



Article Calcination Temperature Reflected Structural, Optical and Magnetic Properties of Nickel Oxide

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Abstract: Stoichiometric compositions of NiO were prepared by the standard chemical co-precipitation method to inspect the effect of the calcination temperature on structures, morphology, and physical properties. The samples were calcined at three different temperatures viz. $350 \,^{\circ}$ C, $550 \,^{\circ}$ C, and $650 \,^{\circ}$ C for 5 h. X-ray diffraction analysis confirmed the cubic (*Fm-3m*) structure of the prepared samples. The average crystalline size increases from 41 nm to above 100 nm as the calcination temperature increases in the same time period. In Fourier transform infrared spectra, the spectral absorption bands were observed at ~413, 434, and 444 cm⁻¹. The bandgap energy of NiO particles is decreased from 3.6 eV to 3.41 eV as the calcination temperature increases. The magnetic analysis confirms that the magnetization value of NiO is invariably decreased with a rise in the calcination temperature.

Keywords: nickel oxide; ceramics; material characterization; optical properties; bandgap engineering



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Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). 1. Introduction

Theory and experimental advances in ferromagnetic semiconductors offer anticipation for a new generation of microelectronic devices. In recent days, nanoscale ferromagnetic semiconductor oxides are gaining significant attention for various potential applications, such as spin-polarized carrier sources, a catalyst for oxygen evolution, and easy integration into semiconductor technology [1–3]. The study of small particles has become of increasing interest due to the new properties in those materials shown when the crystal size is reduced.

Nickel oxide (NiO) in bulk form, crystallized in rhombohedral structure with space group R3m with antiferromagnetic ordering [4]. NiO nanoparticles—a NaCl-type cubic crystalline structure (space group, Fm3m) antiferromagnetic semiconductor—are a transition metal oxide of great interest due to its magnetic properties [5,6] and well dispersion, and are desirable for many applications in designing ceramic, magnetic, and heterogeneous catalytic materials [7]. It was found that the critical annealing temperature for NiO nanoparticles in the air and oxygen atmospheres is ~600 °C [8,9].

To prepare NiO nanoparticles, various synthesis methods were adopted earlier, such as sol-gel [10], impregnation [11], spray pyrolysis [12], microwave [13], oxidation [14], coprecipitation [15], sonochemistry [16], and ion exchange [17]. Some of the above methods suffer from the difficulty in size-homogeneity and well dispersion of NiO nanoparticles. To obtain better properties of NiO, the synthesis conditions must be well controlled to obtain particularly fine powders with a narrow particle size distribution that enhance the performance of materials. It has been reported that NiO nanoparticles show hysteresis below a blocking temperature, irreversibility between the field cooled and zero field cooled magnetization curves [18,19], large coercivities [20], weak ferromagnetism and superparamagnetism, and spin-glass behavior [21,22]. It is found that the calcination temperature significantly affects the characteristic properties of NiO [23,24]. The magnetic properties of the nanostructured nickel oxide are found to be markedly different from those of single-crystalline [23]. It was reported that the increase in calcination temperature resulted in an increase in particle size with a high surface area [25]. Understanding the size-dependent magnetization is not only important for practical applications but also it is of great interest to identify the unprecedented nature of the antiferromagnetic oxide nanostructures between the structure and properties [26].

Motivated through the above, in the present work, we investigate the effect of different calcination temperatures on the structural, optical, and magnetic properties of NiO particles. We have used the chemical co-precipitation method to synthesize the NiO particles. Furthermore, the grown materials were characterized by X-ray powder diffraction (XRD), scanning electron microscopy (SEM), Raman spectroscopy, Fourier transform infrared (FT-IR) spectroscopy, X-ray photoelectron spectroscopy, UV–Vis spectroscopy, and magnetic measurement.

