



Article **Theory on New Fractional Operators Using Normalization and Probability Tools**

Marc Jornet 回

Departament de Matemàtiques, Universitat de València, 46100 Burjassot, Spain; marc.jornet@uv.es

Abstract: We show how a rescaling of fractional operators with bounded kernels may help circumvent their documented deficiencies, for example, the inconsistency at zero or the lack of inverse integral operator. On the other hand, we build a novel class of linear operators with memory effects to extend the L-fractional and the ordinary derivatives, using probability tools. A Mittag–Leffler-type function is introduced to solve linear problems, and nonlinear equations are addressed with power series, illustrating the methods for the SIR epidemic model. The inverse operator is constructed, and a fundamental theorem of calculus and an existence-and-uniqueness result for differintegral equations are proven. A conjecture on deconvolution is raised, which would permit completing the proposed theory.

Keywords: fractional calculus; rescaled operator; probabilistic operator; fractional differential equation; Caputo and L-fractional derivatives; singular and non-singular kernels

MSC: 34A08; 34A25; 60E05; 33E12

1. Introduction

Fractional calculus is concerned with the investigation of linear operators that extend the ordinary derivative by including some sort of memory effect, for example, a continuous delay through integration. These operators are called fractional derivatives because they normally depend on a real positive parameter, the fractional order or index, that gives standard derivatives at integer positive values. If the fractional order is not explicitly written, then one has a memory operator. Given these types of operators, new differential or differintegral equations can be defined that include past history in the system to exhibit non-local properties. Their study, compared to ordinary differential equations, requires the development of new results on existence and uniqueness, explicit solutions (power series, Laplace transforms), and numerical methods. The reader is referred to the monographs [1–6], the review papers [7–10], and the research articles [11–15].

The most famous fractional derivatives are Riemann–Liouville and Caputo. Although the latter was proposed decades ago in viscoelasticity theory [16], they are of use in the current research, both pure and applied [17–19]. In the definition of the operators, the past delay is incorporated by integrating with respect to a singular kernel (i.e., a function with infinite value at the endpoint of the interval). Related to the type of kernel, efforts are also being devoted to non-singular integrators, for example, bounded or continuous [20,21]. However, this kind of fractional operator has certain disadvantages, as proven in some works [22].

In this paper, we consider and modify these operators proposed in the literature. We use the generic notation Dx(t) here for the operator D and its evaluation at functions x, where t is the independent variable (the time). After discussing some disadvantages of classical forms of D, we investigate whether dividing by Dt improves the properties of the operator:

$$\widetilde{D}x(t) := \frac{Dx(t)}{Dt}.$$
(1)



Citation: Jornet, M. Theory on New Fractional Operators Using Normalization and Probability Tools. *Fractal Fract.* 2024, *8*, 665. https:// doi.org/10.3390/fractalfract8110665

Academic Editor: Juan Luis González-Santander

Received: 15 October 2024 Revised: 11 November 2024 Accepted: 13 November 2024 Published: 15 November 2024



Copyright: © 2024 by the author. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). We also address the features of the associated fractional differential equation,

$$\widetilde{D}x(t) = f(t, x(t)).$$
⁽²⁾

In general, with (1) and (2), we see that the values of x'(0) and Dx(0) are more consistent, and that the vector field f has units time⁻¹. For example, if D contains a non-singular kernel, we see that the units of Dx(t) are time⁰, which entails serious drawbacks, but the division by Dt gives rise to a consistent rate time⁻¹. When D has a singular kernel, for instance, of the Caputo type, we see that the division by Dt changes the units time^{- α} to time⁻¹, rendering alternative properties and operators. Even though the normalization with Dt may seem simple, it requires the building of a new theory on fractional calculus, especially with regard to the search of explicit and closed-form solutions. In addition, it provides insight on how to design a general probabilistic definition of operator with memory that extends the ordinary derivative.

We focus, on the one hand, on the Caputo operator $D = {}^{C}D^{\alpha}$, for fractional order $\alpha \in (0, 1)$, due to its accepted used in the literature and its good properties when working with initial states, power series, and Laplace transforms in differential equations. It satisfies $^{C}D^{0^{+}}x(t) = x(t) - x(0)$ and, for fractional differential equations, $|x'(0)| = \infty$ and units time^{- α}. The infinite rate of change in the dynamics at t = 0 and the fractional units may be a problem when dealing with Caputo models, hence, the need for rescaling. The normalization gives rise to the L-fractional derivative, ${}^{L}D^{\alpha}$, suggested a few years ago in the context of geometry and mechanics [23,24] and recently studied from the mathematicalanalysis point of view [25,26]. We aim at generalizing the L-fractional operator from its probabilistic interpretation as the expectation of a certain beta-distributed delay. Albeit not the subject of our contribution, the Riemann–Liouville derivative can also be rescaled by employing the Λ -fractional operator [27]. On the other hand, in addition to Caputo operators, we also devote work to operators with bounded kernels, also named nonsingular. The normalization of operators with bounded kernels may circumvent some of the (right) critiques to their use, such as the inconsistency at zero or the lack of inverse integral operator [22]. We examine the usual exponential and Mittag–Leffler kernels.

The organization of this paper is the following. In Section 2, we revisit the mathematical treatment on L-fractional calculus, from [25]. In Section 3, we develop a novel theory on fractional operators with bounded kernels. We study how the rescaling improves their properties for an appropriate pure and applied use. In Section 4, we build a theory on new operators with memory effects, defined with probability objects. These operators extend the L-fractional and the ordinary derivatives, and they are fractional in some examples. An alternative Mittag–Leffler function is introduced. Finally, Section 5 gives limitations of the paper with open problems for the future.

Concerning notation, we denote by L^p the Lebesgue spaces and by C^p the set of functions with continuous derivatives up to order p. The set of continuous functions is C. All integrals are considered in the Lebesgue sense. The norms of functions and operators use $\|\cdot\|_*$, where * indicates the space. The absolute value of real numbers (\mathbb{R}), the modulus of complex numbers (\mathbb{C}), and the norms of real and complex vectors and matrices are simply expressed with $|\cdot|$. Given a probability space, the expectation is written as $\mathbb{E}[\cdot]$. When the expectation is performed with respect to a certain random variable U, we use $\mathbb{E}_U[\cdot]$. The essential-supremum norm, in different spaces, is $\|\cdot\|_{\infty}$. The composition is denoted with \circ . The symbol \sim captures both the asymptotic behavior of functions and distribution of random variables.

2. The L-Fractional Derivative

This section reviews concepts from [25], which builds a theory on L-fractional operators and linear fractional differential equations. There are several open problems there that might of interest for readers. For nonlinear equations, one may consult [26].

2.1. Caputo Definition

The Caputo fractional derivative is

$${}^{C}D^{\alpha}x(t) = {}^{RL}J^{1-\alpha}x'(t) = \frac{1}{\Gamma(1-\alpha)} \int_{0}^{t} \frac{x'(\tau)}{(t-\tau)^{\alpha}} d\tau,$$
(3)

where

$${}^{RL}J^{1-\alpha}x(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-\tau)^{-\alpha}x(\tau)d\tau$$
(4)

is the Riemann–Liouville integral, $\alpha \in (0, 1)$ is the fractional order, and

$$\Gamma(z) = \int_0^\infty \tau^{z-1} \mathrm{e}^{-\tau} \mathrm{d}\tau$$

is the gamma function (which extends the factorial). The operator (4) is defined for Lebesgue integrable functions $x : [0, T] \to \mathbb{C}$, i.e., $x \in L^1[0, T]$; the new function ${}^{RL}J^{1-\alpha}x$ exists almost everywhere and belongs to $L^1[0, T]$, by the properties of the convolution. The Caputo derivative (3) takes absolutely continuous functions $x : [0, T] \to \mathbb{C}$. Recall that x is said to be absolutely continuous if it is continuous, its derivative x' exists almost everywhere and belongs to $L^1[0, T]$, and Barrow's rule is satisfied in a Lebesgue sense, $x(t) - x(0) = \int_0^t x'(s) ds$.

2.2. L-Fractional Definition

The L-fractional derivative is defined as the normalization of the Caputo operator,

$${}^{L}D^{\alpha}x(t) = \frac{{}^{C}D^{\alpha}x(t)}{{}^{C}D^{\alpha}t} = \frac{\Gamma(2-\alpha)}{t^{1-\alpha}}{}^{C}D^{\alpha}x(t).$$
(5)

For absolutely continuous functions $x : [0, T] \to \mathbb{C}$, the function ${}^{L}D^{\alpha}x$ exists almost everywhere and is in $L^{1}[0, T]$. For $\alpha \in [0, 1]$, it interpolates between

$$\frac{x(t) - x(0)}{t} = \frac{1}{t} \int_0^t x'(s) \mathrm{d}s,$$

which is the mean value of the velocity on [0, T], and the ordinary derivative x'(t) if $x \in C^1[0, T]$.

When $x \in C^2[0, T]$, the operator can be rewritten as

$${}^{L}D^{\alpha}x(t) = x'(0) + \frac{1}{t^{1-\alpha}} \int_{0}^{t} (t-\tau)^{1-\alpha} x''(\tau) d\tau,$$
(6)

pointwise on (0, T]. The kernel $(t - \tau)^{1-\alpha}$ is non-singular, although the denominator $t^{1-\alpha}$ controls the value of ${}^{L}D^{\alpha}x(0+)$ to avoid inconsistencies: when $x \in C^{3}[0, T]$, for example, we have

$$^{L}D^{\alpha}x(0) := \lim_{t \to 0^{+}} ^{L}D^{\alpha}x(t) = x'(0).$$
 (7)

In that case, ${}^{L}D^{\alpha}x \in \mathcal{C}[0,T]$.

