.Mackenzie@cpfs.mpg.de (A.P.M.)
- * Correspondence: jake.bobowski@ubc.ca (J.S.B.); KIKUGAWA.Naoki@nims.go.jp (N.K.); maeno@scphys.kyoto-u.ac.jp (Y.M.)

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Abstract: High-quality single crystals are essentially needed for the investigation of the novel bulk properties of unconventional superconductors. The availability of such crystals grown by the floating-zone method has helped to unveil the unconventional superconductivity of the layered perovskite Sr_2RuO_4 , which is considered as a strong candidate of a topological spin-triplet superconductor. Yet, recent progress of investigations urges further efforts to obtain ultimately high-quality crystalline samples. In this paper, we focus on the method of preparation of feed rods for the floating-zone melting and report on the improvements of the crystal growth. We present details of the improved methods used to obtain crystals with superconducting transition temperatures T_c that are consistently as high as 1.4 K, as well as the properties of these crystals.

Keywords: Sr₂RuO₄; superconductivity; single-crystal growth; floating zone; feed rod

1. Introduction

The availability of high-quality single crystals is essential for the full clarification of the novel bulk properties of quantum materials, especially of unconventional superconductors. The layered perovskite superconductor Sr_2RuO_4 [1] is a typical example. This superconductor has attracted much attention over the last twenty years as a strong candidate of a spin-triplet superconductor [2–6], as well as of a bulk topological superconductor [7,8]. Reflecting its unconventional pairing, its superconductivity is sensitively suppressed by even non-magnetic impurities [9,10]. Thus, it is challenging to consistently obtain a large homogeneous crystal of Sr_2RuO_4 with T_c close to the intrinsic value of 1.50 K, corresponding to the mean free path exceeding 1 µm. Its superconductivity is completely suppressed when the mean-free-path becomes comparable to the superconducting coherence length (\approx 70 nm), corresponding to the impurity level of c.a. 1500 ppm, where the impurities act as strong scattering centers. As another characteristic of unconventional superconductivity, the quasiparticle density of states readily emerges within the superconducting gap even with a small amount of impurities and defects [11]; the resulting large residual density of states often makes the determination of the intrinsic gap anisotropy a challenging issue [12–15]. Furthermore, some unusual superconducting phenomena occur only in very high-purity crystals: the superconducting transition with a magnetic field applied parallel to the RuO_2 plane becomes first order at low temperatures [16–18]. Despite the key experimental results supporting spin-triplet pairing [19–23], such a first-order transition is difficult to explain within the context of spin-triplet superconductivity. This first-order transition becomes second order when the T_c is suppressed below 1.45 K, which corresponds to an impurity

order to deepen our knowledge of the superconducting state of Sr_2RuO_4 . Recent innovations in the design of uniaxial-stress cells enables the T_c of Sr_2RuO_4 to be enhanced up to 3.5 K when the stress is along the crystalline [100] direction [24–27]. The origin of this enhancement of T_c is attributed to the Fermi-level crossing of the van-Hove singularity in one of the three quasi-two-dimensional Fermi surfaces. It is hoped that detailed investigations of this phenomenon will lead to the clarification of the superconducting symmetry and mechanism in Sr_2RuO_4 . The T_c enhancement has actually been known for many years in the eutectic crystals of Sr_2RuO_4 with micron-size metallic Ru platelets, which introduce strong strains in Sr_2RuO_4 near the interfaces [28–30]. In order to investigate the strain-induced superconducting phase, high-quality Ru-inclusion free single crystals of Sr_2RuO_4 with a specific in-plane crystalline direction are in demand.

level of ~50 ppm [10]. Thus, pristine samples with impurity levels less than 10 ppm are required in

In this study, we examine how we can further improve the quality of crystals of Sr_2RuO_4 grown by the floating-zone method with an infrared image furnace. The floating-zone technique has been used to produce high-quality crystals of relatively large size useful for most experimental purposes including inelastic neutron scattering. Because the technique is essentially crucible-free, it can be used to achieve the minimum possible impurity levels. The floating-zone technique has been used for the successful crystal growth of Sr_2RuO_4 [31], as well as other ruthenates [32–37]. Previous reports mainly describe optimization of the atomic compositions of the feed rod and various parameters of the final growth process. On the other hand, there are a number of processes (from powder grinding to feed-rod sintering) involved prior to the zone-melting growth process. These processes turn out to be of critical importance to minimize impurity contamination. In this study, we focus on the process of feed-rod preparation in order to further improve the quality of Sr_2RuO_4 crystals. Although the reexaminations of our previous processes were performed independently in Kyoto and Dresden, we achieved consistent results and reached a consistent set of conclusions. The improved feed-rod preparation processes have allowed for the efficient production of high-quality crystals which will help deepen our understanding of the superconducting state of Sr_2RuO_4 .