2. Materials and Methods

The chemical co-precipitation method was used for the synthesis of polycrystalline NiO samples using a high purity compound of Ni(NO₃)₂·6H₂O. In order to prepare aqueous solutions of nickel salt, 10 gm of Ni(NO₃)₂·6H₂O were dissolved in 500 mL of distilled water. This aqueous solution was freshly prepared in an appropriate molar ratio. After that, the solution was kept at 80 °C and 6 M NaOH was added to the solution. As a result, the pH value of the solution rapidly increased to 12. The particles settled at the bottom were collected and the top water layer with excess salts was discarded. The particles were washed repeatedly with distilled water to remove salt impurities. Later on, the washed particles were dried at room temperature and calcined at different temperatures for 5 h. The synthesized samples of NiO with different calcination temperatures, such as $350 \,^{\circ}$ C, $550 \,^{\circ}$ C, and $650 \,^{\circ}$ C for 5 h were further designated as N-1, N-2, and N-3, respectively

Phase purity and crystal structure of the samples at room temperature were identified using a Bruker D8 Advance X-ray diffractometer. The X-ray diffraction (XRD) data were collected for the range $20^{\circ} < 2\theta < 80^{\circ}$ with a step size of 0.02 and scanning speed of 2° per 50 s. The X-rays were produced (40 kV and 40 mA) using a sealed tube and the wavelength of the X-ray was 1.5406 A (CuK α). The surface morphology and microstructure of the grown ceramics were studied with a scanning electron microscope (SEM) (LEO 430 electron microscope Polaron SC7640 Sputter Coater (gold palladium)). Raman measurements on as-synthesized samples were carried out on a Jobin-Yovn Horiba Labram (System HR800) spectrometer with a 488 nm excitation source equipped with a Peltier cooled charge-coupled device detector. Fourier transform infrared (FT-IR) spectra were recorded in the frequency range of 4000–400 cm⁻¹ employing the KBr pellets technique using Bruker Germany make spectrometer model vertex-70. Furthermore, XPS measurements on the samples were carried out using a VSW ESCA machine equipped with AlK α radiation. The diffuse reflectance spectra of grown samples were measured using the UV-Vis spectrometer (Perkin Elmer, Lambda 950, Waltham, MA, USA). The bandgap energy was calculated from the diffuse reflection spectra according to the Kubelka–Munk theory. The magnetization measurements were performed using a vibrating sample magnetometer (VSM-magnetic analyzer).

3. Results and Discussion

3.1. Structural and Microstructural Properties

Figure 1 shows the XRD patterns of N-1, N-2, and N-3 samples. The peaks' positions appearing at 20 is 37.21°, 43.22°, 63.10°, 75.20°, and 79.39° can be readily indexed as (111), (200), (220), (311), and (222) crystal planes of the NiO, respectively. The main peaks are comparable to that of the standard powder diffraction spectrum of NiO (JCPDS card No. 78-0429) and those reported in the literature for cubic NiO [27,28]. The XRD

pattern shows that the samples are single phase and no other impurities distinct from the diffraction peak except the characteristic peaks of FCC phase NiO were detected. The diffraction peaks are indexed in the cubic phase of NiO with space group *Fm-3m*. Structural analysis confirms the cubic (*Fm-3m*) phase for all the three compounds with the lattice parameters (a = b = c = 4.16 Å; $\alpha = \beta = \gamma = 90^{\circ}$).



Figure 1. Room temperature XRD patterns of N-1, N-2, and N-3 samples are compared with standard powder diffraction of NiO (JCPDS card No. 78-0429).

It can be seen from Figure 1 that the diffraction peaks are low intense and broad at low calcination temperature. The broadening of the peaks indicated that the particles were of nanometer scale. The crystallite size diameter of the NiO sample has been determined from the broadening of the most intense peak (200) in the XRD pattern using Debye-Scherrer's equation:

$$D = \frac{0.9\lambda}{\beta\cos\theta} \tag{1}$$

where λ is the X-ray wavelength, β is the full width of the diffraction line at half its maximum intensity, and θ is the Bragg angle. The crystallite size was found to be 41, 132, 191 nm for N-1, N-2, and N-3, respectively. This is attributed that as we have increased the calcination temperature of the samples, crystallite size also has been increased.

Figure 2 shows the SEM images of N-1, N-2, and N-3 samples. These images reveal a systematic evaluation of the morphology of the prepared samples. For all three NiO samples, the images show aggregate particles. As the particles are aggregated, hence it is difficult to do a size distribution analysis. It is observed that the particles are spherical-shaped nanocrystallites. Previously, it was reported that the crystallite size was found to increase with increasing calcination temperatures [24,25].





Figure 2. SEM micrograph of NiO calcined at a temperature of (A) 350 °C, (B) 550 °C, and (C) 650 °C.