Considering (5), an L-fractional differential equation is

$${}^{L}D^{\alpha}x(t) = f(t, x(t)), \tag{8}$$

for $t \in (0, T]$, with an initial condition or state $x(0) = x_0$, where $f : [0, T] \times \Omega \subseteq [0, T] \times \mathbb{R}^d \to \mathbb{R}^d$, or $f : [0, T] \times \Omega \subseteq [0, T] \times \mathbb{C}^d \to \mathbb{C}^d$, is a continuous function such that $x_0 \in \Omega$. The Equation (8) can be understood almost everywhere or everywhere on [0, T], depending on whether *x* is smooth and (6) and (7) are of use. In contrast to Caputo differential equations, (8) may present smooth solutions, such as power series

$$\sum_{n=0}^{\infty} x_n t^n.$$
(9)

The powers are ordinary, t^n , instead of the fractional powers in the Caputo setting, $t^{\alpha n}$. This is related to the fact that the units of measurement in (8) are time⁻¹, instead of time^{- α}. For smooth solutions, note that the property $x'(0) \in (-\infty, \infty)$ is obtained by (7). Additionally, the result (7) also tells us that (8) is, locally around t = 0, very similar to x'(t) = f(t, x(t)), so the change in the dynamics with α is smoother than in Caputo equations [26].

There is a Mittag–Leffler-type function in this context, alternative to the classical one [28]:

$$\mathcal{E}_{\alpha}(s) = \sum_{n=0}^{\infty} \frac{s^n}{\Gamma(2-\alpha)^n \prod_{j=1}^n \frac{\Gamma(j+1)}{\Gamma(j+1-\alpha)}}$$

$$= \sum_{n=0}^{\infty} \frac{s^n}{\Gamma(2-\alpha)^n \Gamma(1+\alpha)^n \prod_{j=1}^n {j \choose j-\alpha}},$$
(10)

for $s \in \mathbb{C}$. The solution of

$${}^{L}D^{\alpha}x = \lambda x \tag{11}$$

where $\lambda \in \mathbb{C}$ and $t \ge 0$, is

$$x(t) = \mathcal{E}_{\alpha}(\lambda t) x_0. \tag{12}$$

Expressions (10)–(12) can be defined for matrix arguments $\lambda = A \in \mathbb{C}^{d \times d}$, as well as $d \ge 1$.

The integral operator associated with $^{L}D^{\alpha}$ is

$${}^{L}J^{\alpha}x(t) = \frac{1}{\Gamma(\alpha)\Gamma(2-\alpha)} \int_{0}^{t} (t-s)^{\alpha-1} s^{1-\alpha}x(s) \mathrm{d}s.$$
(13)

If $x \in L^1[0, T]$, then ${}^LJ^{\alpha}x \in L^1[0, T]$. If x is continuous on [0, T], then ${}^LJ^{\alpha}x$ is welldefined everywhere on [0, T]. In fact, ${}^LJ^{\alpha} : C[0, T] \to C[0, T]$ is a continuous operator, with norm

$$\|L J^{\alpha}\|_{\infty} \leq T.$$

2.3. Connection with Probability Theory

The link of L-fractional calculus with probability theory is the following: given (5) and (13), we have

 ${}^{L}J^{\alpha}y(t) = t\mathbb{E}[y(tV)]$ ⁽¹⁴⁾

and

$${}^{L}D^{\alpha}y(t) = \mathbb{E}[y'(tW)], \qquad (15)$$

where *V* is a random variable with distribution $\text{Beta}(2 - \alpha, \alpha)$, *W* is a random variable with distribution $\text{Beta}(1, 1 - \alpha)$, and \mathbb{E} is the expectation operator.

In this paper, we investigate fractional operators from the key property (15).

3. Fractional Operators with Bounded Kernels

This part presents a novel theory on fractional operators with non-singular kernels, based on rescaling. Our case studies are exponential and Mittag–Leffler integrators.

3.1. Exponential Kernel

The Caputo–Fabrizio operator [29] is defined as

$${}^{CF}D^{\alpha}x(t) = \frac{1}{1-\alpha} \int_0^t \mathrm{e}^{-\frac{\alpha}{1-\alpha}(t-s)} x'(s) \mathrm{d}s, \tag{16}$$

often in the context of continuously differentiable functions $x : [0, T] \to \mathbb{C}$. Notice that the kernel

$$\mathcal{K}(t-s) = \frac{1}{1-\alpha} e^{-\frac{\alpha}{1-\alpha}(t-s)}$$
(17)

is non-singular, because it does not present any singularity (vanishing denominator, for example) when t = s. The corresponding integral operator is

$${}^{CF}J^{\alpha}x(t) = (1-\alpha)(x(t)-x(0)) + \alpha \int_0^t x(s) \mathrm{d}s, \tag{18}$$

for continuous functions $x : [0, T] \to \mathbb{C}$. It is a convex combination between the discrete change x(t) - x(0) and the continuous change $\int_0^t x(s) ds$.

Although the use of a non-singular kernel may seem appealing, there are many drawbacks associated with the Caputo–Fabrizio operator, which are common in fractional calculus with bounded or continuous kernels [22]. First, it can be proven [30,31] that

$${}^{CF}J^{\alpha} \circ {}^{CF}D^{\alpha}x(t) = x(t) - x(0)$$
⁽¹⁹⁾

and

$${}^{CF}D^{\alpha} \circ {}^{CF}J^{\alpha}x(t) = x(t) - \mathrm{e}^{-\frac{\alpha}{1-\alpha}t}x(0).$$
⁽²⁰⁾

Note that (19) mimics Barrow's rule, but (20) does not correspond to the fundamental theorem of calculus: the derivative of the integral is not the identity operator. This fact is related to the null space of (18), given by

$$\left\langle \mathrm{e}^{-\frac{\alpha}{1-\alpha}t}\right\rangle$$
,

where $\langle \cdot \rangle$ denotes the linear span. Second, when proposing a model of the form

$$CFD^{\alpha}x(t) = f(t, x(t)), \tag{21}$$

there is a clear contradiction at t = 0, since

$$^{CF}D^{\alpha}x(0) = 0$$

is always satisfied. Hence. the only possibility is to work with the integral problem

$$x(t) = x_0 + {}^{CF} J^{\alpha} f(t, x(t)),$$
(22)

which is not entirely equivalent to (21), due to (20). Third, the units of (16) are time⁰, i.e., it has no time units. Indeed, the units of x'(s) cancel out with those of ds, while the exponential kernel (17) is non-singular and does not behave as time^{- β} for any index $\beta > 0$: by Taylor expansion,

$$e^{-\frac{\alpha}{1-\alpha}(t-s)} = 1 - \frac{\alpha}{1-\alpha}(t-s) + \frac{1}{2}\left(\frac{\alpha}{1-\alpha}\right)^2(t-s)^2 + \dots$$

These disadvantages can be resolved by working with the normalized version of (16). Simple computations yield

$${}^{CF}D^{\alpha}t = \frac{1}{\alpha} \left(1 - \mathrm{e}^{-\frac{\alpha}{1-\alpha}t} \right) \sim \frac{1}{1-\alpha}t, \tag{23}$$

where \sim denotes here that the two functions are asymptotically equivalent when $t \rightarrow 0^+$, which is the problematic point. Considering (23), we define

$$c_{\alpha}(t) = \frac{1}{\alpha} \left(1 - \mathrm{e}^{-\frac{\alpha}{1-\alpha}t} \right)$$

and

$${}^{NCF}D^{\alpha}x(t) = \frac{1}{(1-\alpha)c_{\alpha}(t)} \int_{0}^{t} e^{-\frac{\alpha}{1-\alpha}(t-s)}x'(s) ds.$$
(24)

Notice that convergence towards x'(t) when $\alpha \to 1^-$ holds for (24). Furthermore,

$$^{NCF}D^{\alpha}x(0^{+}) = \frac{x(t) - x(0)}{t},$$
 (25)

which is the mean value

$$\frac{1}{t}\int_0^t x'(s)\mathrm{d}s.$$

That is, (24) interpolates between the mean velocity on [0, t] and at t.

It is clear that the units of (24) are time⁻¹, by the division by $c_{\alpha}(t)$. Therefore, the new operator represents some sort of rate of change in contrast to the standard one (16). Furthermore, now (24) is a convolution with respect to a probability density function on [0, t], for each t > 0, so that the weight function seems to be more appropriate.

Definition (24) is well posed for t = 0:

$${}^{NCF}D^{\alpha}x(0) = \lim_{t \to 0^+} e^{-\frac{\alpha}{1-\alpha}t} \cdot \frac{1}{t} \int_0^t e^{\frac{\alpha}{1-\alpha}s} x'(s) ds = x'(0).$$
(26)

Even for $\alpha = 0$, it is true that ${}^{NCF}D^{\alpha}x(0) = x'(0)$ by (25). In (26), we used $(1 - \alpha)c_{\alpha}(t) \sim t$ and decomposed the exponential function to readily apply the fundamental theorem of calculus, but for a general continuous kernel, one can proceed with the Leibniz rule of differentiation for integrals or integration by parts and arrive at the same result. Thus, the new fractional operator has an appropriate value at t = 0, the same as the integer-order one, and one can work with fractional differential equations of the form

$$^{NCF}D^{\alpha}x(t) = f(t, x(t)), \tag{27}$$

as opposed to (21). The associated integral operator (see (18)),

$${}^{NCF}J^{\alpha}x(t) = (1-\alpha)c_{\alpha}(t)x(t) + \alpha \int_0^t c_{\alpha}(s)x(s)\mathrm{d}s, \tag{28}$$

does satisfy the fundamental theorem of calculus, in contrast to (20): for a continuously differentiable function x,

$${}^{NCF}J^{\alpha} \circ {}^{NCF}D^{\alpha}x(t) = {}^{NCF}J^{\alpha}\left[\frac{1}{c_{\alpha}} \cdot {}^{CF}D^{\alpha}x\right](t)$$
$$= {}^{CF}J^{\alpha} \circ {}^{CF}D^{\alpha}x$$
$$= x(t) - x(0)$$

and

$${}^{NCF}D^{\alpha} \circ {}^{NCF}J^{\alpha}x(t) = {}^{NCF}D^{\alpha} \circ {}^{CF}J^{\alpha}[c_{\alpha}x](t)$$

$$= \frac{1}{c_{\alpha}(t)} \cdot {}^{CF}D^{\alpha} \circ {}^{CF}J^{\alpha}[c_{\alpha}x](t)$$

$$= \frac{1}{c_{\alpha}(t)}c_{\alpha}(t)x(t)$$

$$= x(t),$$
(29)

since $c_{\alpha}(t)x(t)$ is 0 at t = 0. We employed both (19) and (20). Observe that t = 0 is not a problem for (29), by (26) and (28):

$${}^{NCF}D^{\alpha}\Big[{}^{NCF}J^{\alpha}x\Big](0) = \Big({}^{NCF}J^{\alpha}x\Big)'(0) = x(0).$$

Equation (27) is now completely equivalent to

$$x(t) = x_0 + {}^{NCF} J^{\alpha} f(t, x(t)),$$
(30)

in the set of $C^1[0, T]$ functions.