2. Experimental

We first describe our standard procedure of feed-rod preparation [31], and identify possible issues. After grinding the starting materials of SrCO₃ and RuO₂ with the molar ratio of Sr:Ru = 2:*n* with n = 1.15, the powder was pressed into pellets using stainless steel molds. The pellets were typically pre-sintered at 1150 °C for 24 h. Next, the pellets were reground and the powder was packed into a balloon to shape the feed rod. For this balloon, we typically use latex tubing with a diameter of 6 mm and a thickness of 0.2 mm. The powder-filled balloon was then pressed under hydrostatic pressure (40 MPa for 5 min) and the resulting rod was sintered at 1420 °C for 2 h. Finally, the sintered rod was suspended in the floating-zone furnace. In this method, many steps are required before arriving at the final feed rod which allows for a number of opportunities for contamination. In particular, the surface of the pellets show a greenish color after pressing in the stainless-steel molds. This contamination is also evident from a black material on the surface of the pellets when pelletizing insulating white oxide powders instead of black RuO₂. Thus, when pellets are really needed, we cover the entire interior of the stainless-steel molds with Teflon sheets and discs. In the newly-developed processes described below, we avoid using stainless-steel molds altogether in order to establish procedures that leave less opportunity for contamination of the feed rod.

We have examined the following Methods (A)–(D), in all of which we avoid using stainless-steel molds. Procedure (A) is the simplest and perhaps the least prone to contamination. Methods (B) through (D) were developed to increase the hardness of the sintered feed rod and at the same time to reduce evaporation during the melt growth. For example, in Method (A) the handling of the powders is

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minimized which should significantly reduce the potential for contamination. However, because there is no pre-sintering, the resulting feed rod is very delicate and contains carbon from the SrCO₃ starting powder. Method (D), on the other hand, requires the most handling of the powder but results in a relatively robust polycrystalline feed rod. Through trials that have been done to date, the intermediate methods have been preferred (Method (B) at Dresden and Method (C) at Kyoto). These two methods offer the best compromise between reducing the risk of contamination of the powders and achieving polycrystalline Sr₂Ru_nO_{2n+2} feed rods that are not too fragile and can, therefore, be suspended in the image furnace with relative ease.

- (A) The ground powder is packed into a balloon, then pressed, and suspended in the floating-zone furnace. Here there is no sintering process and the feed rod contains carbon from pre-reacted SrCO₃.
- (B) The ground powder is packed into a balloon, then pressed, and sintered (1150 °C for 24 h in Kyoto, 1000 °C for 2 h in Dresden). The sintered rod is then suspended in the floating-zone furnace.
- (C) The ground powder is packed into a balloon, then pressed, and sintered at 1150 °C for 24 h and then subsequently at 1420 °C for 2 h. The sintered rod is then suspended in the floating-zone furnace.
- (D) The ground powder is packed into a balloon, then pressed, and pre-sintered at 1150 °C for 24 h. The pre-sintered rod is reground and the powder is again packed into another balloon, then pressed, and sintered at 1420 °C for 2 h. The sintered rod is then suspended in the floating-zone furnace.

All four procedures were examined in Kyoto, with emphasis on (C); Method (B) was adopted in Dresden. From these feed rods, Sr_2RuO_4 crystals were gown by the floating-zone method as described by Mao et al. [31]. Infrared image furnaces with double-elliptical mirrors (Canon Machinery, model SC-E15HD in Kyoto, and model SC-MDH in Dresden) were used. We note that in Dresden a part of the feed rod was used as a seed during the floating-zone growth, whereas oriented single crystals or a part of the feed rod were used as seeds in Kyoto.

We now describe the procedure used to prepare the feed rods in detail, since this is the part we particularly focus on in this study. First, two parts $SrCO_3$ are ground with *n* parts RuO_2 in a dry nitrogen atmosphere using a mortar and pestle for at least one hour. We used 4N+ SrCO₃ containing only 5 ppm Ba. Since $SrCO_3$ readily absorbs moisture from the air, it is first heated to 500 °C for one hour to remove adsorbed water and then weighed before being allowed to cool back to room temperature. We used 3N RuO₂ typically containing a few hundred ppm Cl and ~50 ppm Si as main impurities according to the Glow Discharge Mass Spectroscopy (GDMS) analysis. The metal impurities, such as Fe, Cu, and Zn, in the RuO_2 are less than 20 ppm each. The achieved T_c of as high as 1.495 K suggests that some of these impurities do not enter into the crystal structure of Sr₂RuO₄ during its growth from the melt. Even though the analysis shows less batch dependence, we experience some batch dependence in sintering the rods; consequently, adjustment of the feed speed during the floating-zone growth is necessary. Before packing the ground powder into a latex tubing balloon, both the inner and outer surfaces of the balloon were thoroughly cleaned. The surfaces of the as-purchased latex tubing are coated with a fine TiO₂ powder which prevents the balloon surfaces from sticking to itself. This powder coating must be completely removed in order to avoid unnecessary contamination of the feed rods.

Before the powder is introduced inside, a length of balloon, with both ends open, is slide over a glass rod and its outside surface is cleaned using isopropanol or ethanol and a lint-free wipe. The clean surface of the balloon is then coated with some of the ground $2\text{SrCO}_3 + n\text{RuO}_2$ powder. The balloon is then inverted on the glass rod such that the opposite side can be cleaned and coated. This entire process is repeated for a second time such that the inner and outer surfaces of the balloon are each cleaned and coated twice.