3.2. Raman Analysis

Raman scattering is a useful probe to study the presence of lattice defects in solids, such as structure transition, lattice distortion, cation distribution, and magnetic ordering. The room temperature Raman spectra of pure NiO consists of several bands within the range of 200 to 1400 cm^{-1} , and is shown in Figure 3. The Raman spectrum exhibited a strong, broad peak at ~494–499 cm⁻¹ due to the Ni-O stretching mode [29]. Vibration modes at around 400 cm⁻¹, 499 cm⁻¹ were attributed to one-phonon (1P) and transverse optical (TO) mode, respectively. The frequencies of the LO modes have a blue shift (494–499 cm⁻¹) with an increasing calcination temperature of nanocrystals. This phenomenon was reported previously and can be ascribed to the size-induced phonon confinement effect in corroboration with the strength of the two-phonon coupling and surface relaxation [30–32]. This confirms that the size of NiO nanocrystals has a strong effect on their magnetic properties. Some artifacts are seen in the Raman spectra of N-2 and N-3 samples. These changes are not related to samples and can be related to an instrumental error or some other factor.



Figure 3. Raman scattering spectra of N-1, N-2, and N-3 particles carried out at room temperature.

3.3. FT-IR Spectroscopy

Figure 4 depicts the FT-IR spectrum of N-1, N-2, and N-3 nanomaterials in the range of mid-IR. The broad absorption peak at about 3446 cm⁻¹ is attributed to the O-H stretching vibration of interlayer water molecules and H-bound of OH group and the band near 1621 cm⁻¹ is assigned to H-O-H bending vibrations. The band at 1384 cm⁻¹ is the characteristic of NO³⁻ interlayer [33,34]. The small presence of nitrates in the samples may worsen the characteristic behavior of NiO. The IR band at 1708 cm⁻¹ can be related to NaNO₃ as a by-product of the reaction. The bands at 2923 and 2851 cm⁻¹ are due to C-H vibrations which were incorporated during the calcinations process [35]. The spectrum of the N-1, N-2, and N-3 calcined at different temperatures clearly shows the Ni-O absorption band near ~413, 434, and 444 cm⁻¹, respectively. These absorption bands' peak is undoubtedly assigned to Ni–O stretching and gave clear evidence for the presence of the crystalline behavior [27].



Figure 4. Fourier transforms infrared spectra of N-1, N-2, and N-3 samples. Inset shows the enlarged FT-IR spectra of the N-1 sample.

3.4. XPS Spectroscopy

XRD is a bulk technique that penetrates the entire sample volume. However, XPS is surface-sensitive, with the sampling depth determined by the kinetic energy of the photoelectron. The mean free path is $\sim 1.0-1.3$ nm for the ~ 400 eV (MgK α)–633 eV (AlK α) kinetic energies of the Ni 2*p* photoelectron and assuming an exponential attenuation of the photoelectron with a depth into the sample surface, 95% of the XPS signal arises from within 2.3–3.0 nm of the outermost surface layers of the NiO samples. Thus, the

XPS technique was used to investigate the surface compositions and chemical state of the samples; X-ray photoelectron spectroscopy (XPS) measurements were performed on the N-1 sample (calcined at 350 °C). The full scan of the NiO sample is presented in Figure 5a and performed in a range of binding energy of 0–1400 eV. The spectrum consists of a peak corresponding to C, O, and Ni, only showing the purity of the product. The peak with O1s binding energy around 529.17 eV corresponds to Ni-O bonds, as shown in Figure 5b. This indicates that the oxygen atoms existed as O^{2-} species in as-synthesized NiO compound.



Figure 5. X-ray photoelectron spectra of NiO sample (**a**) Full scan, (**b**) high-resolution spectra of O1s (**c**) high-resolution spectra of Ni 2*p* core level.

The Ni 2*p* region (Figure 5c) comprises four easily discernible features: the Ni 2*p*3/2 main peak and its satellite at ~854 and ~862 eV, and the Ni 2*p*1/2 main peak and its satellite at ~872 and ~879 eV, respectively, and is in agreement with the earlier reported value of NiO [36,37]. No binding energy shifts or peak shape changes are observed, and contrary to other reports for similar-sized particle systems, the nanosized NiO Ni 2*p* XPS is identical to that of well-ordered stoichiometric nickel oxide [38,39].