We do not enter into other possible issues regarding the Caputo–Fabrizio models. For example, (22) is equivalent to an ordinary differential equation by differentiating. The same fact occurs for the alternative model (27), which gives

$$x'(t) = (1 - \alpha)c'_{\alpha}(t)f(t, x(t)) + (1 - \alpha)c_{\alpha}(t)[f_{t}(t, x(t)) + x'(t)f_{x}(t, x(t))] + \alpha c_{\alpha}(t)f(t, x(t)),$$
(31)

by (28) and (30), so the non-local behavior of (27) is debatable. As (31) is well understood, fractional models with an exponential kernel do not seem to give rise to a big new theory, except concrete mathematical results. When we convolve the derivative x' with an exponential in (16), the exponential function decomposes as a product and separates the variables t (outside the integral) and s (inside the integral), and this gives rise to an ordinary differential equation at the end. Since the weight function in (24) is related to the exponential probability distribution, which exhibits the memoryless property, it is intuitive that the non-local behavior of (27) is lost.

Example 1. The model

$$CFD^{\alpha}x(t) = -\lambda x(t), \tag{32}$$

where t > 0 and $\lambda > 0$, extends the exponential decay equation to a fractional setting, for $\alpha \in (0, 1]$. First, observe that if $x(0) = x_0 \neq 0$, then (32) is not well-defined, because ${}^{CF}D^{\alpha}x(0) = 0 \neq -\lambda x_0$. Additionally, as already proven, the units of ${}^{CF}D^{\alpha}x(t)$ are time⁰; hence, there is no valid power for λ ; it does not represent any rate. Finally, the differintegral Equation (32) is not equivalent to the integral equation $x(t) = x_0 + {}^{CF}J^{\alpha}[-\lambda x(t)]$ by (20). The alternative model

$$\mathcal{N}CFD^{\alpha}x(t) = -\lambda x(t) \tag{33}$$

fixes all these issues. It is consistent at t = 0, with ${}^{NCF}D^{\alpha}x(0) = x'(0) = -\lambda x_0$; the units of ${}^{NCF}D^{\alpha}x(t)$ and λ are time⁻¹, thus representing true rates; and finally, the differintegral Equation (33) is equivalent to the integral equation $x(t) = x_0 + {}^{NCF}J^{\alpha}[-\lambda x(t)]$, because the fundamental theorem of calculus holds. The argumentation for this simple example (33) works for any other model defined via a non-singular kernel. During the time the present paper has been in the preprint server ArXiv and under review, specific models based on my proposed normalized operators are being studied [32].

3.2. Mittag-Leffler Kernel

Other possible operators with bounded kernels are based on Mittag–Leffler kernels instead of exponential ones [33,34]. For example,

$${}^{AB}D^{\alpha}x(t) = \frac{1}{1-\alpha} \int_0^t E_{\alpha} \left(-\frac{\alpha}{1-\alpha}(t-s)^{\alpha}\right) x'(s) \mathrm{d}s. \tag{34}$$

Here, one often considers absolutely continuous functions $x : [0, T] \to \mathbb{C}$. The kernel

$$\mathcal{K}(t-s) = \frac{1}{1-\alpha} E_{\alpha} \left(-\frac{\alpha}{1-\alpha} (t-s)^{\alpha} \right)$$

is non-singular. Observe that, again,

$$^{AB}D^{\alpha}x(0)=0,$$

which is a severe restriction in differential equations (exactly the same issue as for Caputo–Fabrizio models) [22]. The units of (34) are also time⁰ (the use of the exponent α does not change this fact). Since

$${}^{AB}D^{\alpha}t^{\gamma} = \frac{1}{1-\alpha}\Gamma(1+\gamma)t^{\gamma}E_{\alpha,\gamma+1}\left(-\frac{\alpha}{1-\alpha}t^{\alpha}\right),$$

where $\gamma > 0$, the rescaled operator associated with (34) is

$$^{NAB}D^{\alpha}x(t) = \frac{1}{(1-\alpha)k_{\alpha}(t)} \int_{0}^{t} E_{\alpha}\left(-\frac{\alpha}{1-\alpha}(t-s)^{\alpha}\right)x'(s)\mathrm{d}s,\tag{35}$$

being

$$k_{\alpha}(t) = rac{1}{1-lpha} t E_{lpha,2} \left(-rac{lpha}{1-lpha} t^{lpha}
ight) \sim rac{1}{1-lpha} t.$$

The new fractional operator (35) satisfies

$${}^{NAB}D^{\alpha}x(0) = x'(0)$$

when calculating the limit as $t \to 0^+$, if x' is continuous, by the Leibniz rule of differentiation for integrals:

$$\begin{split} \lim_{t \to 0^+} {}^{NAB} D^{\alpha} x(t) &= \lim_{t \to 0^+} \frac{1}{t} \int_0^t E_{\alpha} \left(-\frac{\alpha}{1-\alpha} (t-s)^{\alpha} \right) x'(s) \mathrm{d}s \\ &= \frac{\mathrm{d}}{\mathrm{d}t} \bigg|_{t=0^+} \int_0^t E_{\alpha} \left(-\frac{\alpha}{1-\alpha} (t-s)^{\alpha} \right) x'(s) \mathrm{d}s \\ &= x'(0) + \frac{\alpha^2}{1-\alpha} \lim_{t \to 0^+} \int_0^t (t-s)^{\alpha-1} E'_{\alpha} \left(-\frac{\alpha}{1-\alpha} (t-s)^{\alpha} \right) x'(s) \mathrm{d}s \\ &= x'(0). \end{split}$$

Hence there are no issues at t = 0, compared with (34). The units of (35) are time⁻¹. Now (35) is a convolution with respect to a probability density function on [0, t], for each t > 0, so that the weight function is likely more adequate.

The integral operator of (34) is

$${}^{AB}J^{\alpha}x(t) = (1-\alpha)(x(t)-x(0)) + \alpha \cdot {}^{RL}J^{\alpha}x(t),$$
(36)

which is a convex combination of the discrete change x(t) - x(0) and the continuous delayed change ${}^{RL}J^{\alpha}x(t)$. From (36), one deduces that

$$^{NAB}J^{\alpha}x(t) = (1-\alpha)k_{\alpha}(t)x(t) + \alpha \cdot ^{RL}J^{\alpha}[k_{\alpha}x](t)$$
(37)

is the integral counterpart of (35). It is well known that

$${}^{AB}J^{\alpha} \circ {}^{AB}D^{\alpha}x(t) = x(t) - x(0)$$

and

$${}^{AB}D^{\alpha} \circ {}^{AB}J^{\alpha}x(t) = x(t) - E_{\alpha}\left(-\frac{\alpha}{1-\alpha}t^{\alpha}\right)x(0) \neq x(t),$$

but now

$$^{NAB}J^{\alpha} \circ {}^{NAB}D^{\alpha}x(t) = {}^{NAB}J^{\alpha} \left[\frac{1}{k_{\alpha}(t)} \cdot {}^{AB}D^{\alpha}x\right](t)$$
$$= {}^{AB}J^{\alpha} \circ {}^{AB}D^{\alpha}x$$
$$= x(t) - x(0)$$

and

$$\begin{split} {}^{NAB}D^{\alpha} \circ^{NAB}J^{\alpha}x(t) = {}^{NAB}D^{\alpha} \circ^{AB}J^{\alpha}[k_{\alpha}x](t) \\ = &\frac{1}{t} \cdot {}^{AB}D^{\alpha} \circ^{AB}J^{\alpha}[k_{\alpha}x](t) \\ = &\frac{1}{k_{\alpha}(t)}k_{\alpha}(t)x(t) \\ = &x(t), \end{split}$$

because $k_{\alpha}(0)x(0) = 0$. Therefore, the fundamental theorem of calculus for the new operator $^{NAB}D^{\alpha}$ is verified.

We have the equivalence between

$$^{NAB}D^{\alpha}x(t) = f(t, x(t))$$

and

$$x(t) = x_0 + {}^{NAB}J^{\alpha}f(t, x(t))$$

in the set of $C^1[0, T]$ functions. This model is the same as a system with implicit Caputo derivatives only,

$${}^{C}D^{\alpha}x(t) = (1-\alpha) \cdot {}^{C}D^{\alpha}(k_{\alpha}(t)f(t,x(t))) + \alpha k_{\alpha}(t)f(t,x(t))$$

almost everywhere, by (37) and the fundamental theorem of calculus. This fact is natural somehow, because the Mittag–Leffler function is directly related to Caputo fractional differential equations (analogously, the exponential function in the Caputo-Fabrizio model is directly related to ordinary differential equations): it builds the solution of the basic equation $^{C}D^{\alpha}z(t) = \lambda^{\alpha}z(t)$ (analogously, the exponential function builds the solution of the basic equation $z'(t) = \lambda z(t)$).

In conclusion, the factor $(Dt)^{-1}$ seems to resolve major issues associated with fractional operators with bounded kernels. Some authors proposed the use of integration by parts directly in (16) and (34), see [22,35]:

$${}^{CF}\widetilde{D}^{\alpha}x(t) = \frac{1}{1-\alpha} \bigg[x(t) - \mathrm{e}^{-\frac{\alpha}{1-\alpha}t}x(0) - \frac{\alpha}{1-\alpha} \int_0^t \mathrm{e}^{-\frac{\alpha}{1-\alpha}(t-s)}x(s)\mathrm{d}s \bigg], \tag{38}$$

$$^{AB}\widetilde{D}^{\alpha}x(t) = \frac{1}{1-\alpha} \left[x(t) - E_{\alpha} \left(-\frac{\alpha}{1-\alpha} t^{\alpha} \right) x(0) - \frac{\alpha}{1-\alpha} \int_{0}^{t} E_{\alpha} \left(-\frac{\alpha}{1-\alpha} (t-s)^{\alpha} \right) x(s) \mathrm{d}s \right],$$
(39)

for Lebesgue integrable functions *x*. However, since there is an evaluation at 0, the continuity of *x* at t = 0 is needed for uniqueness of the fractional derivative. In that case, the operators (38) and (39) still tend to 0 when $t \rightarrow 0^+$, and the inconsistency persists. Additionally, the issue of the units time⁰ is present again. There is no derivative in (38) and (39), which contradicts the notion of the fractional *derivative*; indeed, some rate of change should be involved in the formulation.