Next, one end of the balloon is tied shut with two knots separated by a few millimeters and then loaded with typically 5 g of 2SrCO₃ + nRuO₂ powder (Figure 1a). The powder is packed into the

bottom of the balloon so as to remove as much air as possible before sealing the opposite end of the balloon with a pair of knots. Once sealed, the powder is distributed evenly throughout the length of the balloon. To form a long narrow rod, the filled balloon is inserted into a tightly-rolled tube of paper. Thick paper is rolled onto a metal rod and then secured with cellophane tape before extracting the rod. The resulting tube of paper is several layers thick and has an inner diameter of 6 mm. A length of string is then tied between the pair of knots at one end of the balloon and used to pull the powder-filled balloon into the paper tube (Figure 1b,c).

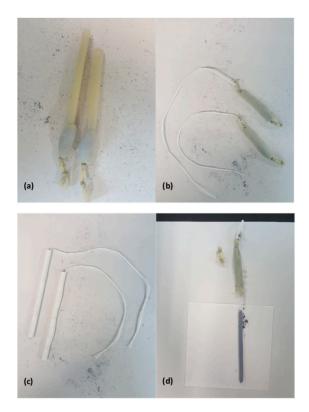


Figure 1. The process of making a feed rod. (**a**) Balloons packed with ground $SrCO_3$ and RuO_2 powder. (**b**) The balloons are sealed and the powder inside is distributed as evenly as possible. A string is tied between the two knots at one end of the balloon. (**c**) The string is used to stretch the balloon and pull it into a paper tube with a 6 mm inner diameter. (**d**) After applying hydrostatic pressure, the feed rod is removed from the balloon.

The paper tube containing the powder-filled balloon is then loaded into a stainless-steel rod which has been partially bored out from one end. The bore is next filled with water and a tight-fitting piston and hydraulic press are used to apply 40 MPa of hydrostatic pressure for five minutes. After removing and drying the paper tube, the cellophane tape is cut and the paper tube carefully uncoiled to expose the balloon. Fine scissors are used to cut one end of the balloon while one's remaining hand lightly pinches the balloon against the pressed-powder rod near the cut. After the cut is made, the pressure applied with the fingers can be slowly released in order to allow the rod to gently slide out of the balloon (Figure 1d). The resulting rod, if in one piece, will be 6 mm in diameter and up to 10 cm long. Typically, only feed rods that are at least 5 cm in length are suitable for floating-zone growth. The rod is very delicate and should be handled with great care.

In Method (A), a wire is attached to this feed rod and one then proceeds with the floating-zone growth. A utility knife can be used to carve a narrow groove near the end of the feed rod. A length of fine chromel wire (90 percent nickel and 10 percent chromium, ϕ 0.2 mm) is wrapped around the groove in order to form a loop that can be used to suspend the feed rod in the floating-zone furnace as shown in Figure 2. In Methods (B)–(D), the feed rod is first sintered before the chromel wire is attached.

For sintering, the feed rods are first transferred onto an alumina boat that has a bedding of pre-sintered $Sr_2Ru_nO_{2n+2}$ powder. This bedding is used to prevent direct contact between the feed rod and alumina boat, which can be a source of unwanted contamination. In Method (B) the rods are sintered for 24 h at 1150 °C (or 2 h at 1000 °C) before the chromel wire is attached. In (C) the rods are sintered at 1150 °C for 24 h followed by 1420 °C for an additional 2 h, after which the chromel wire is attached to the feed rod. Finally, in Method (D) the feed rods are first pre-sintered at 1150 °C for 24 h. These rods were then reground into powder and new feed rods were formed using the procedures described above. This second set of feed rods were then sintered at 1420 °C for 2 h before adding the chromel wire and then suspended in the floating-zone furnace. Weight-loss measurements during sintering processes (B)–(D) show that the evaporation of the excess RuO_2 is less than $\Delta n = 0.01$. The powder X-ray diffraction analysis indicates that the rods sintered at and above 1150 °C contain a slight amount of $Sr_3Ru_2O_7$. We note that the high-temperature sinter has the advantage that the resulting feed rods are less fragile and, therefore, easier to handle and manipulate. As a seed material fixed in a seed holder at the bottom (Figure 2), we use either a part of the polycrystalline feed rod or a crystal of Sr_2RuO_4 .



Figure 2. Photo of a polycrystalline feed rod and seed ready for the floating-zone growth. They are placed at the center of the two elliptical mirrors of the infrared mirror furnace at Kyoto. In this example, a polycrystalline seed is used, rather than a single crystal.