The structure can be analyzed in terms of a simplified model of final-state effects in which photoemission is potentially accompanied by rearrangement of electrons in the valence band:

$$2p^{6}3d^{8} + h\nu \to 2p^{5}3d^{8} + e^{-}$$
⁽²⁾

$$2p^6 3d^8 + h\nu \to 2p^5 3d^9 \underline{L} + e^- \tag{3}$$

where \underline{L} models screening by the O 2*p* valence band and represents a hole at the O 2*p* level. The main feature is generally assigned to the screened core hole photoemission (Equation (3)) with peak 854 eV ascribed to local screening from lattice oxygen adjacent to the Ni 2*p* core hole and peak 855 eV being to nonlocal screening. The satellite region is fit to peak at 861 eV and peak at 865 eV and is dominated by the nonscreened mechanism of Equation (2). Evaluating the peak position and symmetry of the Ni 3*p* line shape, one can reason that nickel was present in a 2+ oxidation state.

3.5. Optical Properties

UV-vis absorption spectroscopy is one of the important methods to reveal the energy structure and optical properties of semiconductor materials. In our measurements, the sample was analyzed through diffuse reflectance arrangement. Figure 6 shows the absorption spectra of the as-prepared NiO samples (i.e., N-1, N-2, and N-3) transformed from the diffuse reflection spectra according to the Kubelka–Munk (K–M) theory [40]. The energy bandgap of NiO particles could be estimated from the Eg measurement using Kubelka–Munk functions [41] and F(R) is estimated from the following equation:

$$F(R) = \frac{(1-R)^2}{2R}$$
(4)

where F(R) is the remission or Kubelka–Munk function, where *R* is the reflectance. A graph is plotted between $[F(R)hv]^2$ and hv, the intercept value is the bandgap energy, as shown in Figure 6. The curves indicate that the values of the direct bandgap (Eg) values are 3.60, 3.57, and 3.41 eV for N-1, N-2, and N-3, respectively. It can be seen that the bandgap values clearly decrease from 3.6 eV (N-1) to 3.41 eV (N-3) with an increase in calcination temperature, which might be ascribed to the additional sub-band-gap energy levels that are induced by the abundant surface and interface defects in the agglomerated particles [42]. The bandgap value is influenced by various factors, such as crystallite size, structural parameter, carrier concentrations, presence of impurities, and lattice strain [43].



Figure 6. Cont.



Figure 6. UV-Vis spectrum for N-1, N-2, and N-3 samples.

3.6. Magnetic Properties

We measured the magnetization (M) of all three N-1, N-2, and N-3 samples as a function of externally applied magnetic field (H) at room temperature. Figure 7 shows the set of hysteresis loops of nickel oxide (i.e., N-1, N-2, and N-3). From the M-H curve, we have found an interesting fact that the magnetization value is invariably decreased with a rise in the calcination temperature. However, the coercivity is increased gradually as the calcination temperature increases (inset of Figure 7).



Figure 7. Room temperature obtained Magnetic (M-H) curves of NiO samples calcined at 350 $^{\circ}$ C, 550 $^{\circ}$ C, and 650 $^{\circ}$ C for 5 h.

It is well known that the bulk NiO has an antiferromagnetic nature [44]. Nevertheless, despite being antiferromagnetic, the N-1 sample obtained after calcination at 350 °C for 5 h shows the weak ferromagnetic component at room temperature in the lower applied field. It is attributed to the small crystallite size and the increasing tendency of uncompensated spins. The very low magnetization value of N-2 and N-3 samples could be attributed to the observed superparamagnetism in nickel oxide, as well as the contribution towards the magnetic moment arising from other sources, such as the presence of non-stoichiometry and/or Ni³⁺ ions [45]. This means that the gradual increase in particle size of the polycrystalline NiO is the favorable condition for antiferromagnetic ordering (or spin compensation).

In summary, we have investigated the effect of calcination temperature on structural, optical, and magnetic properties of NiO. The NiO particles were synthesized by the chemical co-precipitation method. Structural characterization of the samples revealed cubic NaCl structure with *Fm-3m* space group. Structural analysis also confirmed that the samples are free from any other impurity elements. The average crystalline size increases from 41 nm to above 100 nm as the calcination temperature increases. Spectral absorption bands were observed in ~413, 434, and 444 cm⁻¹ for N-1, N-2, and N-3 samples, respectively. The bandgap energy values for N-1, N-2, and N-3 samples are 3.60, 3.57, and 3.41 eV respectively. The bandgap values decrease from 3.6 eV to 3.41 eV as the calcination temperature increases.

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