In terms of explicit or closed-form solutions, a disadvantage of the rescaling of fractional operators in differential equations is that the applicability of the Laplace-transform method is reduced, because the transform of the product or division is not amenable to computing. Nonetheless, the power-series technique may still be of application, as one can find the power series of a product easily.

4. A New Operator with Memory Effects Using Probability Theory

This part of the paper generalizes normalized operators by employing probability concepts. A theory on the new operators is built.

4.1. Definition

By generalizing the probabilistic interpretation of the L-fractional derivative, see (15), we define the new linear operator

$$\mathcal{D}x(t) = \mathbb{E}[x'(tW)], \tag{40}$$

where $t \in [0, T]$, $x : [0, T] \to \mathbb{C}$ is (at least) an absolutely continuous function, and *W* is a random variable with the requirements:

$$support(W) \subseteq [0,1],$$
 (41)

$$\parallel W \parallel_{\infty} = 1, \tag{42}$$

and

$$\lim_{n \to \infty} n \mathbb{E}[W^n] = \infty.$$
(43)

The notation $\|\cdot\|_{\infty}$ is the essential supremum of the random variable. Both (41) and (42) are related to the need of capturing the past history of x' until the present, along with [0, t], i.e., exhibiting memory effects. Physically, the new operator (40) is a weighted average of the velocity on [0, t], where the weight is modeled with a probability distribution. The third condition (43) will be required later, when dealing with differential equations and series, so that the new "exponential" function exists on \mathbb{R} . Observe that (43) implies (42), but we explicitly write (42) for the sake of clarity. Assumption (43) means that the past time near the present t should have enough weight in the memory operator.

We observe that x' exists almost everywhere. In definition (40), we are assuming that

$$\mathbb{E}[|x'(tW)|] < \infty. \tag{44}$$

For example, (44) follows if x' is essentially bounded on [0, T], i.e., $||x'||_{\infty} < \infty$. Some examples of \mathcal{D} are the following:

- When W = 1 is constant, then $\mathcal{D}x = x'$ is the classical derivative. Note that the three assumptions (41)–(43) hold.
- When $W \sim \text{Beta}(1, 1 \alpha)$, for $\alpha \in (0, 1)$, then $\mathcal{D}x = {}^{L}D^{\alpha}x$ is the L-fractional derivative. The third condition (43) is satisfied:

$$\lim_{n \to \infty} n \mathbb{E}[W^n] = \lim_{n \to \infty} n \frac{\Gamma(2 - \alpha) \Gamma(n)}{\Gamma(n + 1 - \alpha)}$$
$$= \lim_{n \to \infty} n \frac{\Gamma(2 - \alpha)}{n^{1 - \alpha}}$$
$$= \infty.$$

If W ~ Beta(α, β), for α, β ∈ (0, 1), then we obtain a generalization of the L-fractional derivative [25], denoted as ^{GL}D^{α,β}. The condition (43) is verified as follows:

$$\lim_{n \to \infty} n \mathbb{E}[W^n] = \lim_{n \to \infty} n \frac{\Gamma(\alpha + \beta)\Gamma(\alpha + n - 1)}{\Gamma(\alpha)\Gamma(\alpha + n - 1 + \beta)}$$
$$= \lim_{n \to \infty} n \frac{\Gamma(\alpha + \beta)}{\Gamma(\alpha)(\alpha + n - 1)^{\beta}}$$
(45)
$$= \infty.$$

11 of 24

Regarding (44), simple computations for integrals yield

$${}^{GL}D^{\alpha,\beta}x(t) = \frac{\Gamma(\alpha+\beta)}{\Gamma(\alpha)\Gamma(\beta)} \int_0^1 s^{\alpha-1} (1-s)^{\beta-1} x'(ts) ds$$

$$= \frac{\Gamma(\alpha+\beta)}{\Gamma(\alpha)\Gamma(\beta)} \frac{1}{t} \int_0^t \frac{s^{\alpha-1}}{t^{\alpha-1}} \left(1-\frac{s}{t}\right)^{\beta-1} x'(s) ds$$

$$= \frac{\Gamma(\alpha+\beta)}{\Gamma(\alpha)\Gamma(\beta)} \frac{1}{t^{\alpha+\beta-1}} \int_0^t (t-s)^{\beta-1} s^{\alpha-1} x'(s) ds$$

$$= \frac{1}{t^{\alpha+\beta-1}} \left(\frac{\Gamma(\alpha+\beta)}{\Gamma(\alpha)\Gamma(\beta)} t^{\beta-1} * \left(t^{\alpha-1} x'\right)\right)(t),$$
(46)

for t > 0. The density function of the beta distribution and the expression for the expectation have been used. Thus, if *x* is absolutely continuous and

$$t^{\alpha - 1} x' \in L^1[0, T], \tag{47}$$

then ${}^{GL}D^{\alpha,\beta}x$ exists almost everywhere and belongs to $L^1[0, T]$ (recall that a convolution $\mathcal{K} * y$ is well-defined almost everywhere and belongs to $L^1[0, T]$ if $\mathcal{K}, y \in L^1[0, T]$). Note that, if x' is essentially bounded on [0, T], that is, $|| x' ||_{\infty} < \infty$, the integrability condition (47) follows. The point t = 0 is not a problem, because

$${}^{GL}D^{\alpha,\beta}x(t) = \mathbb{E}\left[y'(0\cdot W)\right] = y'(0)$$

is well-defined. Note that (46) clearly extends the integral expression for the L-fractional operator, thus modifying the Caputo operator as well.

Based on the L-fractional derivative, when $\alpha \to 1^-$, $\beta \to 0^+$, and *x* is continuously differentiable on [0, T], the ordinary derivative operator is retrieved. If $\alpha \to 1^-$ and $\beta \to 1^-$, then the mean value

$$\frac{1}{t}\int_0^t x'(s)\mathrm{d}s$$

is obtained.

• If $W \sim \text{Uniform}(0, 1)$, i.e., Beta(1, 1), then (43) does not hold: by (45),

$$\lim_{n\to\infty} n\mathbb{E}[W^n] = 1.$$

In consequence, that simple probability distribution is not valid. This is an interesting example, because the uniform distribution should be a good option a priori, as it maximizes the Shannon entropy (the ignorance) on [0, t] when multiplied by t if there is no information available on the weight [36].

• Important cases related to the gamma and exponential distributions are given in Examples 2 and 3, in the context of inverse operators and the fundamental theorem of calculus. New operators with memory will be obtained.

Given (40), one can define the concept of differential equation as

$$\mathcal{D}x(t) = f(t, x(t)). \tag{48}$$

In general, we work in dimension $d \ge 1$, with $x : [0, T] \to \mathbb{C}^d$. Equation (48) has an initial condition or state $x(0) = x_0 \in \mathbb{C}^d$, where $f : [0, T] \times \Omega \subseteq [0, T] \times \mathbb{R}^d \to \mathbb{R}^d$, or $f : [0, T] \times \Omega \subseteq [0, T] \times \mathbb{C}^d \to \mathbb{C}^d$, is a continuous function such that $x_0 \in \Omega$. The units in (48) are time⁻¹. Compared to an ordinary differential equation x'(t) = f(t, x(t)), now we are relating the weighted average of the velocity on [0, t] and the position at *t*. In differential form,

$$\mathbb{E}[x(tW + dt)] = \mathbb{E}[x(tW)] + f(t, x(t))dt.$$

The connection between the past [0, t] and the future (t, t + dt] is stronger than in the ordinary case, similar to the difference between non-Markovian and Markovian processes. If *x* is changed at *t*, then this has an effect on the whole history of *x*, not just on (t, t + dt]. In this work, we deal with (48) abstractly.

4.2. Power Series and Mittag-Leffler-Type Function

We study the differentiation rules of power series.

Proposition 1. Consider the operator D from Section 4.1. If

$$x(t) = \sum_{n=0}^{\infty} x_n t^n \tag{49}$$

is a convergent power series on [0, T], where $x_n \in \mathbb{C}^d$, then

$$\mathcal{D}x(t) = \sum_{n=0}^{\infty} x_n \mathcal{D}t^n = \sum_{n=0}^{\infty} x_{n+1}(n+1)\mathbb{E}[W^n]t^n,$$
(50)

pointwise on [0, T].

Proof. Since *x* is bounded on [0, T], condition (44) is met, and $\mathcal{D}x$ exists pointwise on [0, T]. Let

$$x_N(t) = \sum_{n=0}^N x_n t^n$$

be the truncated sum of (49). It is well-known that

$$\lim_{N\to\infty} \|x_N - x\|_{\infty} = 0$$

and

$$\lim_{N \to \infty} \parallel x'_N - x' \parallel_{\infty} = 0$$

on [0, T]. Furthermore, by linearity,

$$\mathcal{D}x_N(t) = \sum_{n=0}^N x_n \mathcal{D}t^n$$

Then, from the definition (40) for \mathcal{D} , we derive

$$\| \mathcal{D}x - \mathcal{D}x_N \|_{\infty} \leq \| x' - x'_N \|_{\infty} \stackrel{n \to \infty}{\longrightarrow} 0.$$

Thus, the first equality of (50) is proven. For the second equality in (50), compute

$$\mathcal{D}t^n = nt^{n-1}\mathbb{E}\Big[\mathsf{W}^{n-1}\Big],$$

for $n \ge 1$. \Box

In the context of Section 4.1, we define the new Mittag–Leffler-type function

$$\mathcal{E}(s) = \sum_{n=0}^{\infty} \frac{s^n}{n! \prod_{j=1}^{n-1} \mathbb{E}[W^j]},$$
(51)

for $s \in \mathbb{C}$. The empty product is interpreted, as usual, as 1. Of course, when $W \sim \text{Beta}(1, 1 - \alpha)$, we have the L-fractional derivative, and the Mittag–Leffler-type function (10) is retrieved, after computing the moments of that beta distribution. When W = 1, the usual exponential function is obtained. Observe that (51) is convergent on \mathbb{C} , by the ratio test and condition (43):

$$\lim_{n \to \infty} \frac{n! \prod_{j=1}^{n-1} \mathbb{E}[W^j]}{(n+1)! \prod_{i=1}^n \mathbb{E}[W^i]} = \lim_{n \to \infty} \frac{1}{(n+1)\mathbb{E}[W^n]} = 0.$$

The same reasoning can be conducted for matrix arguments $s = A \in \mathbb{C}^{d \times d}$. The function (51) is key to solving linear models.