During the floating-zone growth, some RuO₂ is lost to evaporation via the reaction:

$$Sr_2Ru_nO_{2n+2} \to Sr_2Ru_{n'}O_{2n'+2} + (n-n')RuO_2.$$
 (1)

Therefore, it is necessary to start with n > 1 to produce single crystals with $n' \cong 1$. In Appendix A we describe how the value of n' can be estimated from the mass of the final single crystal after the floating-zone growth and the change in mass of the feed rod. The value of n' is a good indicator to monitor proper growth, and for n = 1.15 we obtain n' that are slightly less than 1. So far, we have not obtained any systematic relation between n' and T_c as long as n = 1.15, perhaps because the precision in determining n' is not sufficiently high. Figure 3a shows photographs of a polycrystalline feed rod prepared following Method (C) after its tip has been melted in the image furnace. The melted tip of the feed rod is suspended above a single-crystal seed. Figure 3b is an image taken during stable growth of single-crystal Sr₂RuO₄ by the floating-zone technique. We use essentially the same growth speed is 43 mm/h in a gas mixture of O₂ (15%) and Ar (85%) at 0.25 MPa. The feed and seed are rotated in opposite directions, each at 30 rpm. In Dresden, the feed speed is typically 28–30 mm/h, and the

crystal-growth speed is 45 mm/h in a gas mixture of O_2 (15% or 10%) and Ar (85% or 90%) at 0.35 MPa. The feed and seed are rotated in opposite directions typically at 30 rpm.

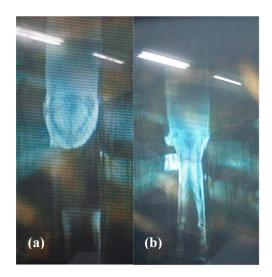


Figure 3. (a) Melted tip of the feed rod just before connecting it to the seed crystal below. The diameter of the feed rod is 6 mm. (b) The melted region between the feed rod and single crystal during stable floating-zone growth.

3. Results

Figure 4 shows a series of back-scattered Laue images from Sr_2RuO_4 samples along the [001] direction. All of these samples were grown using the floating-zone technique starting from feed rods prepared from Method (C). The Laue results confirm the good crystallinity of the samples. We note that high-quality samples from Method (C) have been obtained using both single crystal (C447–C453) and polycrystalline (C442 and C454) seeds. It is much preferred for some key experiments that a particular crystalline axis is along the single-crystalline rod axis. The results in Figure 4 show a tendency of the [110] orientation along the crystalline rod if polycrystalline seeds are used. This tendency was found independently in Dresden using polycrystalline seeds. On the other hand, by using single-crystalline seeds, crystals with the axes matching with those of the seeds can be produced; this facilitates the preparation of single crystals suited for uniaxial-stress experiments. It is also noted that these Laue photos indicate low mosaicity within the X-ray beam spot size of 1 mm; namely, inclusion of crystal domains with slightly different orientations is almost negligible. We took several additional Laue photos in different positions covering the area of $2 \times 2 \text{ mm}^2$ on the crystal piece C454B with the area of $3.2 \times 3.3 \text{ mm}^2$ and confirmed low mosaic spread in a wider part of the crystal.

At Kyoto, low-temperature AC susceptibility measurements are used to assess the quality of the as-grown Sr₂RuO₄ crystals. High-quality samples exhibit a sharp superconducting transition with T_c near 1.50 K. The AC susceptibility measurements are done using a commercial system (Quantum Design, Model PPMS) with an Adiabatic Demagnetization Refrigerator (ADR) option. The coil set and data acquisition system were designed and developed in Kyoto [38]. The dimensions of the samples measured are typically 4 mm × 3 mm × 2 mm where the *c*-axis direction is along the shortest dimension and the AC magnetic field from the primary coil is applied parallel to the *ab*-plane. Figure 5 shows the real and imaginary parts of the AC susceptibility of the sample (batch-sample number C447-B) whose feed rod was prepared following Method (C). For this sample, a single-crystal seed was used during the floating-zone growth. The χ' data show an onset T_c that is greater than 1.45 K and the width of the superconducting transition, as measured from the FWHM of χ'' , is less than 20 mK. As described earlier, we used as a seed either a crystal of Sr₂RuO₄ or a part of the polycrystalline feed rod. High-quality crystals are produced in both cases. However, the samples showing sharpest superconducting transitions with transition widths less than 30 mK were obtained

when using single-crystal seeds. Although most of the experimental work at Kyoto has focused on Method (C), we have also successfully grown high-quality Sr_2RuO_4 crystals using Methods (A), (B), and (D). Method (A) has the advantage that it involves the least handling of the powders prior to the floating-zone growth and therefore is the least susceptible to contamination. However, since there is no pre-sintering the feed rods are very fragile and can be easily broken when adding the wire required to suspend the rod in the image furnace. Furthermore, these feed rods still contain carbon at the start of the floating-zone growth. Method (D) has the disadvantage that it requires the most handling of the powder prior to the floating-zone growth and these feed rods take the most time to prepare.