Proposition 2. Consider the operator \mathcal{D} from Section 4.1. If $A \in \mathbb{C}^{d \times d}$ and $x_0 \in \mathbb{C}^d$, then

$$x(t) = \mathcal{E}(tA)x_0$$

solves

$$\mathcal{D}x(t) = Ax(t)$$

with $x(0) = x_0$, pointwise on $[0, \infty)$.

Proof. By Proposition (1),

$$\mathcal{D}x(t) = \sum_{n=0}^{\infty} A^{n+1} x_0 \frac{(n+1)\mathbb{E}[W^n]}{(n+1)!\prod_{j=1}^n \mathbb{E}[W^j]} t^n$$

= $A \sum_{n=0}^{\infty} A^n x_0 \frac{1}{n!\prod_{j=1}^{n-1} \mathbb{E}[W^j]} t^n$
= $Ax(t).$

4.3. Fundamental Theorem of Calculus

Let *W* be a fixed random variable satisfying (41)–(43). Suppose that there exists another random variable *V*, independent of *W* and with support in [0, 1], such that

$$WV = U \sim \text{Uniform}(0, 1). \tag{52}$$

For example, when $W \sim \text{Beta}(1, 1 - \alpha)$ in L-fractional calculus, $\alpha \in (0, 1)$, then $V \sim \text{Beta}(2 - \alpha, \alpha)$, because

$$\mathbb{E}[U^n] = \mathbb{E}[W^n] \mathbb{E}[V^n]$$

$$= \frac{\Gamma(1+n)\Gamma(1-\alpha)}{\Gamma(2+n-\alpha)} \frac{\Gamma(2-\alpha)}{\Gamma(1)\Gamma(1-\alpha)}$$

$$\times \frac{\Gamma(2-\alpha+n)\Gamma(\alpha)}{\Gamma(2+n)} \frac{\Gamma(2)}{\Gamma(2-\alpha)\Gamma(\alpha)}$$

$$= \frac{1}{1+n}$$

are the moments of a Uniform(0, 1) distribution (which uniquely determine it).

Remark 1. In general, the structure of V in (52) is not simple. Let $W \sim Beta(\alpha, \beta)$ with $0 < \alpha, \beta < 1$. Suppose that $V \sim Beta(\tilde{\alpha}, \tilde{\beta})$. Equate moments

$$\mathbb{E}[V^r] = \frac{\mathbb{E}[U^r]}{\mathbb{E}[W^r]}$$

for $r \ge 1$:

$$\frac{\Gamma(\tilde{\alpha}+r)\Gamma\left(\tilde{\alpha}+\tilde{\beta}\right)}{\Gamma\left(\tilde{\alpha}+r+\tilde{\beta}\right)\Gamma(\tilde{\alpha})} = \frac{\Gamma(\alpha)\Gamma(\alpha+\beta+r)}{\Gamma(\alpha+\beta)(1+r)\Gamma(\alpha+r)}.$$
(53)

Asymptotically, the left-hand side of (53) is of the form

$$\frac{1}{(\widetilde{\alpha}+r)^{\widetilde{\beta}}}\frac{\Gamma\left(\widetilde{\alpha}+\widetilde{\beta}\right)}{\Gamma(\widetilde{\alpha})},$$

while the right-hand side of (53) is

$$\frac{\Gamma(\alpha)}{\Gamma(\alpha+\beta)}\frac{(\alpha+r)^{\beta}}{1+r}$$

 $\widetilde{\beta} = 1 - \beta \in (0, 1)$

In consequence,

and

$$\frac{\Gamma(\widetilde{\alpha}+1-\beta)}{\Gamma(\widetilde{\alpha})} = \frac{\Gamma(\alpha)}{\Gamma(\alpha+\beta)}.$$
(54)

Since $\alpha, \beta \in (0, 1)$ *, we know that*

$$\Gamma(\alpha) > \Gamma(\alpha + \beta). \tag{55}$$

If $\tilde{\alpha} \in (0, 1]$, then we would have $\Gamma(\tilde{\alpha}) > \Gamma(\tilde{\alpha} + 1 - \beta)$, but this is impossible considering (54) and (55). *Hence,*

 $\widetilde{\alpha} > 1.$

On the other hand, (53) becomes

$$\frac{\Gamma(\tilde{\alpha}+r)}{\Gamma(\tilde{\alpha}+r+1-\beta)} = \frac{\Gamma(\alpha+\beta+r)}{(1+r)\Gamma(\alpha+r)},$$
(56)

for all $r \geq 1$. Let $\mu = \tilde{\alpha} - \beta$. Observe that

$$(1+r)\Gamma(\alpha+r) > (\alpha+r)\Gamma(\alpha+r) = \Gamma(\alpha+r+1),$$

so, from (56),

$$\frac{\Gamma(\mu+r+\beta)}{\Gamma(\mu+r+1)} < \frac{\Gamma(\alpha+r+\beta)}{\Gamma(\alpha+r+1)}.$$

Since Γ increases faster and faster as r grows, we deduce that $\mu > \alpha$, that is,

 $\widetilde{\alpha} > \min\{1, \alpha + \beta\}.$

Computations in software show that if $\alpha = 0.99$ and $\beta = 0.5$, for example, then the unique root of (56) for r = 1 is $\tilde{\alpha} \approx 1.51284$, whereas for r = 2, it is $\tilde{\alpha} \approx 1.51183$, which is distinct; therefore, V cannot follow the beta distribution. For $\alpha = 1$ and $\beta \in (0, 1)$, we are in the situation of L-fractional calculus, where V is indeed beta distributed. The root $\tilde{\alpha}$ in that case is unique.

An open problem regarding (52) is the following:

Conjecture 1. *Let W* be a random variable satisfying (41)–(43). Then, under certain conditions on *W*, there exists a random variable V, independent from W and with support in [0, 1], such that (52) holds. Note that from

$$-\log W - \log V = -\log U \sim \text{Exponential}(1), \tag{57}$$

we essentially need to decompose Exponential(1) as a sum of non-negative and independent random variables, where

$$W = -\log W$$

is previously fixed with the properties

 $\mathbb{P}\Big[\widetilde{W}\in[0,\epsilon)\Big]>0$ for all $\epsilon>0$,

where \mathbb{P} is the probability measure, and

$$\lim_{n\to\infty} n\mathbb{E}\left[\mathrm{e}^{-n\widetilde{W}}\right] = 0$$

Comments on Conjecture 1: Note that (52) implies that

$$\mathbb{E}[W^r]\mathbb{E}[V^r] = \mathbb{E}[U^r] = \frac{1}{1+r},$$

i.e.,

$$\mathbb{E}[V^r] = \frac{1}{(1+r)\mathbb{E}[W^r]}.$$
(58)

The question is whether there exists V with those specific moments (58). On the other hand, working with (57) and characteristic functions, the point is whether

$$\varphi(u) = \frac{\varphi_{\text{Ex}(1)}(u)}{\varphi_{\widetilde{W}}(u)}$$
(59)

is a characteristic function of a non-negative random variable. Finally, dealing with (57) and density functions, the question is whether the convolution

$$\left(\rho_{\widetilde{W}}*\rho_{-\log V}\right)(t)=\rho_{\mathrm{Ex}(1)}(t)=\mathrm{e}^{-t}, \quad t>0,$$

is invertible, i.e., computing the deconvolution of functions.

Example 2. Let us see a positive example for Conjecture 1, distinct to the L-fractional context. Consider

$$\widetilde{W} = -\log W \sim \operatorname{Gamma}(\alpha, 1)$$

and

$$V = -\log V \sim \text{Gamma}(1 - \alpha, 1)$$

independent, where $0 < \alpha < 1$ is the shape and 1 is the rate. Then (57) holds. Additionally, for

$$W = e^{-W} \sim e^{-\text{Gamma}(\alpha, 1)},\tag{60}$$

condition (43) is true, from the moment-generating function of the gamma distribution:

$$\lim_{n \to \infty} n \mathbb{E}[W^n] = \lim_{n \to \infty} n \mathbb{E}\left[e^{-n\widetilde{W}}\right]$$
$$= \lim_{n \to \infty} \frac{n}{(1+n)^{\alpha}} = \infty.$$

In conclusion, W defined by (60) is a suitable random variable to define a fractional operator (40) of one parameter α , and its use in modeling shall be investigated. When $\alpha \rightarrow 1^-$, the ordinary derivative is obtained. As now will be seen, (62) becomes its inverse integral operator \mathcal{J} . From the densities

$$\rho_W(w) = \rho_{\widetilde{W}}(-\log w)\frac{1}{w} = \frac{1}{\Gamma(\alpha)}(-\log w)^{\alpha-1},$$

the explicit expressions of the fractional operators are the following:

$$\mathcal{D}x(t) = \frac{1}{\Gamma(\alpha)t} \int_0^t \left(-\log\frac{s}{t}\right)^{\alpha-1} x'(s) ds$$
$$= \frac{1}{\Gamma(\alpha)t} \int_0^t \left(\log\frac{t}{s}\right)^{\alpha-1} x'(s) ds$$

and

$$\mathcal{J}x(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t \left(-\log\frac{s}{t}\right)^{-\alpha} x(s) ds$$
$$= \frac{1}{\Gamma(1-\alpha)} \int_0^t \left(\log\frac{t}{s}\right)^{-\alpha} x(s) ds.$$

Example 3. Let us see another valid distribution for Conjecture 1, distinct to the gamma. Consider

$$\widetilde{W} = -\log W \sim \text{Bernoulli}(1/2) \cdot \text{Exponential}(1)$$

and

$$\widetilde{V} = -\log V \sim \frac{1}{2}$$
Exponential(1),

independent. Recall that $X \sim Bernoulli(1/2)$ if it takes the values 0 and 1 with probability 1/2, respectively, and $Y \sim Exponential(1)$ if its density function is e^{-y} . Observe that \widetilde{W} is not absolutely continuous and its density function is expressed in terms of the Dirac delta function. The moment-generating functions are

$$\mathbb{E}\left[e^{\widetilde{W}u}\right] = \frac{1}{2}\frac{1}{1-u} + \frac{1}{2} = \frac{2-u}{2-2u}$$

and

therefore,

$$\mathbb{E}\left[\mathrm{e}^{\widetilde{V}u}\right] = \frac{1}{1 - u/2};$$

 $\mathbb{E}\left[\mathrm{e}^{(\widetilde{W}+\widetilde{V})u}\right] = \frac{2-u}{2-2u} \times \frac{1}{1-u/2} = \frac{1}{1-u},$

that is, (57) holds. Note that

$$\lim_{n \to \infty} n \mathbb{E}[W^n] = \lim_{n \to \infty} n \mathbb{E}\left[e^{-n\widetilde{W}}\right]$$
$$= \lim_{n \to \infty} n \frac{2+n}{2+2n} = \infty,$$

so (43) is verified. As now will be seen, (62) is the inverse operator \mathcal{J} of (40). The explicit expressions of the operators are

$$\begin{aligned} \mathcal{D}x(t) = & \mathbb{E}\left[x'\left(te^{-\widetilde{W}}\right)\right] \\ = & \frac{1}{2}\mathbb{E}\left[x'(tU)\right] + \frac{1}{2}x'(t) = \frac{1}{2}\frac{x(t) - x(0)}{t} + \frac{1}{2}x'(t) \end{aligned}$$

and

$$\begin{aligned} \mathcal{J}x(t) =& t \mathbb{E}\left[x\left(t\mathrm{e}^{-\widetilde{V}}\right)\right] \\ =& \int_0^t x(s)\rho_V(s/t)\mathrm{d}s = \frac{2}{t}\int_0^t sx(s)\mathrm{d}s \end{aligned}$$

The interpretation of $\mathcal{D}x(t)$ in this case is very simple in terms of x': the average between the historical mean value $\frac{1}{t} \int_0^t x'(s) ds$ and the value at the present time *t*.

Example 4. We present a negative example for Conjecture 1, which shows that certain conditions on *W* are required in general for the existence of *V*. Let

$$\widetilde{W} = -\log W \sim 2 \cdot \text{Bernoulli}(1/2) \cdot \text{Exponential}(1).$$
 (61)

Then

 \mathbf{SO}

$$\mathbb{E}\left[e^{\widetilde{W}u}\right] = \frac{1}{2}\frac{1}{1-2u} + \frac{1}{2} = \frac{1-u}{1-2u},$$
$$\lim_{n \to \infty} n\mathbb{E}[W^n] = \lim_{n \to \infty} n\mathbb{E}\left[e^{-n\widetilde{W}}\right]$$
$$= \lim_{n \to \infty} n\frac{1+n}{1+2n} = \infty,$$

and (43) holds. Now, we have the following ratio (59) of characteristic functions:

$$\begin{split} \varphi(u) &= \frac{\varphi_{\text{Ex}(1)}(u)}{\varphi_{\widetilde{W}}(u)} \\ &= \frac{1}{\frac{1-ui}{1-2ui}} \\ &= \frac{1-2ui}{(1-ui)^2} \\ &= \frac{1-2ui}{1-u^2-2ui} \\ &= \frac{(1-2ui)(1-u^2+2ui)}{(1-u^2-2ui)(1-u^2+2ui)} \\ &= \frac{1+3u^2+2u^3i}{(1+u^2)^2}, \end{split}$$

where *i*is the imaginary unit. Then,

$$|\varphi(u)| = \frac{\sqrt{(1+3u^2)^2 + 4u^6}}{(1+u^2)^2}.$$

However, numerical computations show that

$$|\varphi(0.1)| \approx 1.00971, \quad |\varphi(0.2)| \approx 1.03561, \quad |\varphi(0.7)| \approx 1.15467.$$

This is impossible for a characteristic function, because its modulus should be less than or equal to 1 for every $u \in \mathbb{R}$. Therefore, there is no random variable \tilde{V} , independent from \tilde{W} , that meets $\tilde{W} + \tilde{V} \sim \text{Exponential (1)}$, i.e., (57). We emphasize that it is not necessary for \tilde{W} to have a discrete part; we can take an absolutely continuous random variable \tilde{W}_2 that is sufficiently near to (61) in distribution, such that $\varphi_{\tilde{W}_2}(u)$ tends to $\varphi_{\tilde{W}}(u)$ (by Lévy's continuity theorem [37], Section 18.1), and still have $|\varphi(u)| > 1$. This example is finished.

By taking (14) into account, we define the associated integral operator to (40) as

$$\mathcal{J}x(t) = t\mathbb{E}[x(tV)],\tag{62}$$

where $t \in [0, T]$, $x : [0, T] \to \mathbb{C}$ is a continuous function, and *V* is a random variable, independent from *W* and with support in [0, 1], that satisfies (52) (we assume this condition, as in the L-fractional case or Examples 2 and 3, for instance, in general, we would need Conjecture 1).

With (62), there is a version of the fundamental theorem of calculus.

Theorem 1. If $x : [0, T] \to \mathbb{C}$ is a continuously differentiable function, then

$$\mathcal{J} \circ \mathcal{D}x(t) = x(t) - x(0) \tag{63}$$

and

$$\mathcal{D} \circ \mathcal{J} x(t) = x(t)$$
 (64)

for every $t \in [0, T]$. If x is given by a power series (49), then

$$\mathcal{J}x(t) = \sum_{n=0}^{\infty} x_n \mathcal{J}t^n = \sum_{n=0}^{\infty} x_n \frac{t^{n+1}}{(n+1)\mathbb{E}[\mathbb{W}^n]}.$$
(65)

Proof. First, we note that $\mathcal{D}x \in \mathcal{C}[0, T]$ and $\mathcal{J}x \in \mathcal{C}^1[0, T]$. Indeed, if $t_1 \to t_2$ in [0, T], then $x'(t_1W) \to x'(t_2W)$, and since $|| x' ||_{\infty} < \infty$ on [0, T], the dominated convergence theorem gives $\mathcal{D}x(t_1) \to \mathcal{D}x(t_2)$. For $\mathcal{J}x \in \mathcal{C}^1[0, T]$, note that $t \mapsto \mathbb{E}[x(tV)]$ is $\mathcal{C}^1[0, T]$, because

$$\left|\frac{\mathrm{d}}{\mathrm{d}t}x(tV)\right| = |Vx'(tV)| \le ||x'||_{\infty},$$

which has finite expectation; the dominated convergence theorem is employed to ensure that

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathbb{E}[x(tV)] = \mathbb{E}\left[\frac{\mathrm{d}}{\mathrm{d}t}x(tV)\right].$$

Thus, the conclusion of this remark is that both compositions in (63) and (64) are welldefined when $x \in C^1[0, T]$. Now, we aim at proving the equalities in (63) and (64). For (63), by using the properties of the conditional expectation,

$$\mathcal{I} \circ \mathcal{D}x(t) = t\mathbb{E}[\mathcal{D}x(Vt)]$$
$$= t\mathbb{E}[\mathbb{E}_{W}[x'(WVt)]]$$
$$= t\mathbb{E}[x'(WVt)]$$
$$= t\mathbb{E}[x'(Ut)]$$
$$= t\int_{0}^{1} x'(ut)du$$
$$= \int_{0}^{t} x'(s)ds$$
$$= x(t) - x(0).$$

On the other hand, for (64),

$$\mathcal{D} \circ \mathcal{J}x(t) = \mathcal{D}[t\mathbb{E}_{V}[x(tV)]]$$

= $\mathbb{E}[\mathbb{E}_{V}[x(tVW)]] + t\mathbb{E}[W\mathbb{E}_{V}[Vx'(tVW)]]$
= $\mathbb{E}[x(tU)] + t\mathbb{E}[Ux'(tU)]$
= $\int_{0}^{1} x'(ut)du + t \int_{0}^{1} ux'(ut)du$
= $\frac{1}{t} \int_{0}^{t} x(s)ds + \frac{1}{t} \int_{0}^{t} sx'(s)ds$
= $\frac{1}{t} \int_{0}^{t} x(s)ds + \frac{1}{t} \left(tx(t) - \int_{0}^{t} x(s)ds\right)$
= $x(t).$

Integration by parts has been used in the last line. Finally, for (65), observe that \mathcal{J} is a continuous operator from $\mathcal{C}[0, T]$ into $\mathcal{C}[0, T]$: if

$$\| x - y \|_{\infty} \to 0$$

then

$$\parallel \mathcal{J}x - \mathcal{J}y \parallel_{\infty} \leq T \parallel x - y \parallel_{\infty} \to 0.$$

Hence, (65) holds, considering that

$$\mathcal{J}t^n = t\mathbb{E}[t^n V^n] = t^{n+1} \frac{\mathbb{E}[U^n]}{\mathbb{E}[W^n]} = \frac{t^{n+1}}{(n+1)\mathbb{E}[W^n]}$$

We give a version of Picard's theorem for (48), based on fixed-point theory.