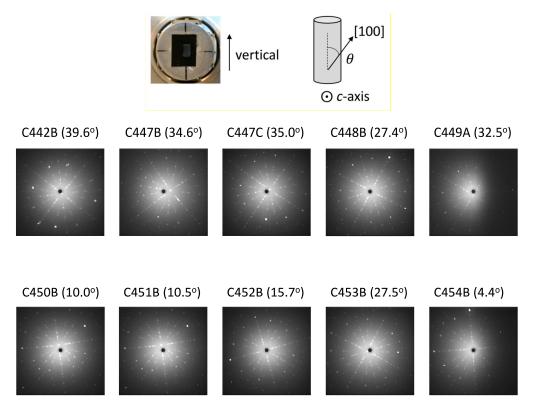


Figure 4. Laue photos of the Sr_2RuO_4 crystals grown at Kyoto using Method (C) for the feed-rod preparation. These are back-scattered Laue photos taken with X-ray beam nearly along the [001] direction. A beam collimator of 1-mm diameter was used. The values in the parentheses indicate the angle of the [100] direction with respect to the crystalline-rod axis. Polycrystalline seed was used for C442, resulting in the [110] direction within 6 degrees from the crystalline-rod axis. Single-crystalline seeds were used in all other cases except C454. The same single crystal (C339) was used as seeds for C447 and C448; C441 was used as seeds for both C450 and C451, resulting in the [100] axis close to the crystalline-rod axis. For C448, crystal was disconnected during the growth and the crystal does not preserve the orientation of the seed crystal. Otherwise, the seed crystal orientations are well preserved in the crystals grown.

A photograph of a crystal grown from a feed rod prepared by Method (B) at Dresden is shown in Figure 6a. Because we used a polycrystalline seed, rather than a single-crystalline seed, for this crystal, the growth at the first stage was necked. A back-scattered Laue-image from another batch of sample along the [001] direction is shown in Figure 6b, after cleaving the as-grown pieces along the basal plane. It is worth mentioning that crystals grown using polycrystalline seeds have a tendency of their [110] axis along the crystalline rod within 10 degrees. This tendency has also been observed at Kyoto. We note, however, that despite this tendency there are exceptions. Sample C454 was grown at Kyoto using a polycrystalline seed and its [100] axis was found to be misaligned with the crystal-growth axis by only 4.4°. The bottom-right image in Figure 4 shows the back-scattered Laue image obtained

from this sample. In order to characterize the bulk quality of the crystal prepared following Method (B) in Dresden, we measured the specific heat (C_p) as a function of temperature (T) below 2 K using a PPMS equipped with a ³He cooling system. Figure 6c shows the measured specific heat divided by temperature (C_p/T) as a function of temperature. We see a clear and sharp superconducting transition at 1.5 K, along with a relatively small residual density of states, deduced by extrapolating C_p/T to T = 0 K. These observations guarantee that the grown crystal is of high quality. We note that the typical sample size used for the C_p measurement is 3 mm \times 3 mm \times 1 mm. This size is sufficient for a wide variety of experiments.

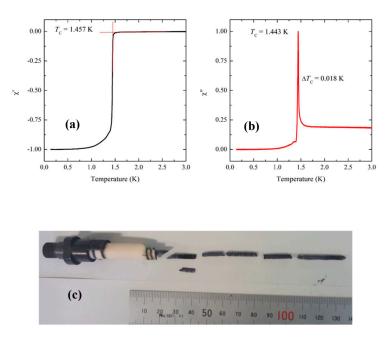


Figure 5. Chi-AC and the corresponding photo of the Kyoto crystal C447B (Method (C)). (**a**) Real and (**b**) imaginary parts of the AC susceptibility of a Sr_2RuO_4 crystal that was prepared following Method (C). The sample that was measured was cut from the crystal shown in (**c**).

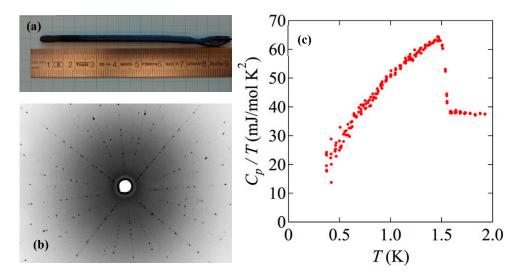


Figure 6. (a). A photograph of a crystal grown by Method (B) in Dresden. A polycrystalline rod was used for the seed. (b) A back-scattered Laue photograph taken with X-ray beam nearly along the [001] direction. A beam collimator of 0.5-mm diameter was used. As described in the main text, the crystalline [110] direction for this sample is close to the rod axis. (c) Specific heat divided by temperature (C_p/T) as a function of temperature below 2 K. We can see the clear transition at high- T_c with $T_c = 1.50$ K as well as the small residual C_p/T to T = 0 K.

4. Conclusions

We examined several feed-rod preparation procedures with the aim of avoiding impurity contamination to improve our conventional method. In particular, in all of the new procedures the feed-rod powder never comes into direct contact with any metallic surfaces. Method (B) and (C) offer the best compromise between minimizing the potential for contamination of the powders during the feed-rod preparation and producing robust feed rods of polycrystalline Sr₂Ru_nO_{2n+2}. We demonstrated that the process of making pellets is unnecessary, and the crystals obtained from the feed rods prepared by the new procedures were of high quality. At Dresden, the success rate for obtaining crystals with $T_c > 1.45$ K was 57% among 23 different batch crystals, and at Kyoto superconducting transition widths less than 20 mK as measured by AC susceptibility were obtained. The new procedures have the additional advantage that the time required to prepare the feed rods has been substantially reduced.