Theorem 2. If $f : [0, T] \times \Omega \subseteq [0, T] \times \mathbb{R}^d \to \mathbb{R}^d$ is continuous and is locally Lipschitz continuous with respect to the second variable (i.e., for every $(t_1, x_1) \in [0, T] \times \Omega$, it is Lipschitz on a neighborhood of it), then the integral problem associated with the differential Equation (48),

$$x(t) = x_0 + \mathcal{J}f(t, x(t)) = x_0 + t\mathbb{E}[f(tV, x(tV))],$$
(66)

presents local existence and uniqueness of solution in the set of continuous functions. Specifically, if

$$Q = [0, a] \times B(x_0, b) \subseteq [0, T] \times \Omega$$

is a rectangle where f is M-Lipschitz and $\overline{B}(x_0, b)$ *is the closed ball of* $\|\cdot\|_{\infty}$ *-radius b centered at* x_0 *, then an interval of definition of the solution is* $[0, T^*]$ *, such that* $T^* \leq a$ *and*

$$T^* < \min\left\{\frac{b}{\|f\|_{\infty}}, \frac{1}{M}\right\},\$$

where $\| f \|_{\infty}$ is the maximum of f on Q. Finally, if the solution of (66) is continuously differentiable, then it solves (48) locally, by Theorem 1.

Proof. Consider the Banach space

$$\Theta = \{ y \in \mathcal{C}[0, T^*] : \| y - x_0 \|_{\infty} \le b \}.$$

Let

$$\Lambda: \Theta \to \Theta,$$

$$\Lambda x(t) = x_0 + t \mathbb{E}[f(tV, x(tV))].$$

$$f(t_1V, x(t_1V)) \to f(t_2V, x(t_2V))$$

almost surely by continuity, and since *f* is bounded on $[0, T] \times x([0, T^*])$, the dominated convergence theorem ensures that

$$\Lambda x(t_1) \to \Lambda x(t_2),$$

that is, $\Lambda x \in C[0, T^*]$. On the other hand, if $t \in [0, T^*]$, then

$$|\Lambda x(t) - x_0| \le t \parallel f \parallel_{\infty} \le b;$$

therefore,

$$\| \Lambda x - x_0 \|_{\infty} \leq b$$

and $\Lambda x \in \Theta$. Hence, Λ is well defined. Let us see that Λ is a contraction: if $_1x, _2x \in \Theta$, then

$$\begin{aligned} |\Lambda[_{1}x](t) - \Lambda[_{2}x](t)| &\leq t \mathbb{E}[|f(tV,_{1}x(tV)) - f(tV,_{2}x(tV))|] \\ &\leq Mt \mathbb{E}[||_{1}x(tV) - _{2}x(tV)|] \\ &\leq MT^{*} \|_{1} |x - _{2}x| \|_{\infty}. \end{aligned}$$

 $MT^* < 1$, Λ is a contraction and, by Banach fixed-point theorem ([38], Chapter 1), it has a unique fixed point $x \in \Theta$. \Box

4.4. Power-Series Solutions of Nonlinear Equations: The Cauchy–Kovalevskaya Theorem

We work in the context of Sections 4.1 and 4.2. We do not need the integral operator \mathcal{J} from the previous part.

We focus on scalar problems with the polynomial vector field:

$$\mathcal{D}x = \sum_{i=0}^{m} a_i x^i,\tag{67}$$

where $a_0, ..., a_m \in \mathbb{R}$. For example, a logistic vector field would be $\mu x(1 - x/K)$, $\mu \in \mathbb{R}$, and K > 0, with quadratic nonlinearity m = 2, which has previously been studied in the Caputo sense [11].

Theorem 3. Equation (67) has a power-series solution that converges on a neighborhood $[0, \epsilon)$.

Proof. By using a power series of the form (49) and Proposition 1, we have

$$\sum_{n=0}^{\infty} x_{n+1}(n+1)\mathbb{E}[W^n]t^n = \sum_{n=0}^{\infty} \left(\sum_{i=0}^m a_i C_i(x_0, \dots, x_n) \right) t^n,$$

where $C_i(x_0, ..., x_n)$ are *n*-th terms of Cauchy products of power *i*. The coefficients then satisfy

$$x_{n+1} = \frac{1}{(n+1)\mathbb{E}[W^n]} \sum_{i=0}^m a_i C_i(x_0, \dots, x_n), \quad x_0 = x(0),$$

where $C_i(x_0, ..., x_n)$ are *n*-th terms of Cauchy products of power *i*. We know that there exists a constant K > 0 such that

$$\frac{1}{(n+1)\mathbb{E}[W^n]} \le K$$

for all $n \ge 0$, by the stronger condition (43). Then, by the triangular inequality,

$$|x_{n+1}| \leq K \sum_{i=0}^{m} |a_i| C_i(|x_0|, \dots, |x_n|).$$

The formal majorizing series $\psi(t) = \sum_{n=0}^{\infty} y_n t^n$ has coefficients

$$y_{n+1} = C \sum_{i=0}^{m} |a_i| C_i(y_0, \dots, y_n), \quad y_0 = |x_0|.$$

We have the following functional, algebraic equation:

$$\psi(t) = y_0 + t \sum_{n=0}^{\infty} y_{n+1} t^n = y_0 + Kt \sum_{i=0}^{m} |a_i| \psi(t)^i.$$

With

$$\phi(t,w) = w - y_0 - Kt \sum_{i=0}^m |a_i| w^i, \quad (t,w) \in \mathbb{R}^2,$$

one has $\phi(0, y_0) = 0$ and $\frac{\partial \phi}{\partial w}(0, y_0) = 1 \neq 0$. The implicit-function theorem ([39], Section 8, Chapter 0) shows that there exists a unique analytic function $w \equiv w(z)$ on a neighborhood of zero such that $w(0) = y_0$ and $\phi(z, w(z)) = 0$. This implies that $w(z) = \psi(z)$ is analytic at 0. \Box

Example 5. Fractional differential equations have been used to extend customary epidemic models [40]. We consider a generalization of the susceptible–infected–recovered (SIR) model,

$$\begin{cases} \mathcal{D}S = -\beta SI, \\ \mathcal{D}I = \beta SI - \gamma I, \\ \mathcal{D}R = \gamma I, \end{cases}$$
(68)

where $\beta > 0$ and $\gamma > 0$ control the infection and the recovery rates, respectively. With the form (49) and the differentiation rule from Proposition 1, one has

$$\mathcal{D}S(t) = \sum_{n=0}^{\infty} s_{n+1}(n+1)\mathbb{E}[W^n]t^n,$$
$$\mathcal{D}I(t) = \sum_{n=0}^{\infty} i_{n+1}(n+1)\mathbb{E}[W^n]t^n,$$
$$\mathcal{D}R(t) = \sum_{n=0}^{\infty} r_{n+1}(n+1)\mathbb{E}[W^n]t^n.$$

On the other hand, we have the Cauchy product

$$S(t)I(t) = \sum_{n=0}^{\infty} \left(\sum_{m=0}^{n} s_m i_{n-m} \right) t^n.$$

Then,

$$s_{n+1} = -\frac{1}{(n+1)\mathbb{E}[W^n]}\beta \sum_{k=0}^n s_k i_{n-k},$$
(69)

$$i_{n+1} = \frac{1}{(n+1)\mathbb{E}[W^n]} \left[\beta \sum_{k=0}^n s_k i_{n-k} - \gamma i_n \right],$$
(70)

$$r_{n+1} = \frac{1}{(n+1)\mathbb{E}[W^n]}\gamma i_n \tag{71}$$

are the coefficients of the solution. The proof of convergence with (69)–(71) is very similar to Theorem 3, now adapted to the case of a system of three equations [41]. Based on (43), we bound

$$|s_{n+1}| \le C\beta \sum_{k=0}^{n} |s_k| |i_{n-k}|,$$
$$|i_{n+1}| \le C\beta \sum_{k=0}^{n} |s_k| |i_{n-k}| + C\gamma |i_n|,$$
$$|r_{n+1}| \le C\gamma |i_n|,$$

where C > 0 is a constant. Then we consider the majorizing sequences $\tilde{s}_0 = |s_0|$, $\tilde{i}_0 = |i_0|$, $\tilde{r}_0 = |r_0|$,

$$\widetilde{s}_{n+1} = C\beta \sum_{k=0}^{n} \widetilde{s}_{k} \widetilde{i}_{n-k},$$
$$\widetilde{i}_{n+1} = C\beta \sum_{k=0}^{n} \widetilde{s}_{k} \widetilde{i}_{n-k} + C\gamma \widetilde{i}_{n},$$
$$\widetilde{r}_{n+1} = C\gamma \widetilde{i}_{n}.$$

If we define the formal power series

$$\widetilde{S}(t) = \sum_{n=0}^{\infty} \widetilde{s}_n t^n,$$
$$\widetilde{I}(t) = \sum_{n=0}^{\infty} \widetilde{i}_n t^n,$$
$$\widetilde{R}(t) = \sum_{n=0}^{\infty} \widetilde{r}_n t^n,$$

then

$$\widetilde{S}(t) = |s_0| + Ct\beta \widetilde{S}(t)\widetilde{I}(t),$$

$$\widetilde{I}(t) = |i_0| + Ct \Big(\beta \widetilde{S}(t)\widetilde{I}(t) + \gamma \widetilde{I}(t)\Big),$$

$$\widetilde{R}(t) = |r_0| + Ct\gamma \widetilde{I}(t).$$

The analytic function

$$\begin{split} \phi : \mathbb{R}^4 \to \mathbb{R}^3, \\ (t, x, y, z) \mapsto (x - |s_0| - Ctxy, y - |i_0| - Ct(\beta xy + \gamma y), z - |r_0| - Cty) \\ \phi \Big(0, \widetilde{s}_0, \widetilde{i}_0, \widetilde{r}_0 \Big) = 0, \end{split}$$

satisfies

$$\phi\left(0,\widetilde{s}_{0},i_{0},\widetilde{r}_{0}
ight)=0,$$

 $J\phi\left(0,\widetilde{s}_{0},\widetilde{i}_{0},\widetilde{r}_{0}
ight)=1,$

where *J* is the Jacobian with respect to (x, y, z). By the implicit-function theorem, there are analytic functions x(t), y(t), and z(t) in a neighborhood of t = 0 such that $\phi(t, x(t), y(t), z(t)) = 0$ and $x(0) = \tilde{s}_0$, $y(0) = \tilde{i}_0$, $z(0) = \tilde{r}_0$. Then, $\tilde{S} = x$, $\tilde{I} = y$ and $\tilde{R} = z$ are analytic functions, as wanted. Once this model (68) is well-defined, further extensions of it could be investigated, such as the incorporation of a Brownian perturbation giving rise to a stochastic fractional differential equation.