For most experiments using single crystals, it is desirable to control the direction of the crystalline axes with respect to the crystalline-rod direction. For the investigation of the uniaxial-strain effects, for example, it is desirable to have samples with the [100] direction parallel to the long axis of the crystal. In this study, we confirmed that the seed-crystal orientation is well preserved in the crystal growth and found a tendency of [110] orientation when polycrystalline feed rods are used. The improvements and additional insights acquired in this study have helped advanced us towards our goal of obtaining the "ultimate" crystals of Sr_2RuO_4 to be used for the full clarification of its superconducting state.

Author Contributions: J.S.B. and Y.M. planned the project at Kyoto, whereas N.K., D.A.S., and A.P.M. planned it independently at Dresden. J.S.B., T.M., H.S., S.Y. grew the crystals and characterized them at Kyoto, and N.K. and D.A.S. at Dresden. The manuscript was written mainly by Y.M., J.S.B., and N.K. with the input from all the authors.

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Conflicts of Interest: The authors declare no conflict of interest.

Appendix A. Estimating n' in Sr₂Ru_{n'}O_{2n'+2}</sub>

The floating-zone growth converts a polycrystalline feed rod with the composition $Sr_2Ru_nO_{2n+2}$ into a single crystal with the composition $Sr_2Ru_{n'}O_{2n'+2}$. This appendix describes how the value of n' can be estimated from set of masses that can be easily measured. The estimated value of n' in turn helps one to optimize the value of the starting composition n. Before the floating-zone growth, one can measure the combined mass of the seed holder and the seed material (m_{s1}) and the mass of the polycrystalline feed rod plus the wire that is used to suspend it in the floating-zone furnace (m_{f1}). After the crystal growth, the total mass of the single crystal, seed and seed holder (m_{s2}), as well as the mass of the remaining feed rod plus suspending wire (m_{f2}) can both be measured. The difference between the final and initial seed masses results in the mass of the crystal produced during the floating-zone growth:

$$\Delta m_{\rm s} = m_{\rm s2} - m_{\rm s1} = N M_{n'},\tag{A1}$$

where $M_{n'}$ is the molar mass of the formula unit $\text{Sr}_2\text{Ru}_{n'}\text{O}_{2n'+2}$ and N is the molar number. The difference in the initial and final feed rod masses specifies the total mass of polycrystalline $\text{Sr}_2\text{Ru}_n\text{O}_{2n+2}$ that was converted into single crystal $\text{Sr}_2\text{Ru}_n'\text{O}_{2n'+2}$:

$$\Delta m_{\rm f} = m_{\rm f2} - m_{\rm f1} = -NM_n,\tag{A2}$$

where M_n is the molar mass of the formula unit Sr₂Ru_nO_{2n+2}. Equation (A2) can be used to determine N since n is known from the initial masses of the 2SrCO₃ and nRuO₂ powders (typically, n = 1.15). Making use of the fact that $M_n - M_{214} = (n - 1)M_{RuO_2}$ allows one to write:

$$N = \frac{-\Delta m_{\rm f}}{(n-1)M_{\rm RuO_2} + M_{214}},\tag{A3}$$

where M_{214} and M_{RuO_2} are the molar masses of Sr₂RuO₄ and RuO₂, respectively. During the single-crystal growth, the mass of RuO₂ that is lost to evaporation is given by:

$$-\Delta m_{\rm f} - \Delta m_{\rm s} = N(n - n')M_{\rm RuO_2}.$$
(A4)

Eliminating *N* from Equations (A3) and (A4) and solving for n' leads to:

$$n' = n \left| \frac{\Delta m_{\rm s}}{\Delta m_{\rm f}} \right| + \left(\frac{M_{214}}{M_{\rm RuO_2}} - 1 \right) \left(\left| \frac{\Delta m_{\rm s}}{\Delta m_{\rm f}} \right| - 1 \right),\tag{A5}$$

where the ratio $M_{214}/M_{RuO_2} = 2.5574$.

References

- 1. Maeno, Y.; Hashimoto, H.; Yoshida, K.; Nishizaki, S.; Fujita, T.; Bednorz, J.G.; Lichtenberg, F. Superconductivity in a layered perovskite without copper. *Nature* **1994**, *372*, 532–534. [CrossRef]
- Mackenzie, A.P.; Maeno, Y. The superconductivity of Sr₂RuO₄ and the physics of spin-triplet pairing. *Rev. Mod. Phys.* 2003, 75, 657–712. [CrossRef]
- 3. Maeno, Y.; Kittaka, S.; Nomura, T.; Yonezawa, S.; Ishida, K. Evaluation of Spin-Triplet Superconductivity in Sr₂RuO₄. *Phys. Soc. Jpn.* **2012**, *81*, 11009. [CrossRef]
- 4. Liu, Y.; Mao, Z.-Q. Unconventional superconductivity in Sr₂RuO₄. *Phys. C* 2015, 514, 339–353. [CrossRef]
- 5. Kallin, C.; Berlinsky, J. Chiral superconductors. Rep. Prog. Phys. 2016, 79, 054502. [CrossRef] [PubMed]
- 6. Mackenzie, A.P.; Scaffidi, T.; Hicks, C.W.; Maeno, Y. Even odder after twenty-three years: the superconducting order parameter puzzle of Sr₂RuO₄. *npj Quantum Mater.* **2017**, *2*, 40. [CrossRef]
- Alicea, J. New directions in the pursuit of Majorana fermions in solid state systems. *Rep. Prog. Phys.* 2012, 75, 076501. [CrossRef]
- 8. Sato, M.; Ando, Y. Topological superconductors: A review. Rep. Prog. Phys. 2017, 80, 076501. [CrossRef]
- 9. Mackenzie, A.; Haselwimmer, R.; Tyler, A.; Lonzarich, G.; Mori, Y.; Nishizaki, S.; Maeno, Y. Extremely Strong Dependence of Superconductivity on Disorder in Sr₂RuO₄. *Phys. Rev. Lett.* **1998**, *80*, 161–164. [CrossRef]
- Kikugawa, N.; Mackenzie, A.P.; Maeno, Y. Effects of In-Plane Impurity Substitution in Sr₂RuO₄. J. Phys. Soc. Jpn. 2003, 72, 237–240. [CrossRef]
- Dodaro, J.F.; Wang, Z.; Kallin, C. Effects of deep superconducting gap minima and disorder on residual thermal transport in Sr₂RuO₄. *arxiv*. 2018. Available online: https://arxiv.org/abs/1810 (accessed on 28 December 2018). [CrossRef]
- 12. NishiZaki, S.; Maeno, Y.; Mao, Z.Q. Effect of Impurities on the Specific Heat of the Spin-Triplet Superconductor Sr₂RuO₄. *J. Low Temp. Phys.* **1999**, *117*, 1581–1585. [CrossRef]
- Deguchi, K.; Mao, Z.Q.; Yaguchi, H.; Maeno, Y. Gap Structure of the Spin-Triplet Superconductor Sr₂RuO₄ Determined from the Field-Orientation Dependence of the Specific Heat. *Phys. Rev. Lett.* 2004, *92*, 047002. [CrossRef] [PubMed]
- Hassinger, E.; Bourgeois-Hope, P.; Taniguchi, H.; René de Cotret, S.; Grissonnanche, G.; Anwar, M.S.; Maeno, Y.; Doiron-Leyraud, N.; Taillefer, L. Vertical Line Nodes in the Superconducting Gap Structure of Sr₂RuO₄. *Phys. Rev. X* 2017, *7*, 011032. [CrossRef]
- Kittaka, S.; Nakamura, S.; Sakakibara, T.; Kikugaa, N.; Terashima, T.; Uji, S.; Sokolov, D.A.; Mackenzie, A.P.; Irie, K.; Tsutsumi, Y.; Suzuki, K.; Machida, K. Searching for Gap Zeros in Sr₂RuO₄ via Field-Angle-Dependent Specific-Heat Measurement. J. Phys. Soc. Jpn. 2018, 87, 093703. [CrossRef]
- Yonezawa, S.; Kajikawa, T.; Maeno, Y. First-Order Superconducting Transition of Sr₂RuO₄. *Phys. Rev. Lett.* 2013, 110, 77003. [CrossRef] [PubMed]
- Kittaka, S.; Kasahara, A.; Sakakibara, T.; Shibata, D.; Yonezawa, S.; Maeno, Y.; Tenya, K.; Machida, K. Sharp magnetization jump at the first-order superconducting transition in Sr₂RuO₄. *Phys. Rev. B* 2014, *90*, 220502. [CrossRef]