5. Conclusions and Open Problems

The main novelties of the presented paper were:

- The definition of normalized fractional operators with non-singular kernel, for the first time, in Section 3. In the literature, there have been documented deficiencies of fractional operators with bounded kernels, see [22]. Our work showed that a rescaling fixes the issues of these operators: inconsistency at zero, units time⁰, and lack of fundamental theorem of calculus. Our new operators and equations are mathematically valid and could be further investigated, beyond exponential and Mittag–Leffler kernels. During the time the present paper has been in the preprint server ArXiv and under review, new works citing my suggested normalized operators are being conducted [32].
- The definition of a general class of fractional operators, based on a probabilistic approach, in Section 4. My previous articles [25,26] dealt with the L-fractional derivative, defined as the normalization of the Caputo derivative. The L-fractional derivative was relevant as it offered alternative properties: units time⁻¹ instead of time^{-α}, smoothness of solutions, finite ordinary derivative at the origin, etc. These properties were reviewed in Section 2. The novel Section 4 generalized, for the first time, the L-fractional derivative and the normalization of the Caputo operator via probability theory by defining fractional operators with an averaged probabilistically distributed past. Many properties were stated and demonstrated: the validity and consistency of the definition, the associated Mittag–Leffler function, existence and uniqueness of solution by fixed-point theory, and an example related to the SIR model.

In addition to the future research lines suggested in [25], some open problems from the present paper are the following:

- The study of more properties of rescaled operators with bounded kernels. See Section 3.
- The development of more theory on operators with memory (40). Specifically, it would be of great relevance to obtain a complete resolution of Conjecture 1. This would better characterize when the fundamental theorem of calculus and existence-and-uniqueness results hold; see Section 4.3.
- The study of the new Mittag–Leffler-type function (51): representation formulas, dynamics, asymptotic values, etc.
- The investigation of dynamical systems based on the new operator (40).
- The design of numerical methods for (48). The search for applications in modeling.

Funding: This research received no external funding.

Data Availability Statement: Data is contained within the article.

Conflicts of Interest: The author declares no conflicts of interest.

References

- 1. Podlubny, I. Fractional Differential Equations—An Introduction to Fractional Derivatives, Fractional Differential Equations, to Methods of Their Solution and Some of Their Applications, 1st ed.; Academic Press: Cambridge, MA, USA, 1998; Volume 198.
- Kilbas, A.A.; Srivastava, H.M.; Trujillo, J.J. Theory and Applications of the Fractional Differential Equations; North-Holland Mathematics Studies; Elsevier: Amsterdam, The Netherlands, 2006.
- 3. Diethelm, K. The Analysis of Fractional Differential Equations—An Application-Oriented Exposition Using Differential Operators of Caputo Type; Lecture Notes in Mathematics; Springer: Berlin/Heidelberg, Germany, 2010.
- 4. Abbas, S.; Benchohra, M.; Lazreg, J.E.; Nieto, J.J.; Zhou, Y. *Fractional Differential Equations and Inclusions*; Classical and Advanced Topics; World Scientific: Singapore, 2023.
- 5. Yong, Z. Basic Theory of Fractional Differential Equations, 3rd ed.; World Scientific: Singapore, 2023.
- 6. Ascione, G.; Mishura, Y.; Pirozzi, E. *Fractional Deterministic and Stochastic Calculus*; De Gruyter Series in Probability and Stochastics; Walter de Gruyter: Berlin, Germany, 2024; Volume 4.
- 7. Oliveira, E.C.D.; Machado, J.A.T. A review of definitions for fractional derivatives and integral. *Math. Probl. Eng.* 2014, 2014, 238459. [CrossRef]
- 8. Ortigueira, M.D.; Machado, J.T. What is a fractional derivative? J. Comput. Phys. 2015, 293, 4–13. [CrossRef]
- 9. Teodoro, G.S.; Machado, J.T.; Oliveira, E.C.D. A review of definitions of fractional derivatives and other operators. *J. Comput. Phys.* **2019**, *388*, 195–208. [CrossRef]

- 10. Diethelm, K.; Kiryakova, V.; Luchko, Y.; Machado, J.T.; Tarasov, V.E. Trends, directions for further research, and some open problems of fractional calculus. *Nonlinear Dyn.* **2022**, *107*, 3245–3270. [CrossRef]
- 11. Area, I.; Nieto, J.J. Power series solution of the fractional logistic equation. Phys. A 2021, 573, 125947. [CrossRef]
- 12. D'Ovidio, M.; Lai, A.C.; Loreti, P. Solutions of Bernoulli equations in the fractional setting. Fractal Fract. 2021, 5, 57. [CrossRef]
- 13. Jornet, M. Power-series solutions of fractional-order compartmental models. *Comput. Appl. Math.* 2024, 43, 67. [CrossRef]
- 14. Kexue, L.; Jigen, P. Laplace transform and fractional differential equations. Appl. Math. Lett. 2011, 24, 2019–2023. [CrossRef]
- 15. Garrappa, R. On linear stability of predictor-corrector algorithms for fractional differential equations. *Int. J. Comput. Math.* **2010**, *87*, 2281–2290. [CrossRef]
- 16. Caputo, M. Linear model of dissipation whose Q is almost frequency independent-II. Geophys. J. Int. 1967, 13, 529–539. [CrossRef]
- 17. Beghin, L.; Cristofaro, L.; Garrappa, R. Renewal processes linked to fractional relaxation equations with variable order. *J. Math. Anal. Appl.* **2024**, *531*, 127795. [CrossRef]
- Hazarika, D.; Borah, J.; Singh, B.K. Existence and controllability of non-local fractional dynamical systems with almost sectorial operators. J. Math. Anal. Appl. 2024, 532, 127984. [CrossRef]
- Fernandez, A.; Buranay, S.C. The Peano-Sard theorem for Caputo fractional derivatives and applications. *J. Comput. Appl. Math.* 2024, 441, 115705. [CrossRef]
- Fernandez, A.; Al-Refai, M. A rigorous analysis of integro-differential operators with non-singular kernels. *Fractal Fract.* 2023, 7, 213. [CrossRef]
- Onitsuka, M.; L-Fassi, I.I.E. Generalized Caputo-Fabrizio fractional differential equation. J. Appl. Anal. Comput. 2024, 14, 964–975. [CrossRef]
- 22. Diethelm, K.; Garrappa, R.; Giusti, A.; Stynes, M. Why fractional derivatives with nonsingular kernels should not be used. *Fract. Calc. Appl. Anal.* **2020**, *23*, 610–634. [CrossRef]
- Lazopoulos, A.K.; Karaoulanis, D. On L-fractional derivatives and L-fractional homogeneous equations. *Int. J. Pure Appl. Math.* 2016, 21, 249–268.
- 24. Lazopoulos, K.A.; Karaoulanis, D.; Lazopoulos, A.K. On fractional modelling of viscoelastic mechanical systems. *Mech. Res. Commun.* **2016**, *78*, 1–5. [CrossRef]
- 25. Jornet, M. Theory on Linear L-Fractional Differential Equations and a New Mittag-Leffler-Type Function. *Fractal Fract.* 2024, *8*, 411. [CrossRef]
- 26. Jornet, M.; Nieto, J.J. Power-series solution of the L-fractional logistic equation. Appl. Math. Lett. 2024, 154, 109085. [CrossRef]
- 27. Lazopoulos, K.A.; Lazopoulos, A.K. Equilibrium of Λ-fractional liquid crystals. Mech. Res. Commun. 2024, 136, 104243. [CrossRef]
- 28. Mainardi, F. Why the Mittag-Leffler function can be considered the queen function of the fractional calculus? *Entropy* **2020**, 22, 1359. [CrossRef]
- 29. Caputo, M.; Fabrizio, M. A new definition of fractional derivative without singular kernel. Prog. Fract. Differ. Appl. 2015, 1, 73–78.
- 30. Losada, J.; Nieto, J.J. Properties of a new fractional derivative without singular kernel. Prog. Fract. Differ. Appl. 2015, 1, 87–92.
- Losada, J.; Nieto, J.J. Fractional integral associated to fractional derivatives with nonsingular kernels. *Prog. Fract. Differ. Appl.* 2021, 7, 137–143.
- 32. Al-Refai, M.; Baleanu, D. On new solutions of the normalized fractional differential equations. *Fractals* **2024**, *32*, 2450115. [CrossRef]
- 33. Atangana, A.; Baleanu, D. New fractional derivatives with non-local and non-singular kernel. Theory and application to heat transfer model. *Therm. Sci.* 2016, 20, 763–769. [CrossRef]
- 34. Area, I.; Nieto, J.J. Fractional-order logistic differential equation with Mittag-Leffler-type Kernel. *Fractal Fract.* **2021**, *5*, 273. [CrossRef]
- 35. Al-Refai, M.; Baleanu, D. On an extension of the operator with Mittag-Leffler kernel. Fractals 2022, 30, 2240129. [CrossRef]
- 36. Jornet, M. Theory and methods for random differential equations: A survey. SeMA J. 2023, 80, 549–579. [CrossRef]
- 37. Williams, D. Probability with Martingales; Cambridge University Press: Cambridge, UK, 1991.
- 38. Granas, A.; Dugundji, J. Fixed Point Theory; Springer: New York, NY, USA, 2003.
- 39. Kaup, L.; Kaup, B. Holomorphic Functions of Several Variables: An Introduction to the Fundamental Theory; Walter de Gruyter: Berlin, Germany; New York, NY, USA, 2011; Volume 3.
- Area, I.; Batarfi, H.; Losada, J.; Nieto, J.J.; Shammakh, W.; Torres, Á. On a fractional order Ebola epidemic model. *Adv. Differ. Equ.* 2015, 1, 1–12. [CrossRef]
- 41. Srivastava, H.M.; Area, I.; Nieto, J.J. Power-series solution of compartmental epidemiological models. *Math. Biosci. Eng.* **2021**, *18*, 3274–3290. [CrossRef] [PubMed]

Disclaimer/Publisher's Note: The statements, opinions and data contained in all publications are solely those of the individual author(s) and contributor(s) and not of MDPI and/or the editor(s). MDPI and/or the editor(s) disclaim responsibility for any injury to people or property resulting from any ideas, methods, instructions or products referred to in the content.