- Kikugawa, N.; Terashima, T.; Uji, S.; Sugii, K.; Maeno, Y.; Graf, D.; Baumbach, R.; Brooks, J. Superconducting subphase in the layered perovskite ruthenate Sr₂RuO₄ in a parallel magnetic field. *Phys. Rev. B* 2016, *93*, 184513. [CrossRef]
- 19. Ishida, K.; Mukuda, H.; Kitaoka, Y.; Asayama, K.; Mao, Z.Q.; Mori, Y.; Maeno, Y. Spin-triplet superconductivity in Sr₂RuO₄ identified by ¹⁷O Knight shift. *Nature* **1998**, *396*, 658–660. [CrossRef]
- 20. Ishida, K.; Mukuda, H.; Kitaoka, Y.; Mao, Z.; Mori, Y.; Maeno, Y. Anisotropic Superconducting Gap in the Spin-Triplet Superconductor Sr₂RuO₄: Evidence from a Ru-NQR Study. *Phys. Rev. Lett.* **2000**, *84*, 5387–5390. [CrossRef]
- 21. Murakawa, H.; Ishida, K.; Kitagawa, K.; Mao, Z.; Maeno, Y. Measurement of the ¹⁰¹Ru-Knight Shift of Superconducting Sr₂RuO₄ in a Parallel Magnetic Field. *Phys. Rev. Lett.* **2004**, *93*, 167004. [CrossRef] [PubMed]
- 22. Duffy, J.; Hayden, S.; Maeno, Y.; Mao, Z.; Kulda, J.; McIntyre, G. Polarized-Neutron Scattering Study of the Cooper-Pair Moment in Sr₂RuO₄. *Phys. Rev. Lett.* **2000**, *85*, 5412–5415. [CrossRef]
- 23. Ishida, K.; Manago, M.; Yamanaka, T.; Fukazawa, H.; Mao, Z.Q.; Maeno, Y.; Miyake, K. Spin polarization enhanced by spin-triplet pairing in Sr₂RuO₄ probed by NMR. *Phys. Rev. B* **2015**, *92*, 100502(R). [CrossRef]
- Hicks, C.W.; Brodsky, D.O.; Yelland, E.A.; Gibbs, A.S.; Bruin, J.A.N.; Barber, M.E.; Edkins, S.D.; Nishimura, K.; Yonezawa, S.; Maeno, Y.; et al. Strong Increase of T_c of Sr₂RuO₄ Under Both Tensile and Compressive Strain. *Science* 2014, 314, 283–285. [CrossRef] [PubMed]
- Steppke, A.; Zhao, L.; Barber, M.E.; Scaffidi, T.; Jerzembeck, F.; Osner, H.; Gibbs, A.S.; Maeno, Y.; Simon, S.H.; Mackenzie, A.P.; et al. Strong peak in *T*_c of Sr₂RuO₄ under uniaxial pressure. *Science* 2017, 355, eaaf9398. [CrossRef]
- 26. Barber, M.E.; Gibbs, A.S.; Maeno, Y.; Mackenzie, A.P.; Hicks, C.W. Resistivity in the Vicinity of a van Hove Singularity: Sr₂RuO₄ under Uniaxial Pressure. *Phys. Rev. Lett.* **2018**, *120*, 076602. [CrossRef] [PubMed]
- Luo, Y.; Guzman, Y.; Dioguardi, A.P.; Pustogow, A.; Thomas, S.M.; Ronning, F.; Kikugawa, N.; Sokolov, D.; Jerzembeck, F.; Mackenzie, A.P.; et al. Normal state ¹⁷O NMR studies of Sr₂RuO₄ under uniaxial stress. *arXiv.* 2018. Available online: https://arxiv.org/abs/1810.01209 (accessed on 2 October 2018).
- Maeno, Y.; Ando, T.; Mori, Y.; Ohmichi, E.; Ikeda, S.; NishiZaki, S.; Nakatsuji, S. Enhancement of Superconductivity of Sr₂RuO₄ to 3 K by Embedded Metallic Microdomains. *Phys. Rev. Lett.* 1998, *81*, 3765–3768. [CrossRef]
- 29. Ando, T.; Akima, T.; Mori, Y.; Maeno, Y. Upper Critical Fields of the 3-K Superconducting Phase of Sr₂RuO₄. *J. Phys. Soc. Jpn.* **1999**, *68*, 1651–1656. [CrossRef]
- 30. Yaguchi, H.; Wada, M.; Akima, T.; Maeno, Y.; Ishiguro, T. Interface superconductivity in the eutectic Sr₂RuO₄-Ru: 3-K phase of Sr₂RuO₄. *Phys. Rev. B* **2003**, *67*, 214519. [CrossRef]
- 31. Mao, Z.Q.; Fukazawa, H.; Maeno, Y. Crystal growth of Sr₂RuO₄. Mater. Res. Bull. 2000, 35, 1813–1824. [CrossRef]
- 32. Perry, R.S.; Maeno, Y. Systematic approach to the growth of high-quality single crystals of Sr₃Ru₂O₇. *J. Cryt. Growth* **2004**, *271*, 942–950. [CrossRef]
- Zhou, M.; Hooper, J.; Fobes, D.; Mao, Z.Q.; Golub, V.; O'Connor, C.J. Electronic and magnetic properties of triple-layered ruthenate Sr₄Ru₃O₁₀ single crystals grown by a floating-zone method. *Mater. Res. Bull.* 2005, 40, 1813–1824. [CrossRef]
- 34. Kikugawa, N.; Balicas, L.; Mackenzie, A.P. Physical Properties of Single-Crystalline CaRuO₃ Grown by a Floating-Zone Method. *J. Phys. Soc. Jpn.* **2009**, *78*, 014701. [CrossRef]
- 35. Kikugawa, N.; Baumbach, R.; Brooks, J.S.; Terashima, T.; Uji, S.; Maeno, Y. Single-Crystal Growth of a Perovskite Ruthenate SrRuO₃ by the Floating-Zone Method. *Cryst. Growth Des.* **2015**, *15*, 5573–5577. [CrossRef]
- 36. Yoshida, Y.; Nagai, I.; Ikeda, S.-I.; Shirakawa, N.; Kosaka, M.; Môri, N. Quasi-two-dimensional metallic ground state of Ca₃Ru₂O₇. *Phys. Rev. B* **2004**, *69*, 220411(R). [CrossRef]
- 37. Kikugawa, N.; Rost, A.W.; Hicks, C.W.; Schofield, A.J.; Mackenzie, A.P. Ca₃Ru₂O₇: Density Wave Formation and Quantum Oscillations in the Hall Resistivity. *J. Phys. Soc. Jpn.* **2010**, *79*, 024704. [CrossRef]
- 38. Yonezawa, S.; Higuchi, T.; Sugimoto, Y.; Sow, C.; Maeno, Y. Compact AC susceptometer for fast sample characterization down to 0.1 K. *Rev. Sci. Instrum.* **2015**, *86*, 093903. [CrossRef]



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